

# Collective Modes

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Feb 20, 2008



# Learning Goals

- \* Idea of probing collective modes
    - \* phonons, plasmons,...
  - \* Equations of motion (collective coordinate) approach
  - \* Virial Theorem
  - \* Sum Rules
  - \* Linear Response
- Very specific example

# Setup

N atoms in harmonic trap (anisotropic)

$$H = \sum_i \left[ \frac{p_i^2}{2m} + \sum_\alpha \frac{1}{2} m \omega_\alpha^2 r_{i\alpha}^2 \right] + \frac{U_0}{2} \sum_{i < j} \delta(\mathbf{r}_i - \mathbf{r}_j)$$

Change omega's and watch response

Measure in an experiment:

$$Q_\alpha = \sum_j r_{j\alpha}^2$$

(breathing modes,  
quadrupole modes)

Goal today:  
calculate frequencies

# Equation of motion approach

$$\dot{Q}_\alpha = \frac{1}{i\hbar} [Q_\alpha, H] = \frac{1}{m} \sum_i r_{i\alpha} p_{i\alpha} + p_{i\alpha} r_{i\alpha}$$

$$\ddot{Q}_\alpha = \frac{1}{i\hbar} [\dot{Q}_\alpha, H] = \frac{4}{m} \left( T_\alpha - V_\alpha + \frac{1}{2} U \right)$$

$$T_\alpha = \sum_i \frac{p_{i\alpha}^2}{2m} \quad V_\alpha = \frac{1}{2} m \omega_\alpha^2 Q_\alpha \quad U = \frac{1}{2} U_0 \sum_i \delta(r_i - r_j)$$

(AKA collective coordinates)

# Equation of motion approach

$$\dot{Q}_\alpha = \frac{1}{i\hbar} [Q_\alpha, H] = \frac{1}{m} \sum_i r_{i\alpha} p_{i\alpha} + p_{i\alpha} r_{i\alpha} = \text{"virial"}$$

$$\ddot{Q}_\alpha = \frac{1}{i\hbar} [\dot{Q}_\alpha, H] = \frac{4}{m} \left( T_\alpha - V_\alpha + \frac{1}{2} U \right)$$

$$T_\alpha = \sum_i \frac{p_{i\alpha}^2}{2m} \quad V_\alpha = \frac{1}{2} m \omega_\alpha^2 Q_\alpha \quad U = \frac{1}{2} U_0 \sum_i \delta(r_i - r_j)$$

(AKA collective coordinates)

# Virial Theorem

$$\ddot{Q}_\alpha = \frac{1}{i\hbar} [\dot{Q}_\alpha, H] = \frac{4}{m} \left( T_\alpha - V_\alpha + \frac{1}{2} U \right)$$

In equilibrium  $Q$  is time independent

$$\langle T_\alpha \rangle - \langle V_\alpha \rangle + \frac{1}{2} \langle U \rangle = 0$$

Also useful to note

$$H = \sum_{\alpha} T_{\alpha} + \sum_{\alpha} V_{\alpha} + U$$

# Closing the equations

$$\ddot{Q}_\alpha = \frac{1}{i\hbar} [\dot{Q}_\alpha, H] = \frac{4}{m} \left( T_\alpha - V_\alpha + \frac{1}{2} U \right)$$

**Case 1:  $U=0$**

$$H_\alpha = T_\alpha + V_\alpha \quad \text{is constant of motion}$$

$$\ddot{Q}_\alpha = -4\omega_\alpha^2 \left[ Q_\alpha - \frac{H_\alpha}{m\omega_\alpha^2} \right]$$

**Shift Q:**  $\tilde{Q}_\alpha = Q_\alpha - \frac{H_\alpha}{m\omega_\alpha^2}$

$$\ddot{\tilde{Q}}_\alpha = -4\omega_\alpha^2 \tilde{Q}_\alpha \quad \text{Oscillates at twice trap frequency}$$



# Closing the equations

$$\ddot{Q}_\alpha = \frac{1}{i\hbar} [\dot{Q}_\alpha, H] = \frac{4}{m} \left( T_\alpha - V_\alpha + \frac{1}{2} U \right)$$

Case 2:  $T=0$  (good approximation for a BEC)

$$U = H - \sum_{\beta} V_{\beta}$$

$$\frac{d^2}{dt^2} \begin{pmatrix} Q_x \\ Q_y \\ Q_z \end{pmatrix} = \frac{2}{m} H - \begin{pmatrix} 3\omega_x^2 & \omega_x\omega_y & \omega_x\omega_z \\ \omega_x\omega_y & 3\omega_y^2 & \omega_y\omega_z \\ \omega_x\omega_z & \omega_y\omega_z & 3\omega_z^2 \end{pmatrix} \begin{pmatrix} Q_x \\ Q_y \\ Q_z \end{pmatrix}$$

constant  
(shift away)

eigenvalues give oscillation frequencies

Q: if all omegas are equal,  
what are oscillation frequencies?

$$\frac{d^2}{dt^2} \begin{pmatrix} Q_x \\ Q_y \\ Q_z \end{pmatrix} = \frac{2}{m} H - \begin{pmatrix} 3\omega_x^2 & \omega_x\omega_y & \omega_x\omega_z \\ \omega_x\omega_y & 3\omega_y^2 & \omega_y\omega_z \\ \omega_x\omega_z & \omega_y\omega_z & 3\omega_z^2 \end{pmatrix} \begin{pmatrix} Q_x \\ Q_y \\ Q_z \end{pmatrix}$$

constant  
(shift away)

eigenvalues give oscillation frequencies

# Closing the equations

Generically the equations do not close so easily  
-- introduce more formal tools for dealing with it

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Changing trap constants for short time:

$$H_{\text{pert}} = \lambda_{\alpha}(t)Q_{\alpha}$$

**Response**

$$\begin{aligned}\langle Q_{\beta}(t) \rangle &= \langle T e^{i \int^t d\tau (H_0 + H_{\text{pert}})} Q_{\beta} e^{-i \int^t d\tau (H_0 + H_{\text{pert}})} \rangle \\ &= \langle Q_{\beta} \rangle_0 + \int dt' \chi_{\alpha\beta}^R(t - t') \lambda_{\alpha}(t') + \dots\end{aligned}$$

$$\chi_{\alpha\beta}^R = \frac{\theta(t)}{i} \langle [Q_{\beta}(t), Q_{\alpha}(0)] \rangle_0$$

# Interaction picture

(only if requested)

$$U = T e^{-i \int^t d\tau (H_0 + H')}$$

$$\bar{U} = e^{iH_0 t} U$$

$$\begin{aligned} i\partial_t \bar{U} &= e^{iH_0 t} H' U \\ &= H'(t) \bar{U} \end{aligned}$$

so

$$\bar{U}(t) \approx 1 - i \int^t d\tau H'(\tau)$$

# Equations of motion

$$\chi_{\alpha\beta}^R = \frac{\theta(t)}{i} \langle [Q_\beta(t), Q_\alpha(0)] \rangle_0$$

$$\partial_t^2 \chi_{\alpha\beta} = \frac{\delta'(t)}{i} \langle [Q_\beta(0), Q_\alpha(0)] \rangle_0 + \frac{\delta(t)}{i} \langle [\dot{Q}_\beta(0), Q_\alpha(0)] \rangle + \frac{\theta(t)}{i} \langle [\ddot{Q}_\beta(t), Q_\alpha(0)] \rangle$$

Substitute in EOM for Q

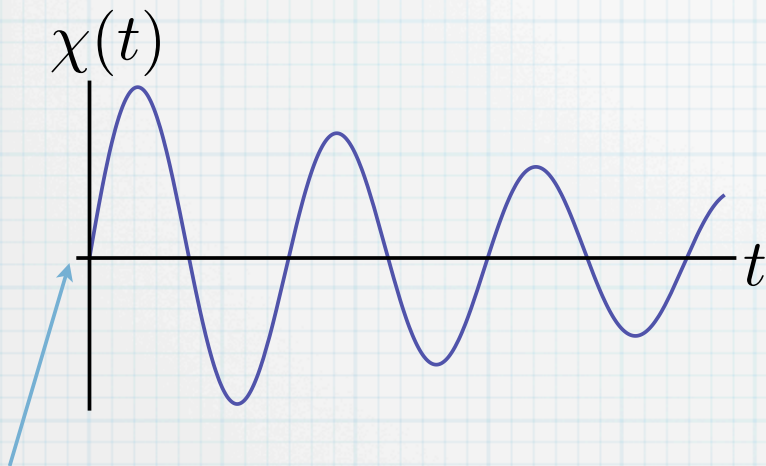
When EOM for Q close -- so do EOM for chi

$$\langle [\dot{Q}_\beta(0), Q_\alpha(0)] \rangle = \delta_{\alpha\beta} \frac{4\langle Q_\alpha \rangle}{m}$$

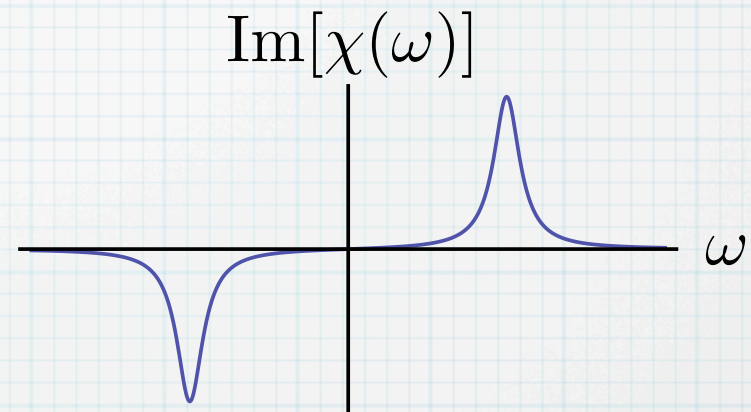
$$\dot{Q}_\alpha = \frac{1}{i\hbar} [Q_\alpha, H] = \frac{1}{m} \sum_i r_{i\alpha} p_{i\alpha} + p_{i\alpha} r_{i\alpha}$$

# Expected Structure:

Damped Harmonic oscillator



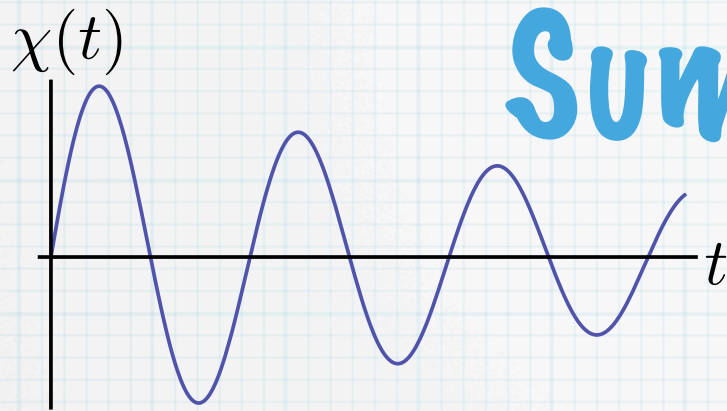
Probe and response commute  
at  $t=0$



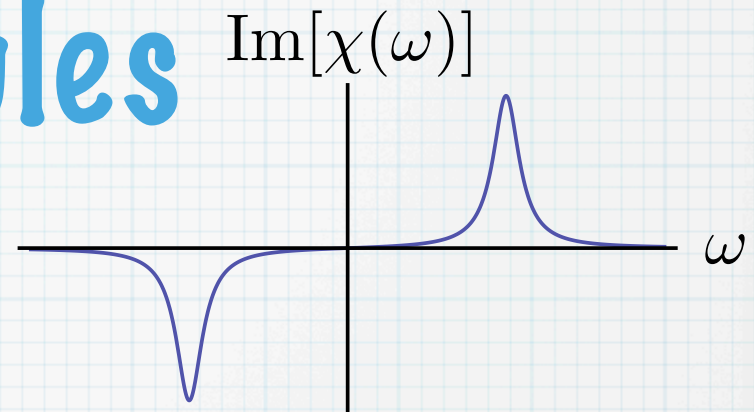
Sum rules:

Make Ansatz for  $\chi(t)$

Fit parameters from  $t=0$   
 $\chi$  and derivatives



# Sum Rules

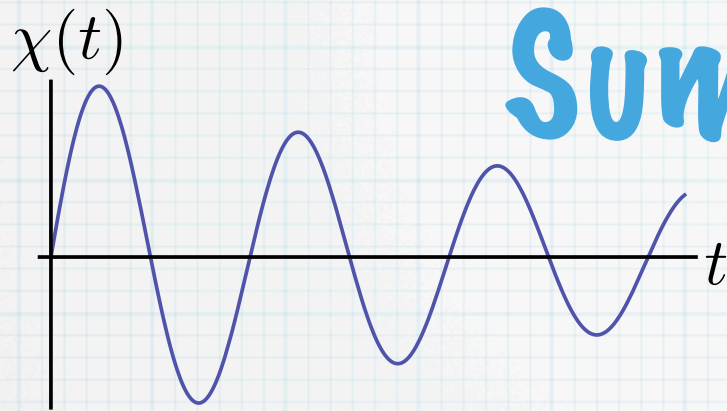


**f-sum rule: (always set by kinematics)**

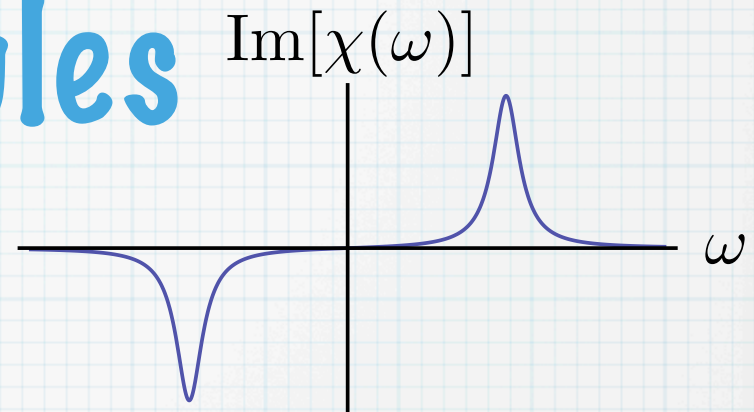
$$\begin{aligned}
 M_1 &= \int \frac{d\omega}{2\pi} (i\omega) \chi(\omega) = -i \langle [Q'(0), Q(0)] \rangle \\
 &= \frac{4}{m} \langle Q_\alpha \rangle \delta_{\alpha\beta}
 \end{aligned}$$


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$$\begin{aligned}
 M_3 &= \int \frac{d\omega}{2\pi} (i\omega)^3 \chi(\omega) = -i \langle [Q''(0), Q'(0)] \rangle \\
 &= \frac{8}{m^2} \left[ \delta_{\alpha\beta} (2\langle T_\alpha \rangle + 2\langle V_\alpha \rangle) + \frac{1}{2} \langle U \rangle \right]
 \end{aligned}$$



# Sum Rules



Estimate frequency

$$-\omega^{-2} \approx \text{eigenvalues}(M_1^{-1} M_3)$$

Other tools: Compressibility sum rule

$$M_{-1} = \lim_{t \rightarrow 0} \int \frac{d\omega}{2\pi} e^{i\omega t} \frac{\chi(\omega)}{i\omega} = \lim_{t \rightarrow 0} \int_{-\infty}^t \chi(t') dt'$$

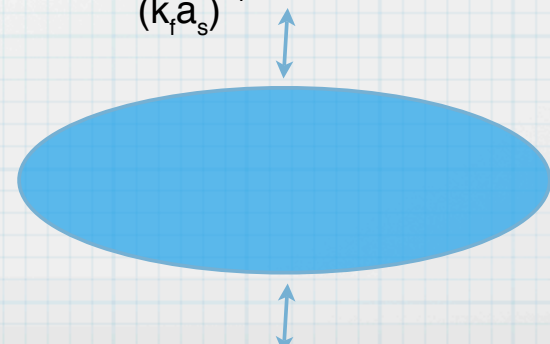
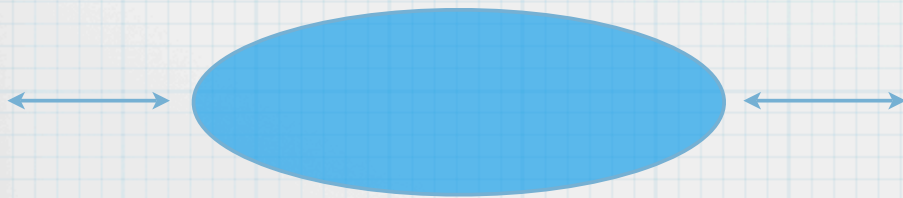
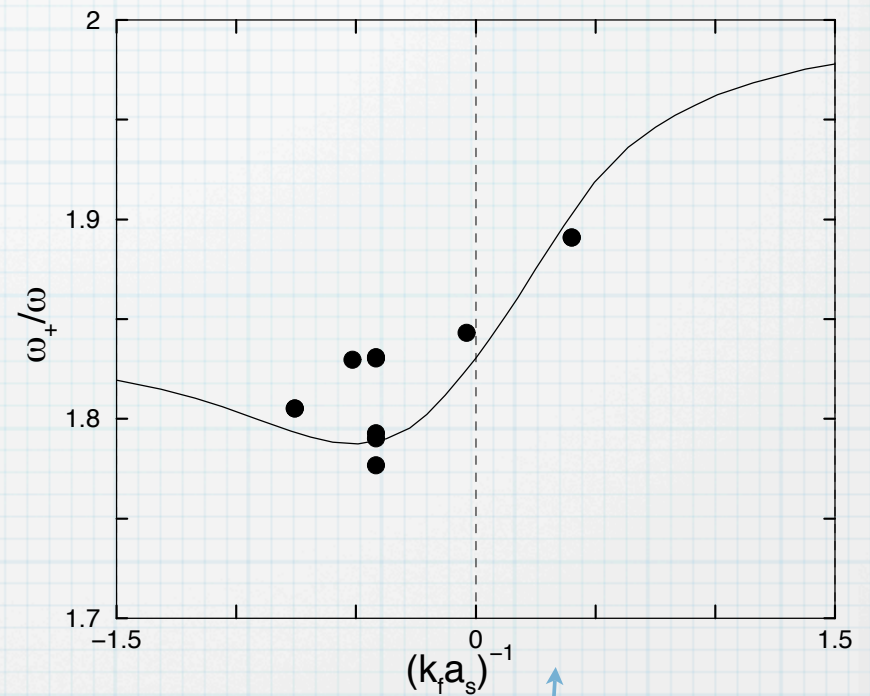
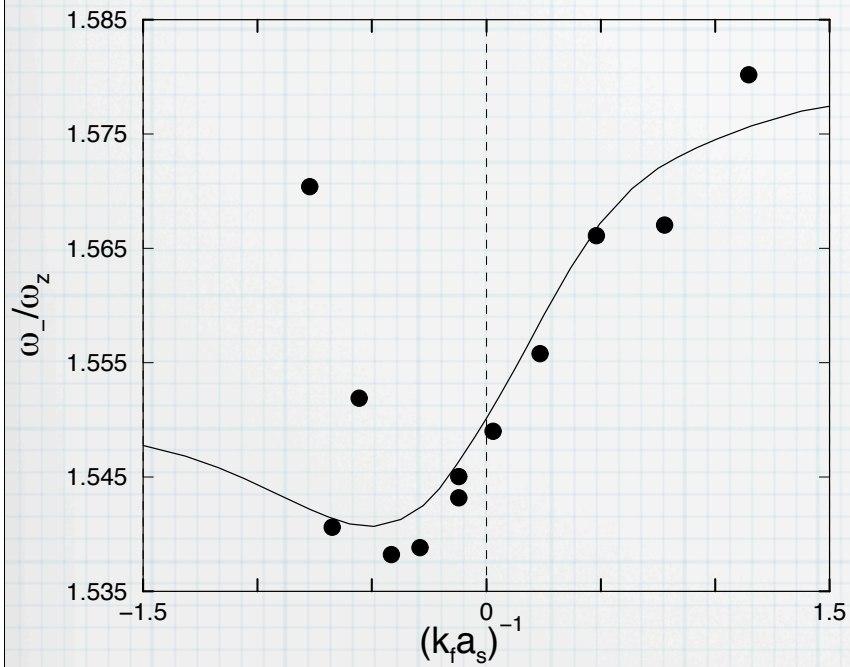
$$= \chi(\omega = 0)$$

Static Response

(or use Kramers-Kronig)



# 2-component Fermi Gas



# So What



Aside from learning that our sum rules work --  
what do we learn?

A: Modes depend on equation of state --  
learn about it

# Simpler Case

## Dipole Mode (Kohn Mode)

$$X = \sum_i x_i$$

$$P = \sum_i p_{xi}$$

Shift Trap center:  $H_{\text{pert}} = \lambda(t)X$

$$\chi = \frac{\theta(t)}{i} \langle X(t)X(0) - X(0)X(t) \rangle$$

Equations of motion:

$$\partial_t X = P/m \quad \partial_t P = -m\omega_x^2 X$$

Solution

$$X(t) = \cos(\omega t)X(0) + (m/\omega) \sin(\omega t)P(0)$$

# Response function

$$\chi = \frac{\theta(t)}{i} \langle X(t)X(0) - X(0)X(t) \rangle$$

$$\chi = \frac{m}{\omega} \theta(t) \sin(\omega t)$$

Useful to introduce “structure factors”

(we will see these next when we discuss Scattering)

$$\begin{aligned} \chi^{\gt} &= \langle X(t)X(0) \rangle \\ &= \langle X^2 \rangle \cos(\omega t) + (m/\omega) \langle XP \rangle \sin(\omega t) \\ &= \frac{m}{2\omega} (\coth(\beta\omega/2) \cos(\omega t) + i \sin(\omega t)) \end{aligned}$$

$$\begin{aligned} \chi^{\lt} &= \langle X(0)X(t) \rangle \\ &= \frac{m}{2\omega} (\coth(\beta\omega/2) \cos(\omega t) - i \sin(\omega t)) \end{aligned}$$

know about T

# Detailed Balance

$$\begin{aligned}\chi^>(t) &= \text{Tr} e^{-\beta H} e^{iHt} X e^{-iHt} X \\ &= \text{Tr} e^{-\beta H} X e^{-\beta H} e^{iHt} X e^{-iHt} e^{\beta H} \\ &= \chi^<(t + i\beta)\end{aligned}$$

Fourier Transform:

$$\chi^>(\nu) = \chi^<(-\nu) = e^{\beta\nu} \chi^>(-\nu)$$

