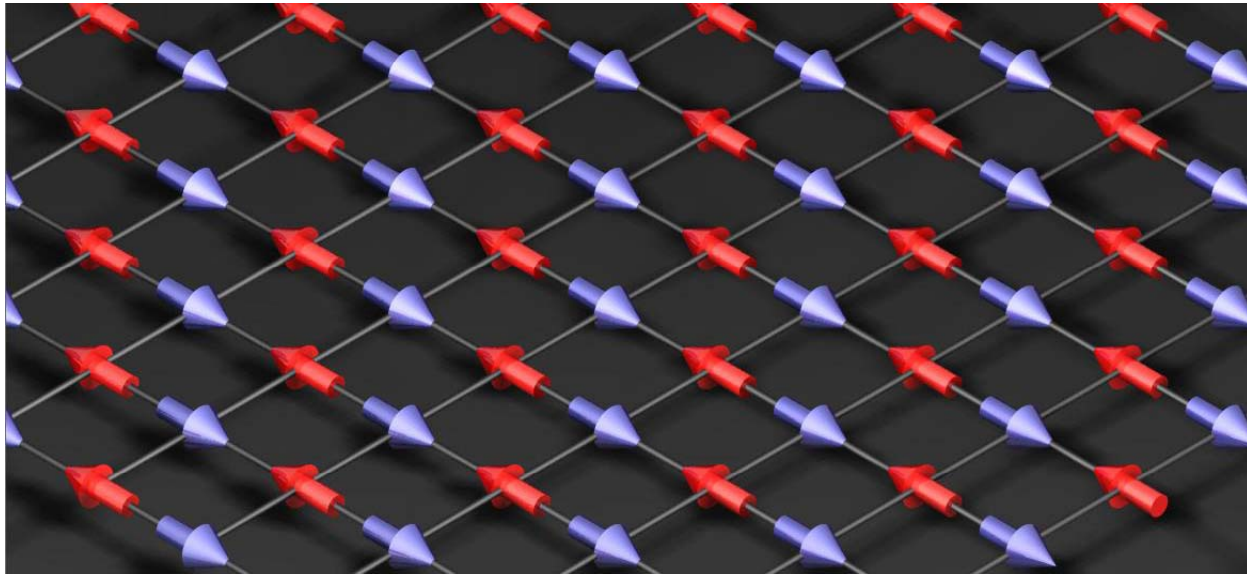


Quantum Simulations using Fermionic Atoms in Optical Lattices



Ken O'Hara

The Pennsylvania State University



Fermi Gas Group at Penn State



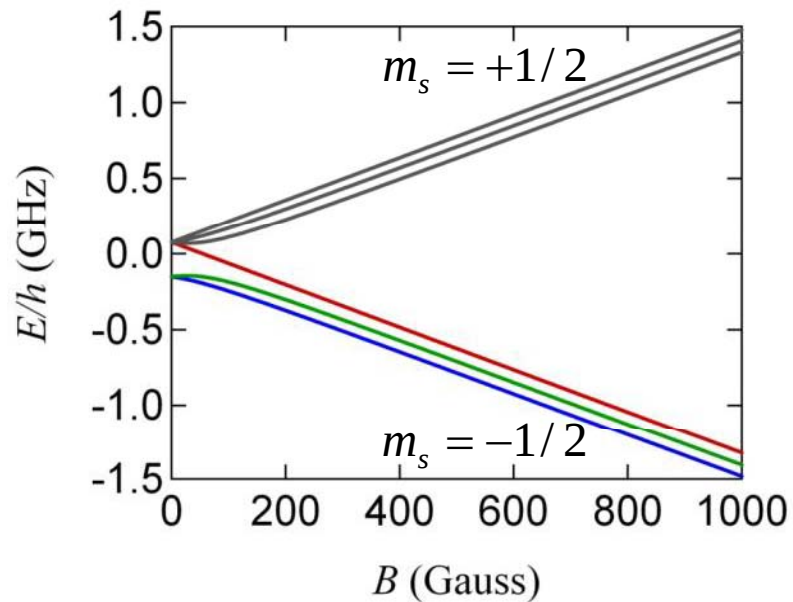
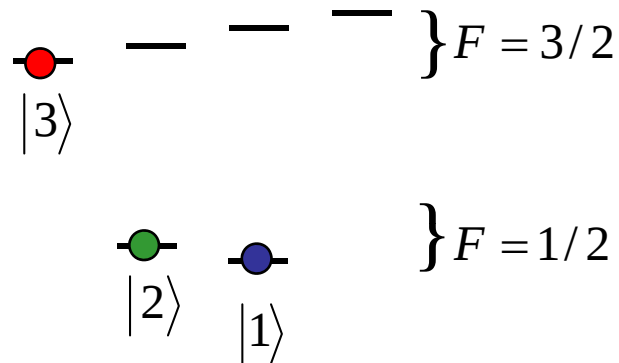
Ken O'Hara John Huckans Ron Stites Eric Hazlett Jason Williams Yi Zhang

Outline

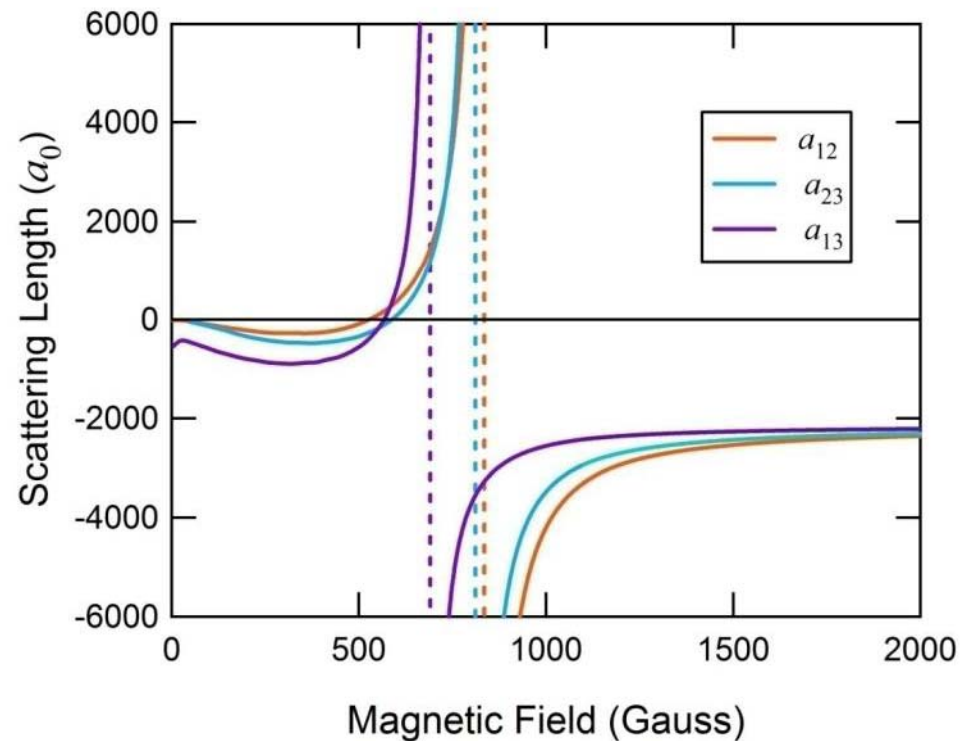
- o Ultracold Fermionic Atoms
- o The Hubbard Model
 - Quantum Magnetism
 - High-temperature superconductivity
- o Reducing the Entropy of Atoms in a Lattice
- o Three-Component Fermi Gases
 - The quantum 3-body problem
 - $SU(3)$ Hubbard model: Color Superfluidity

Two-State Mixtures in ^6Li

Hyperfine States



Feshbach Resonances

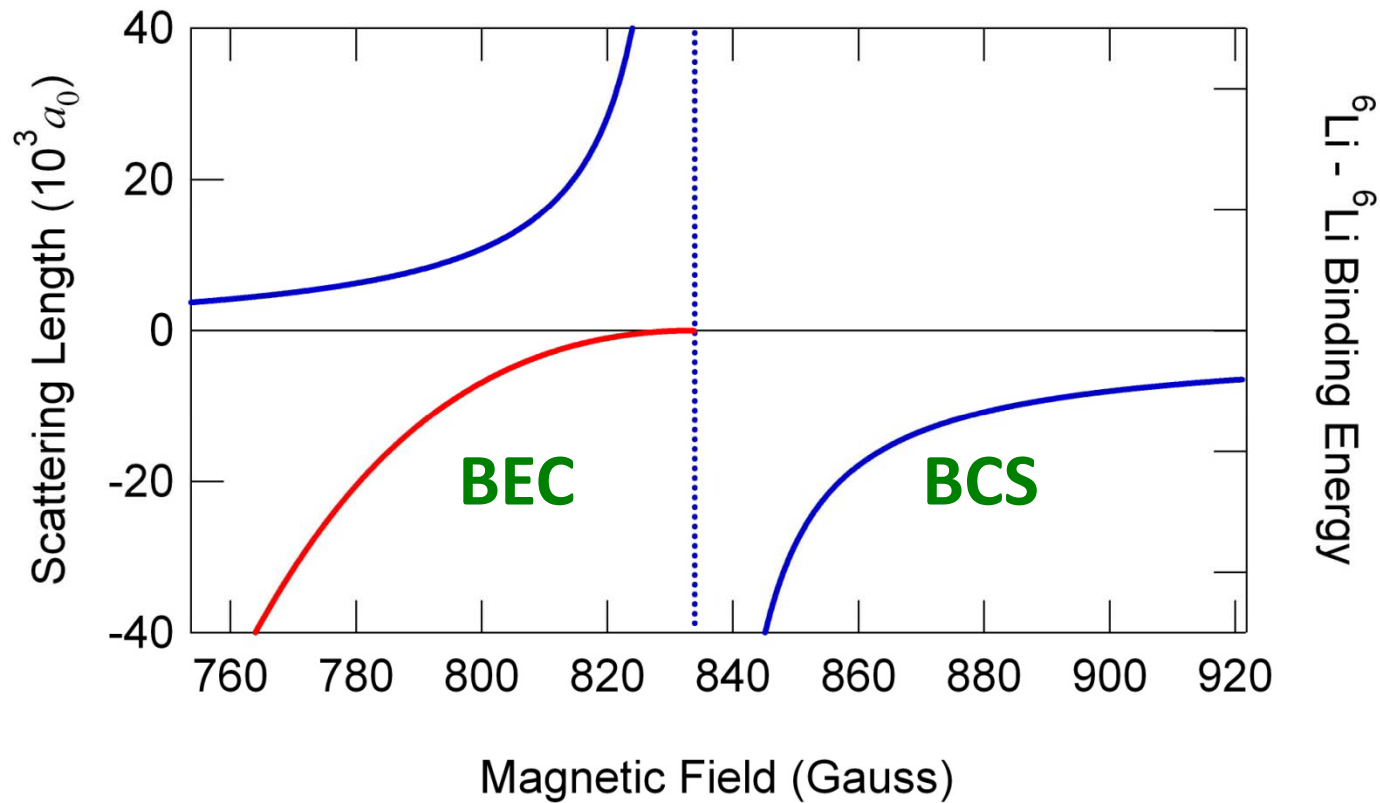


Interactions at High Field

$$a_{12}, a_{23}, a_{13} \rightarrow a_t \quad \text{as} \quad B \rightarrow \infty$$

$$a_t = -2140a_0$$

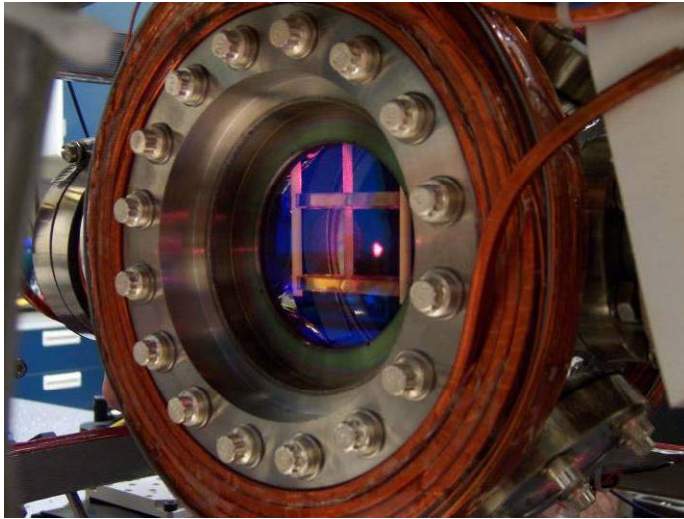
BEC-BCS Crossover with a Feshbach Resonance



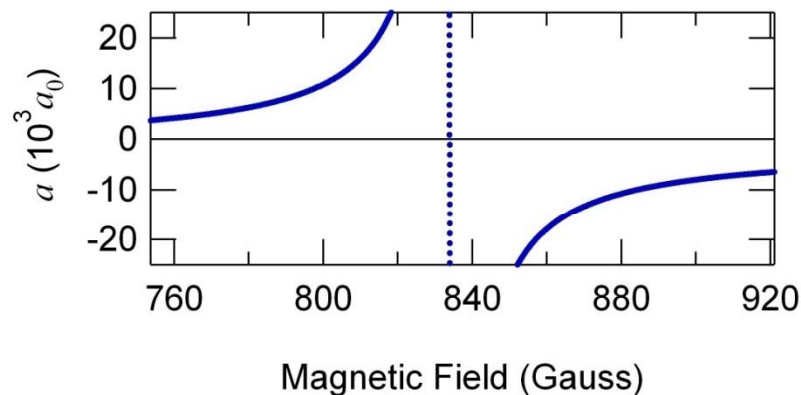
$$V(\mathbf{r}_1 - \mathbf{r}_2) = g \delta^{(3)}(\mathbf{r}_1 - \mathbf{r}_2)$$

$$g = \frac{4\pi\hbar^2 a}{m}$$

Making Degenerate Fermi Gases



- Magneto-Optical Trap
- Optical Trap
- Evaporative Cooling (at 760 Gauss)



Vertical crossed-dipole trap:
Two 80 Watt 1064 nm beams



$$U_f = 38 \mu\text{K}/\text{beam}$$

$$\nu_y = 106 \text{ Hz}$$

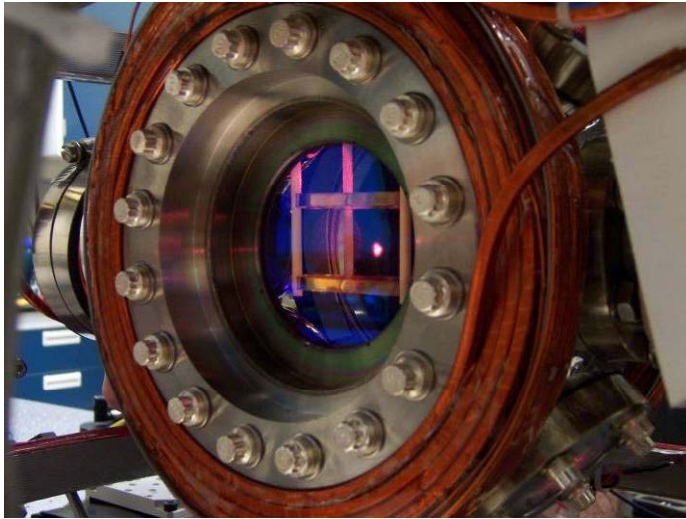
$$\nu_z = 965 \text{ Hz}$$

$$\nu_x = 3.84 \text{ kHz}$$

$$N = 2 \times 10^5 \text{ atoms}$$

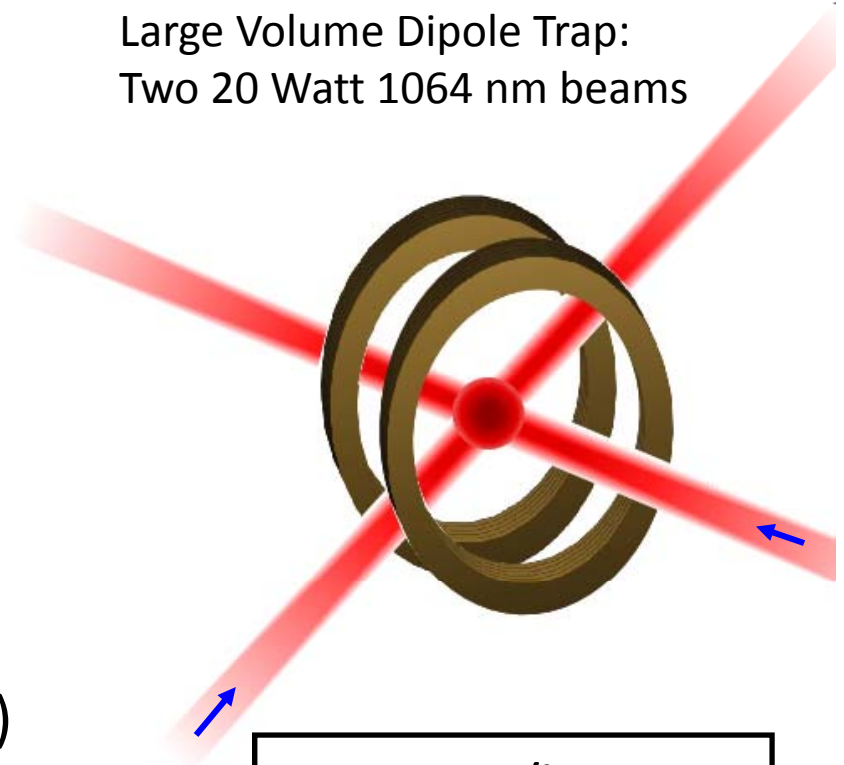
$$T_F = 3 \mu\text{K}$$

Making Degenerate Fermi Gases



- Magneto-Optical Trap
- Optical Trap
- Evaporative Cooling (at 760 Gauss)
- Transfer to Large Volume Trap

Large Volume Dipole Trap:
Two 20 Watt 1064 nm beams



$$U_f = 2 \mu\text{K}/\text{beam}$$

$$\nu_y = 21 \text{ Hz}$$

$$\nu_z = 35 \text{ Hz}$$

$$\nu_x = 92 \text{ Hz}$$

$$N = 2 \times 10^5 \text{ atoms}$$

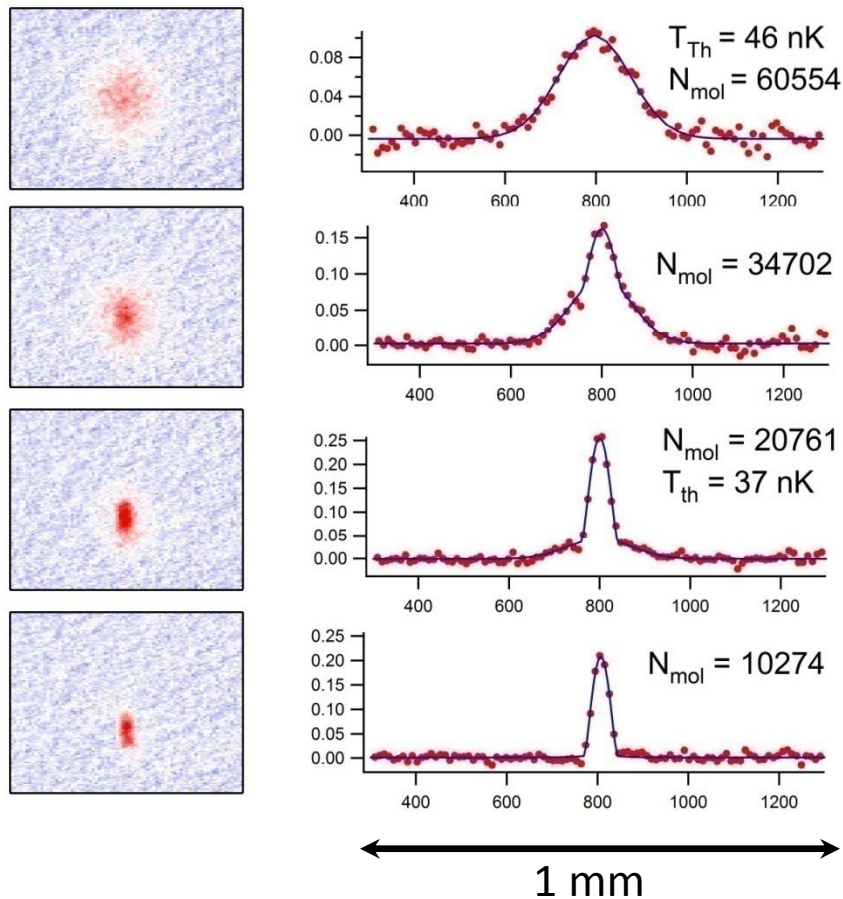
$$T_F = 160 \text{ nK}$$

BEC and DFG of ^6Li

BEC of Li_2 Molecules

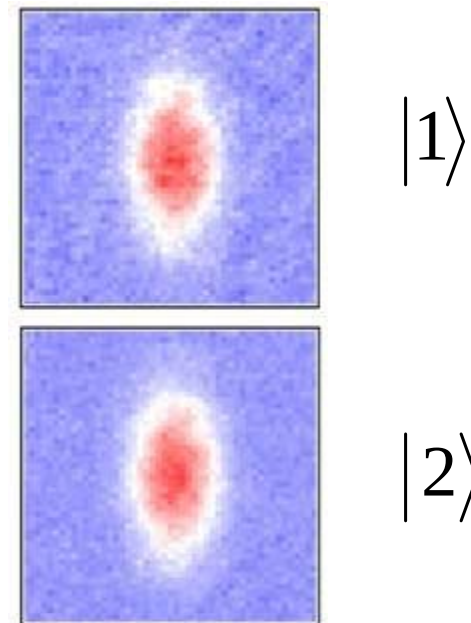
Absorption Image after Expansion

$t_{\text{TOF}} = 10\text{ms}$



2-State Degenerate Fermi Gas

Absorption Image after Expansion

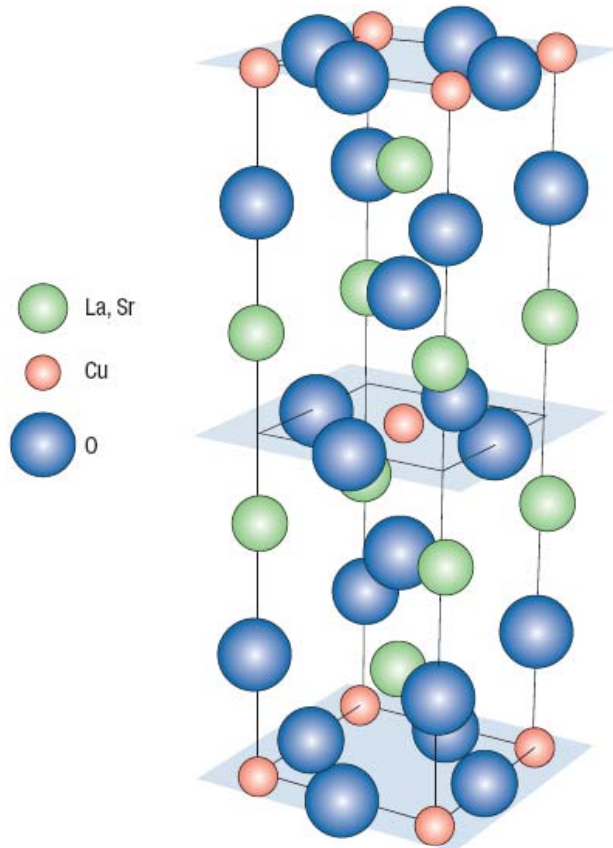


$$N_{\uparrow} = N_{\downarrow} = 10^5 \quad T \leq 0.1 T_F$$

$$T_F = 160 \text{ nK} \quad T \leq 16 \text{ nK}$$

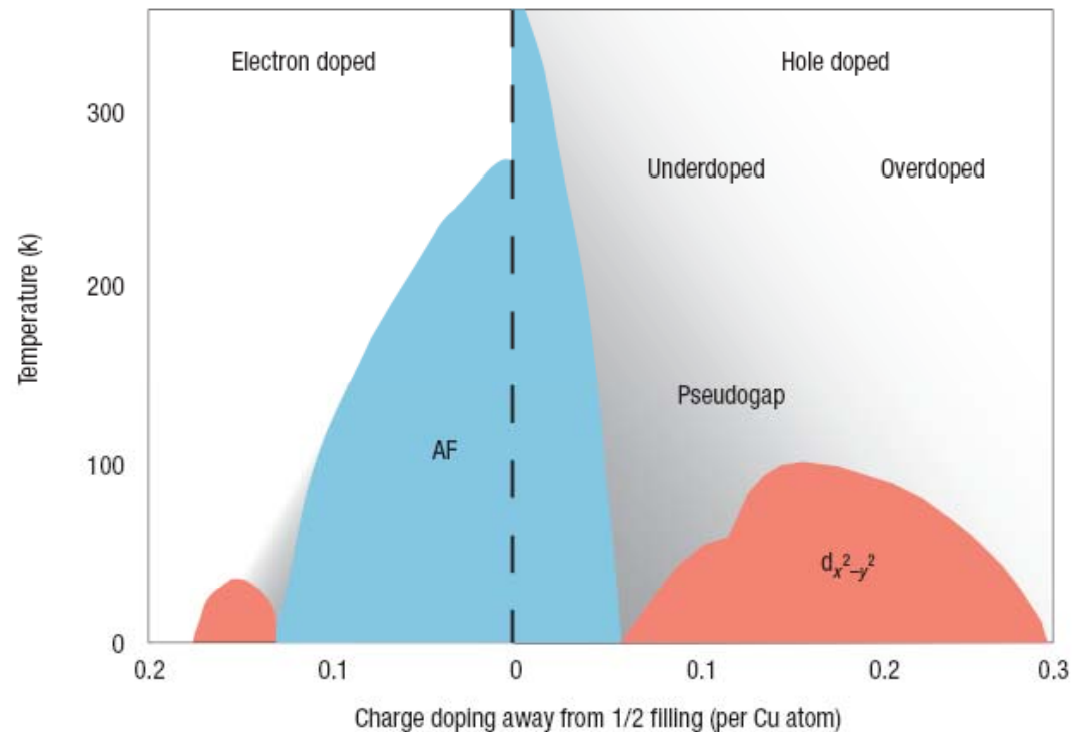
High-Temperature Superconductors

$\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$
Crystalline Structure



Weakly Coupled
Square-Planar Sheets of CuO

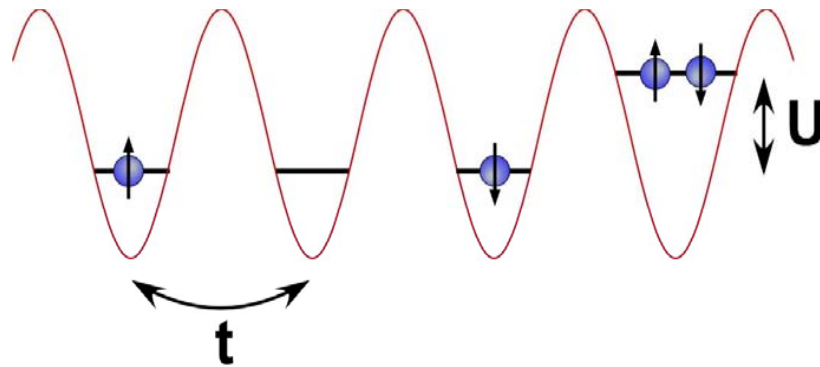
Typical Phase Diagram
of the Cuprates



Antiferromagnetism
Adjacent to Superconductivity

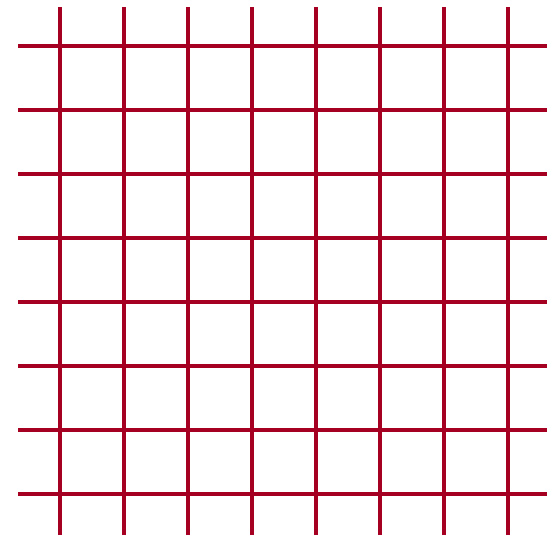
The 2D Hubbard Model

- Simplest Model that includes Band Structure & Interactions



$$H = -t \sum_{\langle i,j \rangle, \sigma} (c_{i,\sigma}^\dagger c_{j,\sigma} + h.c.) + U \sum_j n_{j,\uparrow} n_{j,\downarrow}$$

Square 2D Lattice



t = hopping term, describes tunneling between lattice sites

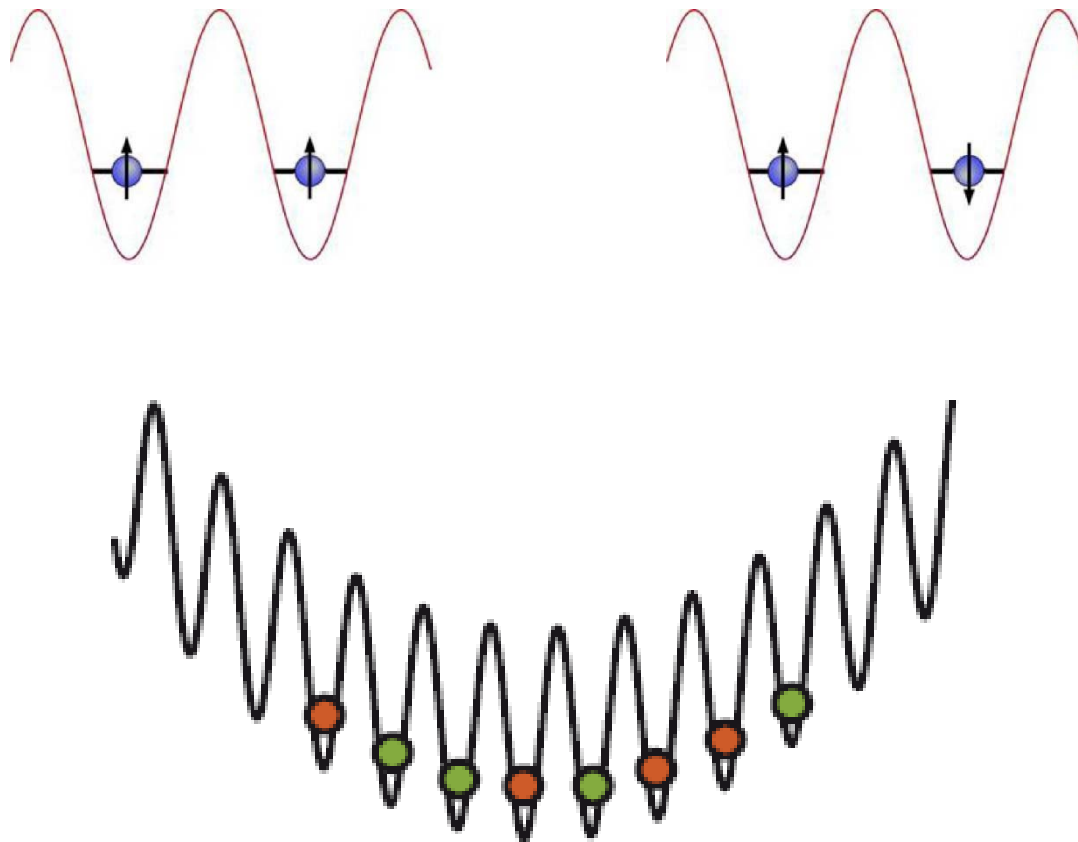
U = Interaction term, describes energy cost of filling doubly occupied sites

- Repulsive- U Hubbard Model

$U > 0$: Mimics Coulomb Repulsion between Electrons

Large U : Mott Insulator

- $U/t \rightarrow \infty$ Insulator
- Ground State: Mott Insulator

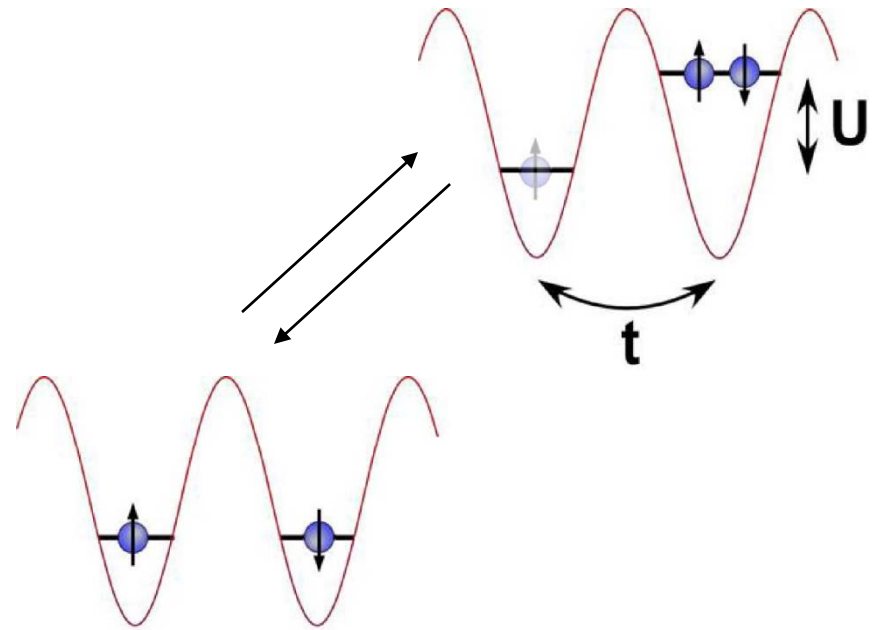
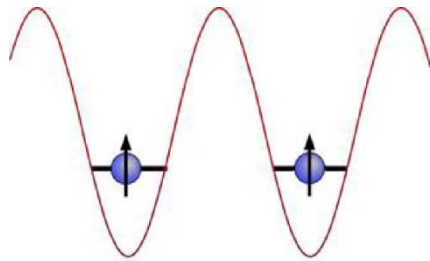


Non-zero t : Exchange Interactions

- Small, but nonzero t
- Matrix element t couples to doubly occupied sites

$$H_0 = U \sum_j n_{j,\uparrow} n_{j,\downarrow}$$

$$H_t = -t \sum_{\langle i,j \rangle, \sigma} (c_{i,\sigma}^\dagger c_{j,\sigma} + h.c.)$$

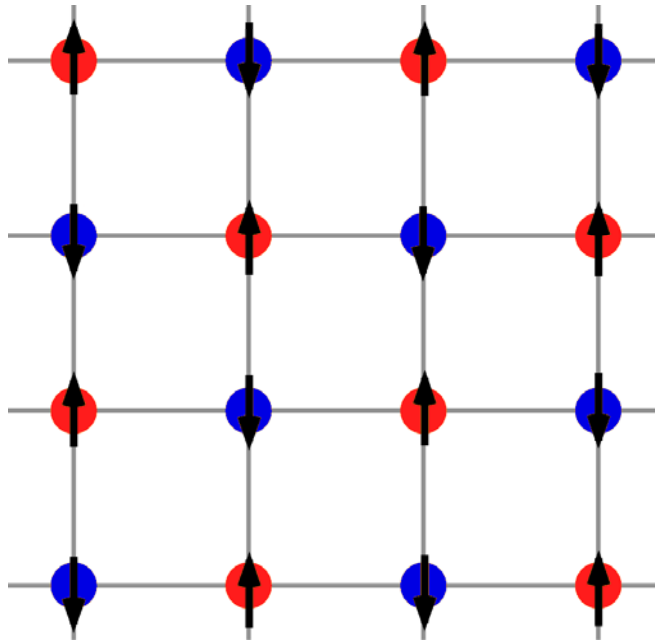


$$H^{(2)} = -\frac{\langle g | H_t | e \rangle \langle e | H_t | g \rangle}{\langle e | H_0 | e \rangle} = -\frac{t^2}{U}$$

$$H_{AF} = J \mathbf{S}_1 \cdot \mathbf{S}_2 \quad J = \frac{4t^2}{U}$$

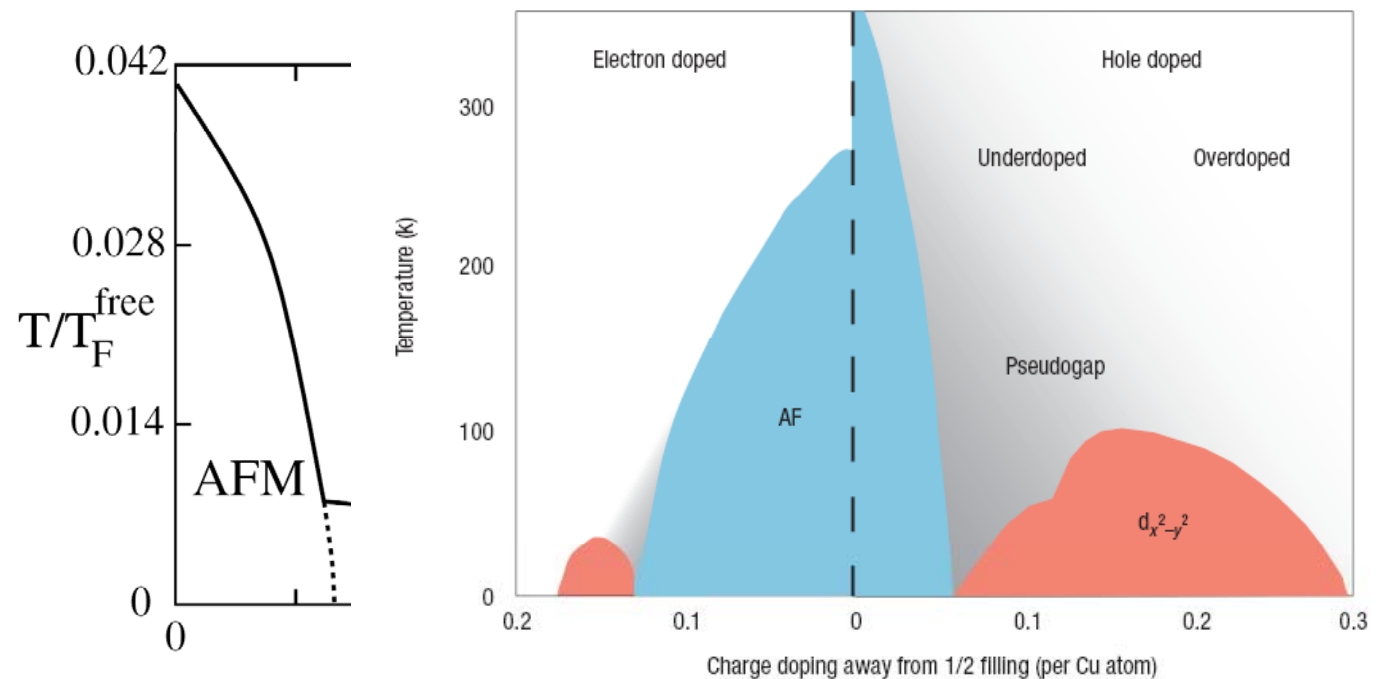
Anti-Ferromagnetic Mott Insulator

$$H_{AF} = JS_1 \cdot S_2 \qquad J = \frac{4t^2}{U}$$



Phase Diagram of the Hubbard Model

Phase Diagram in Fluctuation-Exchange Approximation



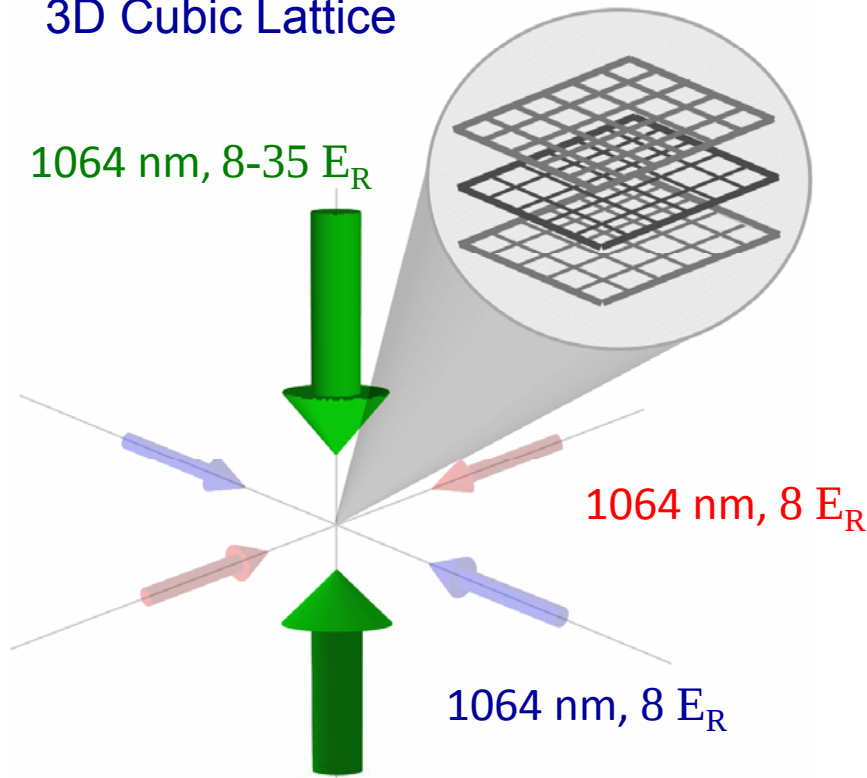
W. Hofstetter et al., PRL 89, 220407 (2002)

Realization of the Hubbard Model

Spin-1/2 Fermions:

${}^6\text{Li}$ in $F = 1/2$, $m_F = \pm 1/2$

3D Cubic Lattice



Tunneling Time ($8 E_R$): $\tau = \frac{h}{4t} = 455 \mu\text{s}$

Experimental Parameters

Scattering Length: $a_s = 63a_0$

$$\frac{t}{h} = 548 \text{ Hz}$$

$$\frac{U}{h} = 4.37 \text{ kHz}$$

$$\frac{U}{t} = 8$$

$$\frac{E_{\text{gap}}}{U} = 29$$

Temperatures of Interest

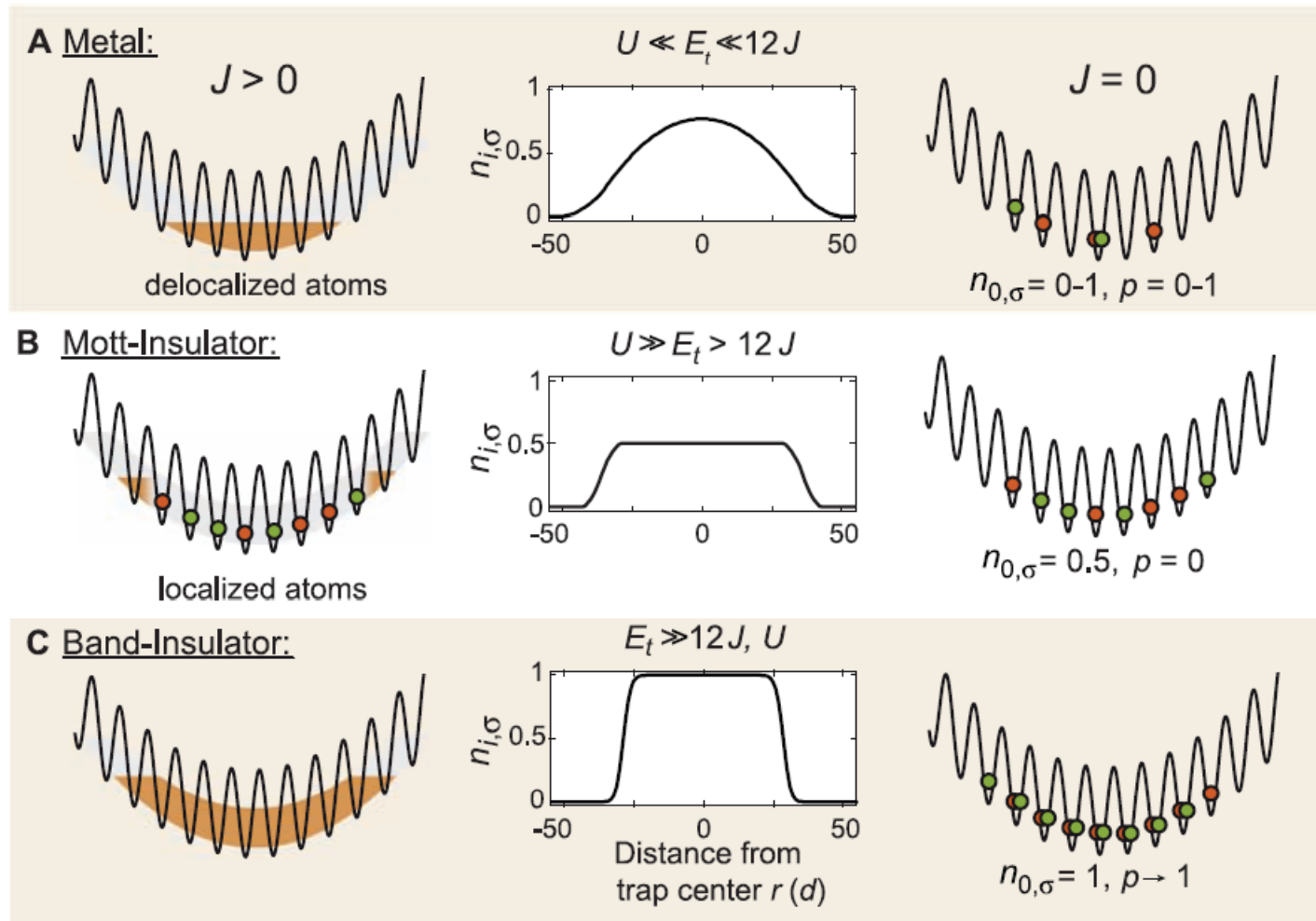
$$\begin{aligned} \text{Néel Temp: } \frac{4t^2}{U} &\cong (275\text{Hz}) h \\ &= 0.07 k_B T_F \\ &= 13 \text{ nK } k_B \end{aligned}$$

Superfluidity:

$$\begin{aligned} T_c &\sim 0.01t \sim 0.002k_B T_F \\ &= 350 \text{ pK } k_B \end{aligned}$$

State of the Art: $T \approx 0.25T_F$

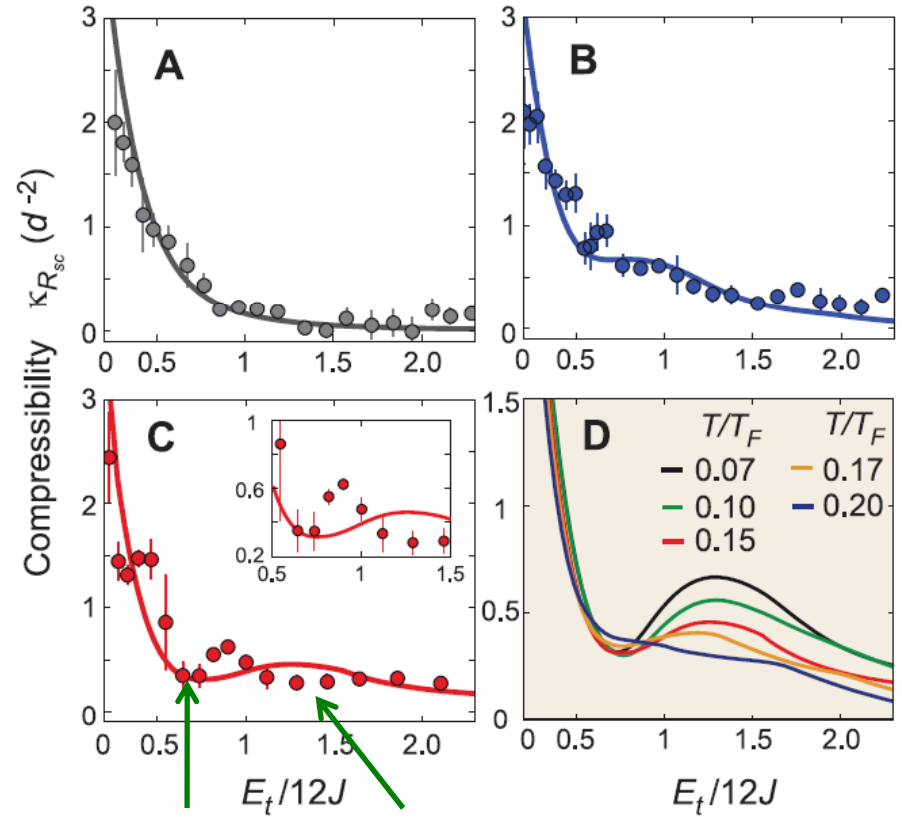
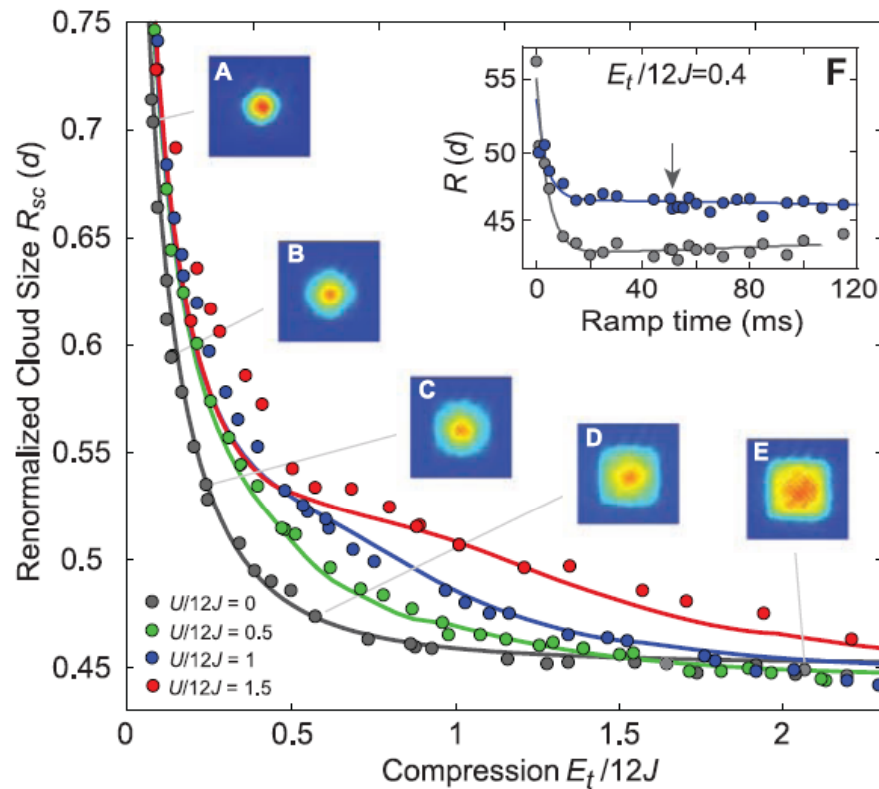
Observation of the Mott Insulator



T. Esslinger *et al.*, Science **322**, 1520 (2008)

I. Bloch *et al.*, Nature **455**, 204 (2008)

Observation of the Mott Insulator



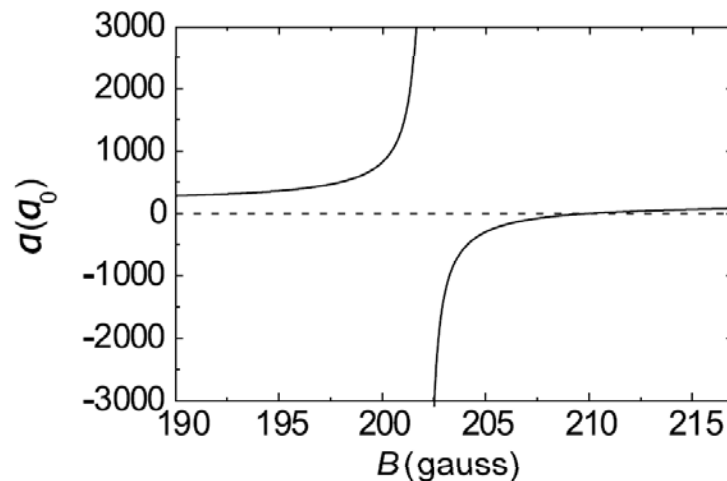
Mott Insulator Band Insulator

Requirement for Anti-Ferromagnet

Entropy of AF: $\frac{S}{N k_B} < \ln(2) = 0.69$

Entropy of Fermi Gas: $\frac{S}{N k_B} = \pi^2 \frac{T}{T_F}$

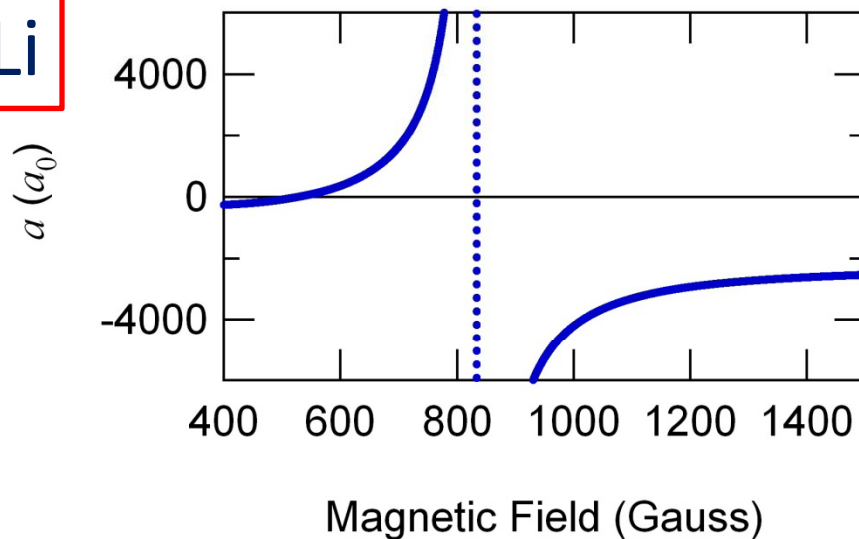
^{40}K



Large inelastic decay rates
3-body
p-wave resonance

$T/T_F \sim 0.15$ $\frac{S}{N k_B} = 1.48$

^6Li



Cooling near Feshbach Res.

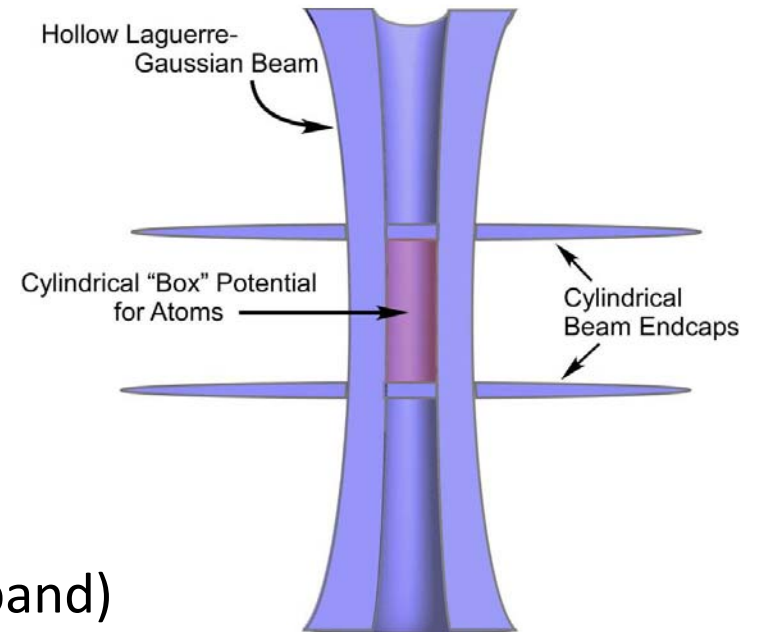
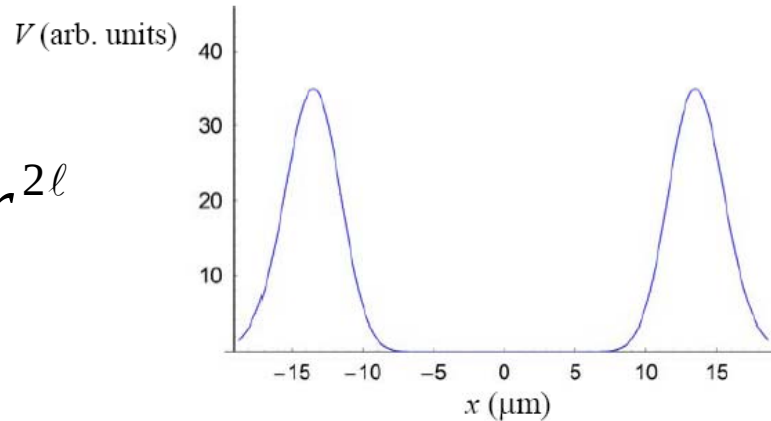
$T/T_F < 0.05$ $\frac{S}{N k_B} < 0.49$

Cannot access weak
repulsive interactions

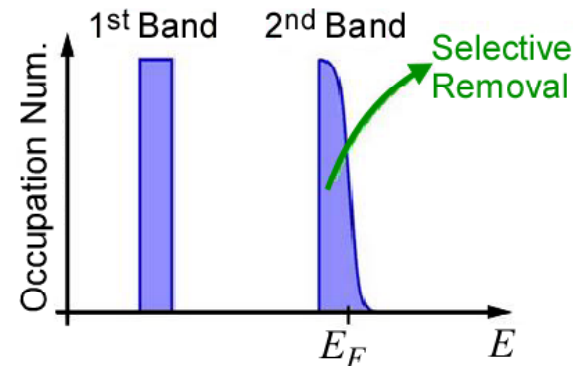
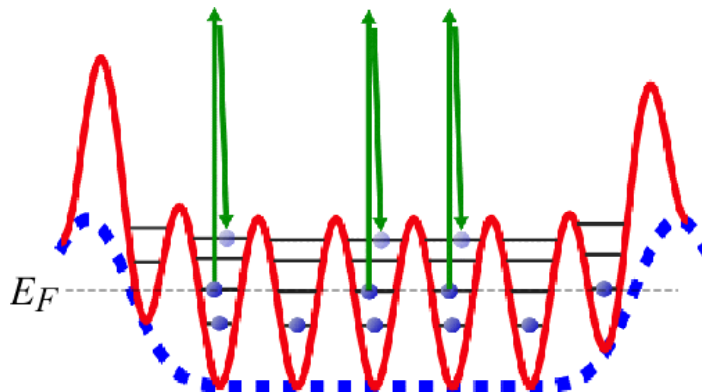
Removing Entropy from the System

- External confinement: box-like potential

$$V \propto r^{2\ell}$$



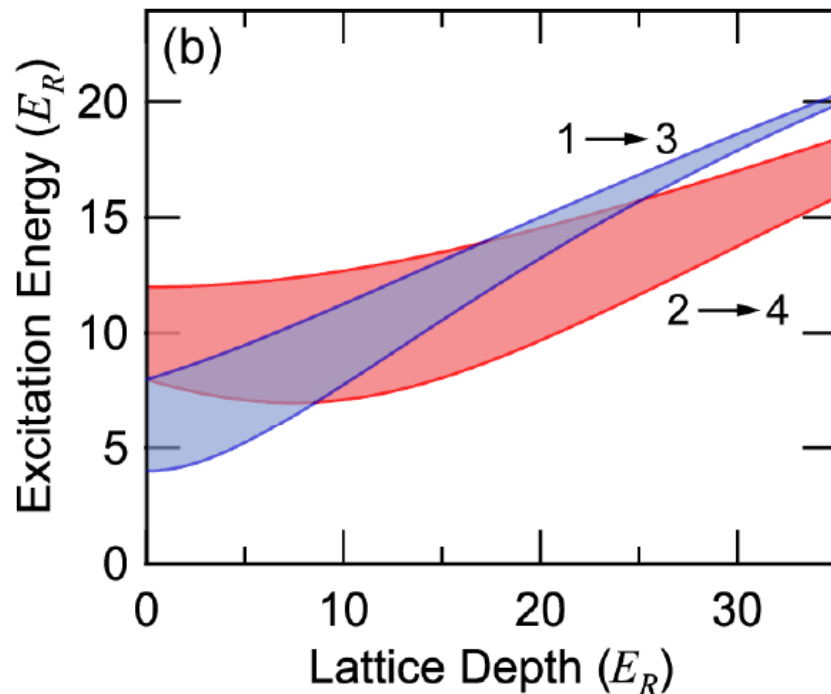
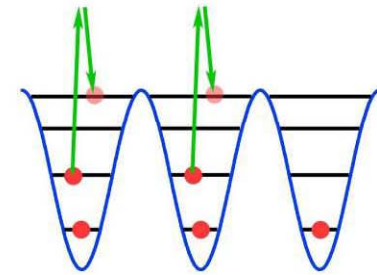
- Load atoms at high density (partially fill 2nd band)
- Selectively remove atoms from higher bands



Filtering higher bands in 1D w/ harmonic confinement J. I. Cirac *et al.*, New J. Phys. **8**, 164 (2006).

Selective Removal from 2nd Band

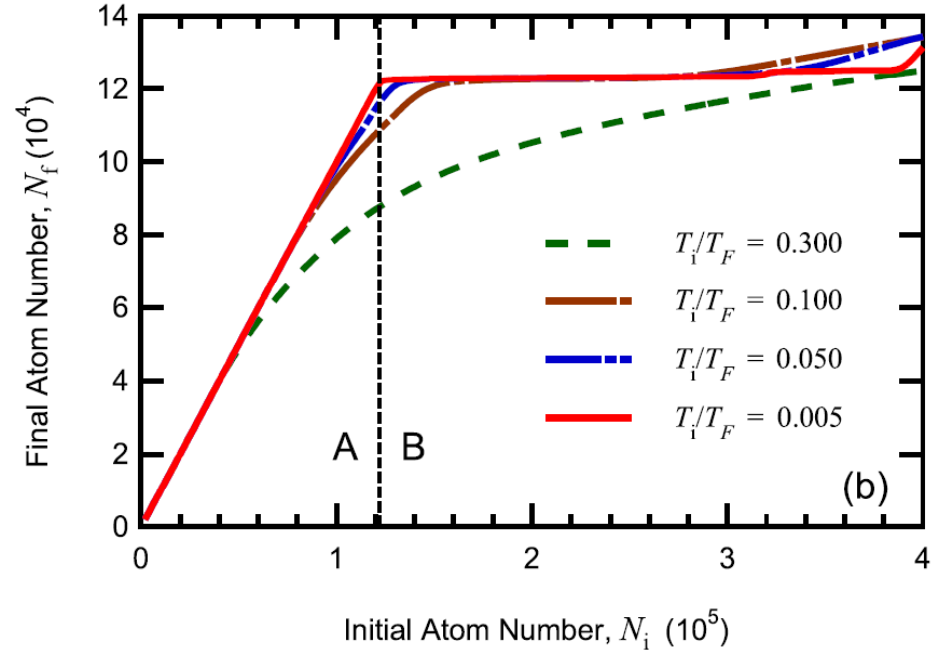
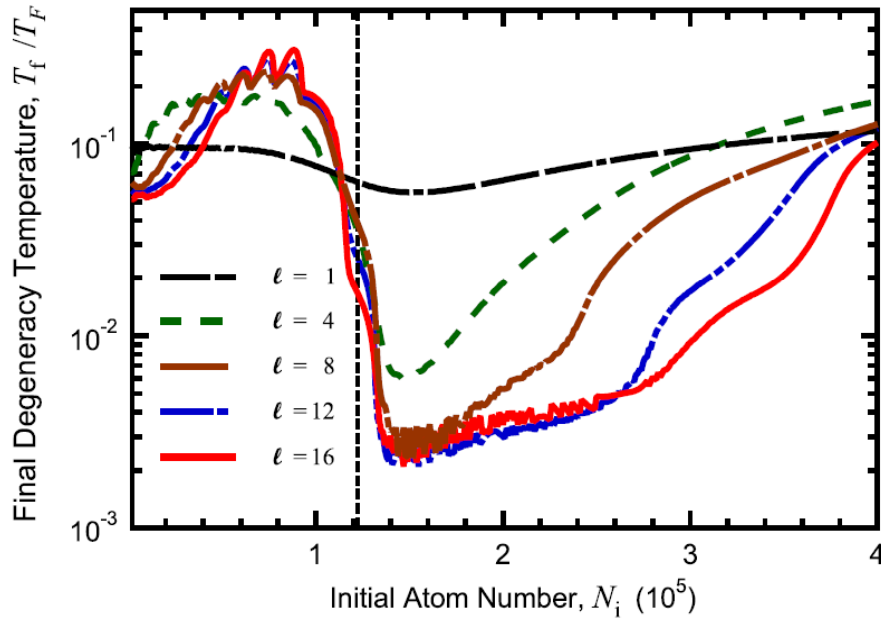
- Amplitude Modulation
 - Selectively Excite 2nd $\rightarrow$ 4th Band



- Band Excitation by Rapid Adiabatic Passage
- Irreversible Removal of Atoms
- Reduction in Entropy

Projected Final Temperatures

- Numerical Results for Loading $35 E_R$ Lattice and Filtering



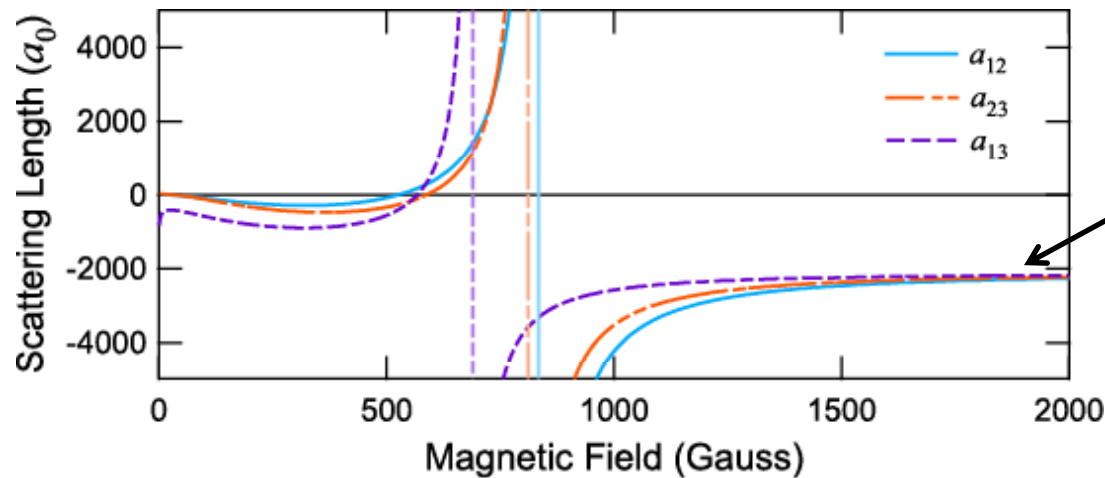
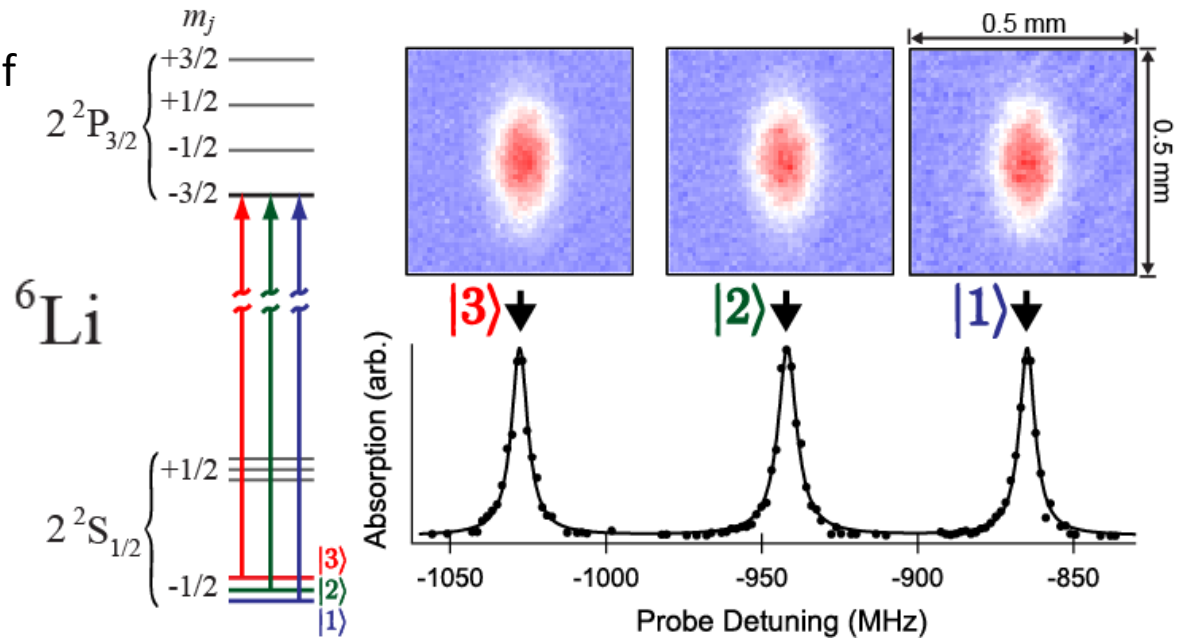
- Achieve $T = 0.003 T_F$ with 13% atom loss
- Initial $S/N = 0.28 k_B$, Final $S/N = 0.024 k_B$
- Adiabatically Lowering Lattice to $5 E_R$ gives $T < 0.001 T_F$

Outline

- o Ultracold Fermionic Atoms
- o The Hubbard Model
 - Quantum Magnetism
 - High-temperature superconductivity
- o Reducing the Entropy of Atoms in a Lattice
- o Three-Component Fermi Gases
 - SU(3) Hubbard model: Color Superfluidity
 - The quantum 3-body problem

A Three-Component Fermi Gas

- Three different internal states of ${}^6\text{Li}$ – labeled by color
- Mixture is stable against 2-body inelastic collisions
- Tunable interactions

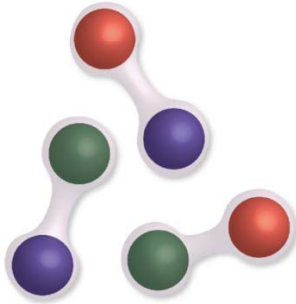


- Equal & strongly attractive interactions at high field ($a_t = -2140 a_0$)
- BCS Pairing with 3 types of Cooper pairing

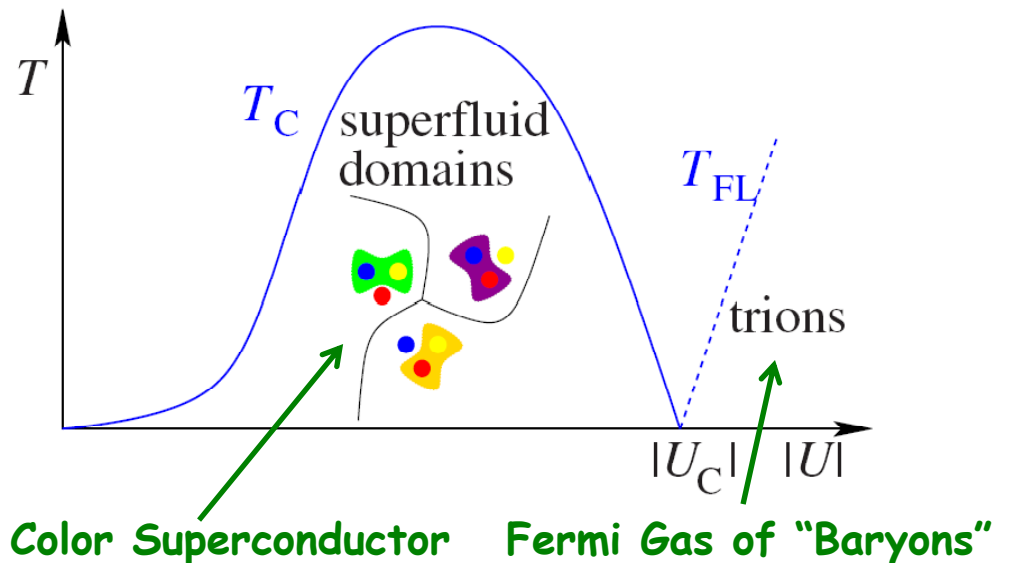
Huckans, Williams, Hazlett, Stites & O'Hara, PRL **102**, 163502 (2009)

“Color” Superconductivity

- BCS superfluid with three types of Cooper pairs



- Optical lattice: controls ratio of interaction energy to kinetic energy
- Quantum phase transition from “color superfluid” to “baryons” (RGB bound states) as interaction energy is increased



QUANTUM CHROMODYNAMICS

Lifestyles of the small and simple

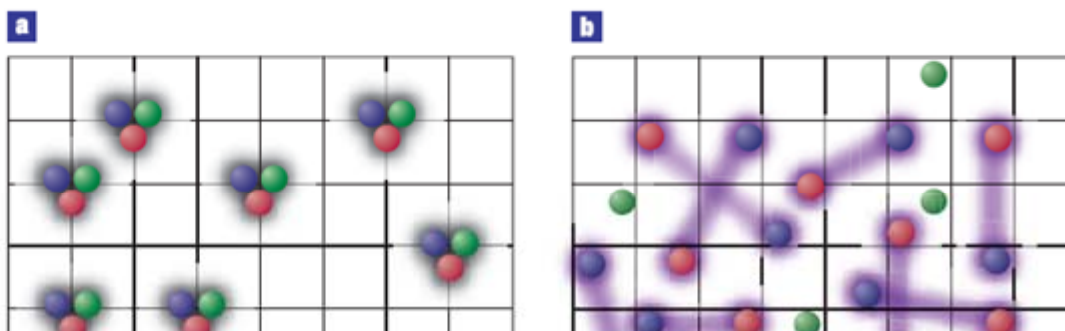
Ultracold atoms in optical lattices are already used to simulate complex solid-state phenomena. But could the same platform also give us a better grasp of how quarks group together?

Frank Wilczek

Is in the Department of Physics, Massachusetts Institute of Technology, 5 Cambridge Center, Cambridge, Massachusetts 02139, USA.

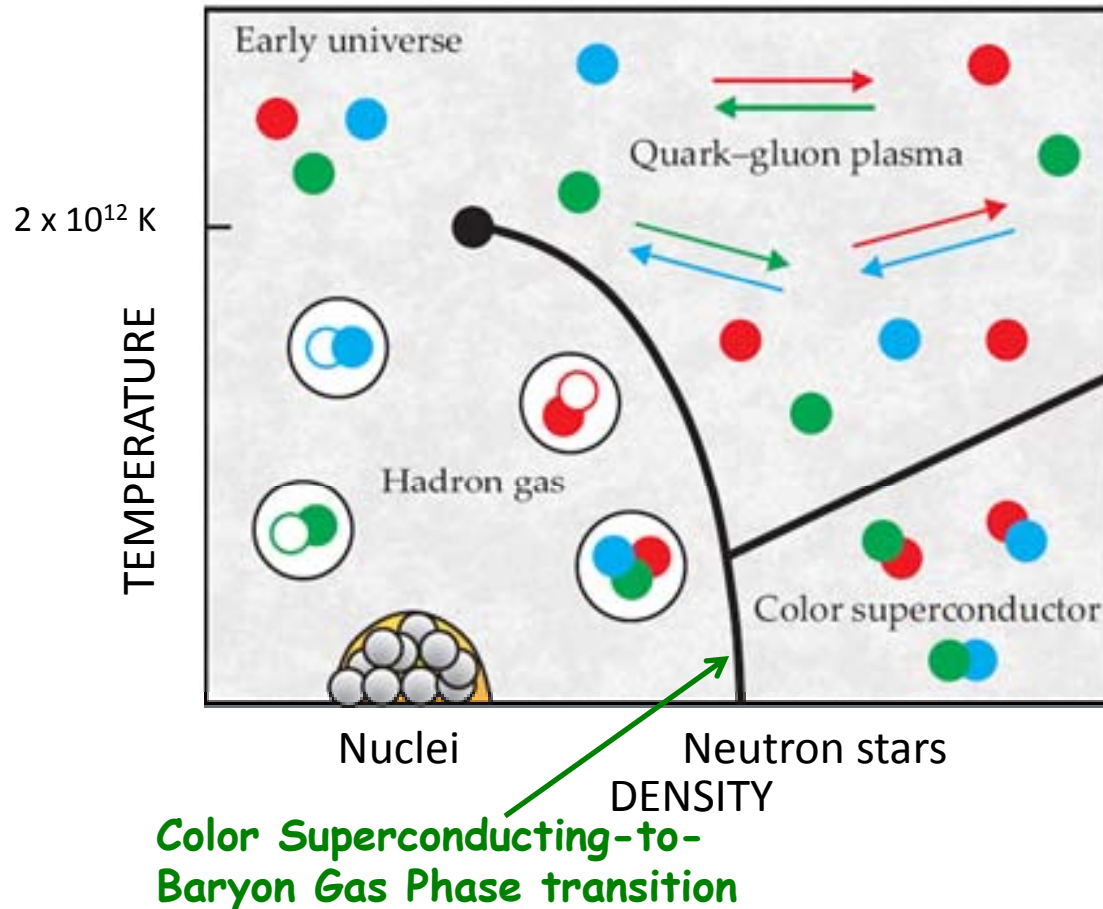
e-mail: wilczek@mit.edu

Quarks find other quarks — with different flavours and colours — attractive. They can choose either



QCD Phase Diagram

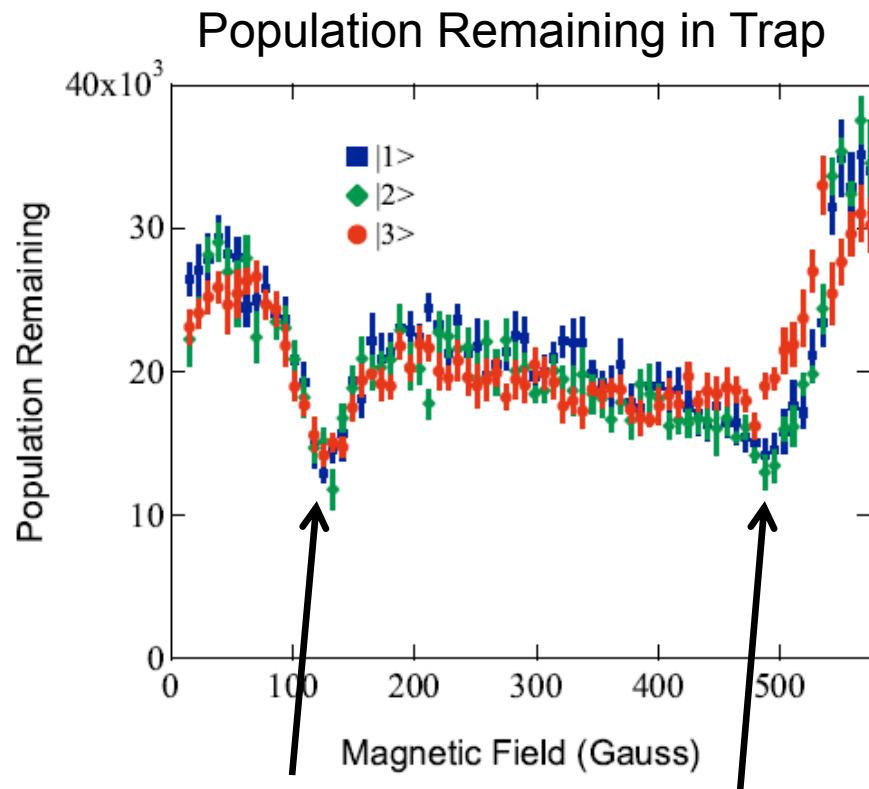
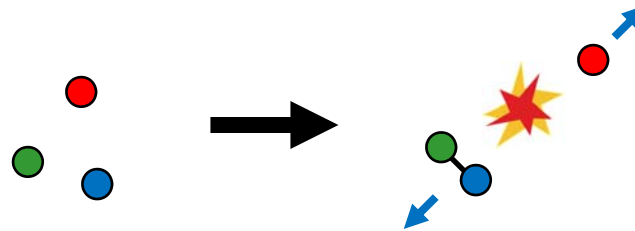
Conjectured phase diagram for QCD



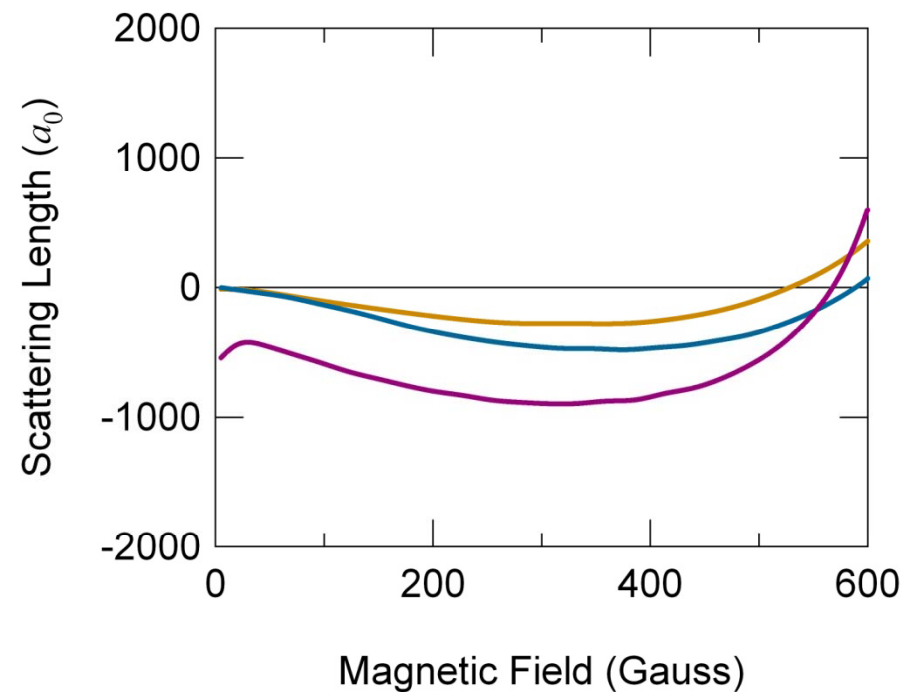
- Ultracold atoms can be used to simulate part of the QCD phase diagram
- Is the baryon-to-CSF phase transition first or second order?
- Does the color superfluid exhibit domain formation?
- What phases are possible in a spin-imbalanced 3-state Fermi gas?
(which mimics the higher mass of the strange quark)

Three-Body Recombination

Three-body recombination:



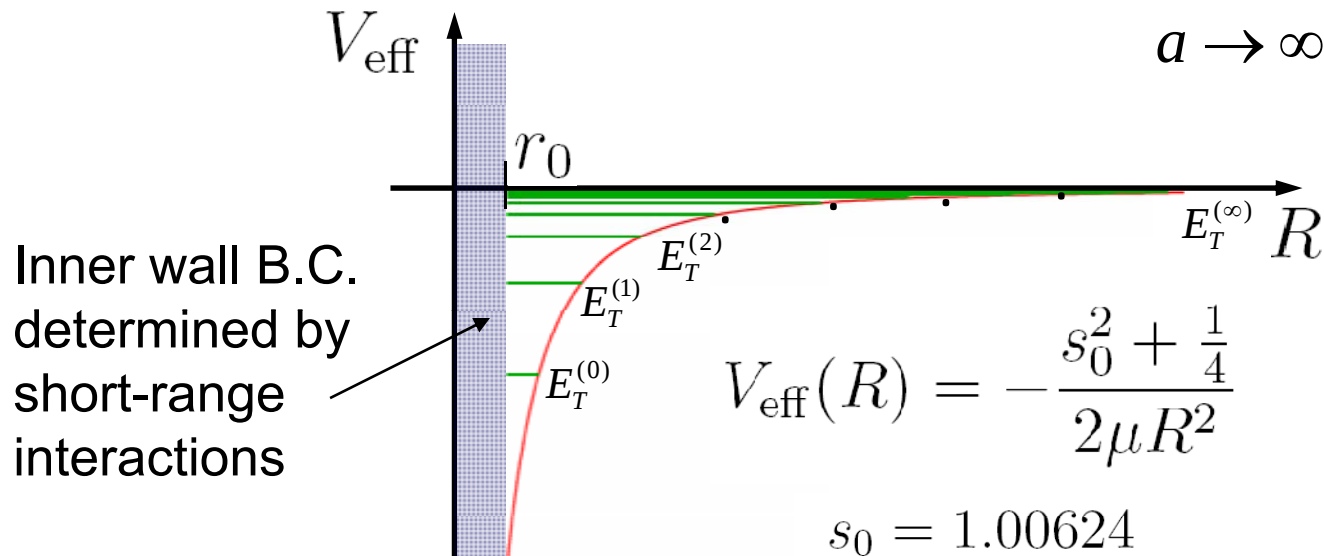
Low Field Scattering Lengths



Resonances in the 3-Body Recombination Rate!

3-Body Problem: Efimov's Solution

(1970) Efimov: An infinite number of bound 3-body states



Vitaly Efimov
circa 1970

Infinitely many 3-body bound states (universal scaling):

$$E_T^{(n+1)} / E_T^{(n)} \rightarrow 1/e^{2\pi/s_0} = 1/515.03$$

A single 3-body parameter: κ_*

$$E_T^{(n)} \rightarrow \left(\frac{1}{515.03} \right)^{n-n_*} \frac{\hbar^2 \kappa_*^2}{m}$$

3-Body QM Problem: Efimov's Solution

(1970 & 1971) Efimov: Identical Bosons in Universal Regime $a \gg r_0$

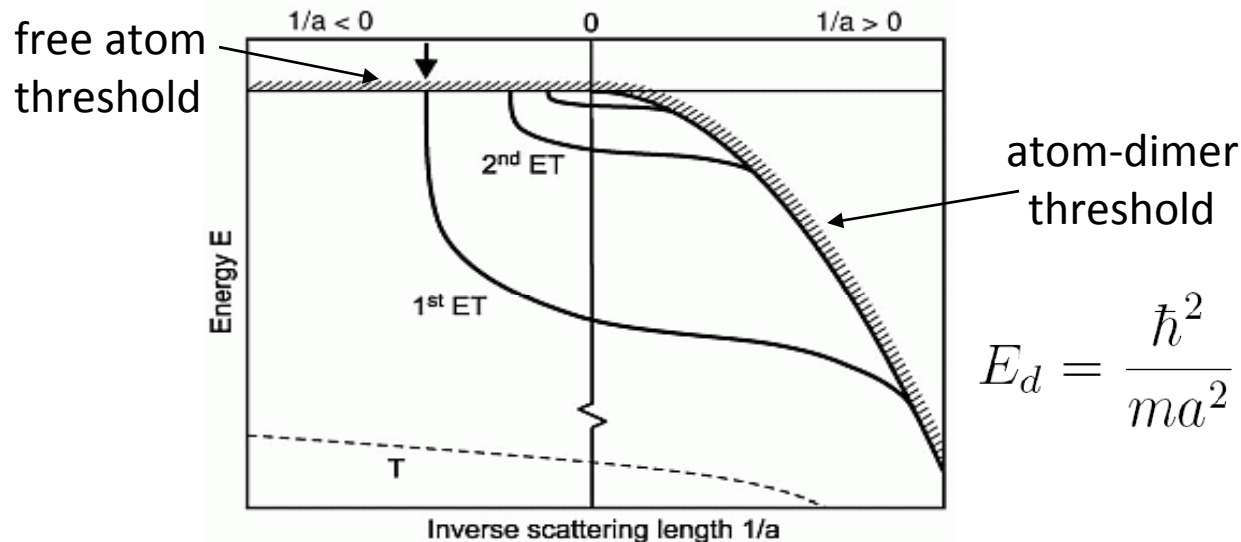


Diagram from:
T. Kraemer *et al.*,
Nature **440** 315 (2006)

Note:

Only two free parameters:

κ^* and η^*

Observable:

Enhanced 3-body recombination ↓

$$K_3 = \frac{4677 \sinh(2\eta_*)}{\sin^2[s_0 \ln(0.6642|a|\kappa_*)] + \sinh^2 \eta_*} \frac{\hbar a^4}{m} \quad (a < 0)$$

Log-periodic scaling

$$a < 0 \quad \& \quad a > 0$$

$$a_*^{(n+1)} / a_*^{(n)} \rightarrow 22.7$$

E. Braaten, H.-W. Hammer, D. Kang and L. Platter, arXiv:0811.3578

3-Body QM Problem: Efimov's Solution

(1970 & 1971) Efimov: Identical Bosons in Universal Regime

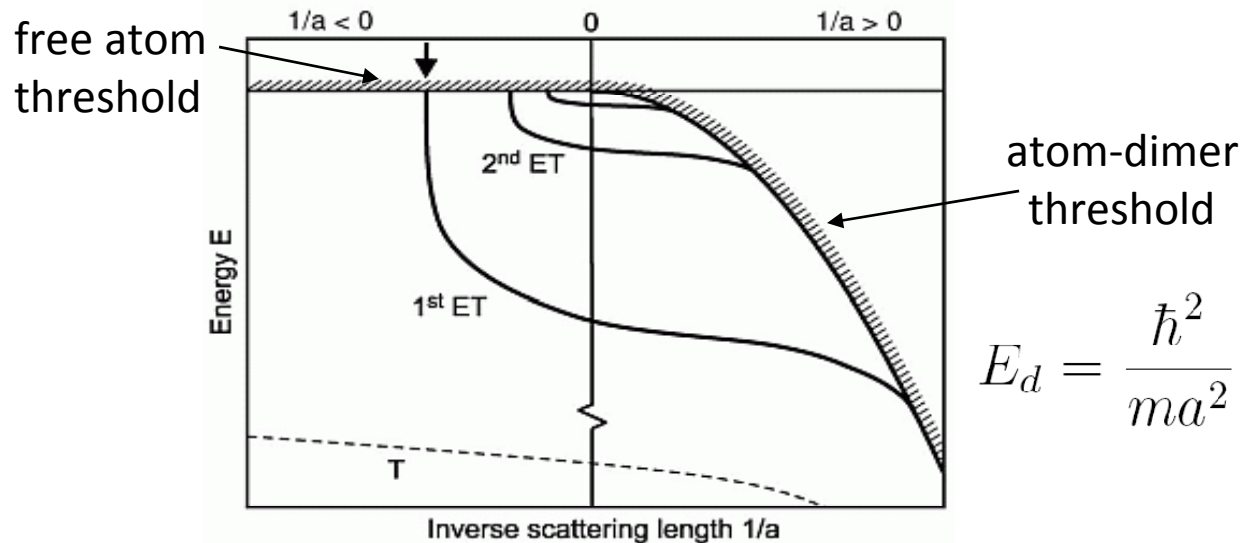


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E. Braaten, H.-W. Hammer, D. Kang and L. Platter, arXiv:0811.3578

Observed Efimov Features: (ultracold Bosons)

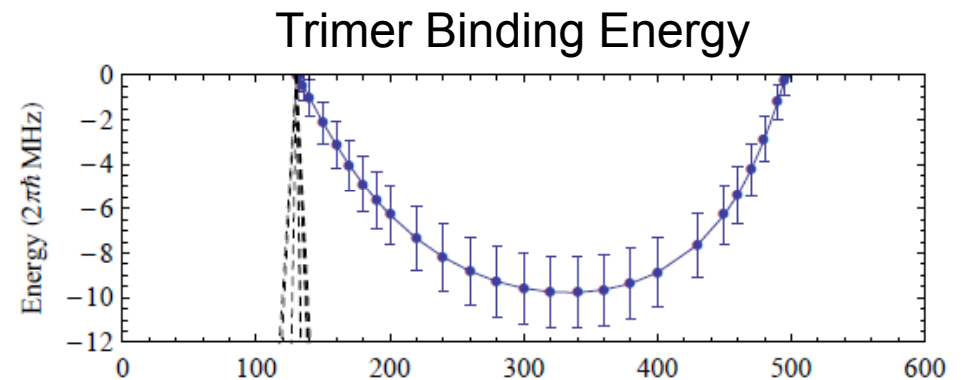
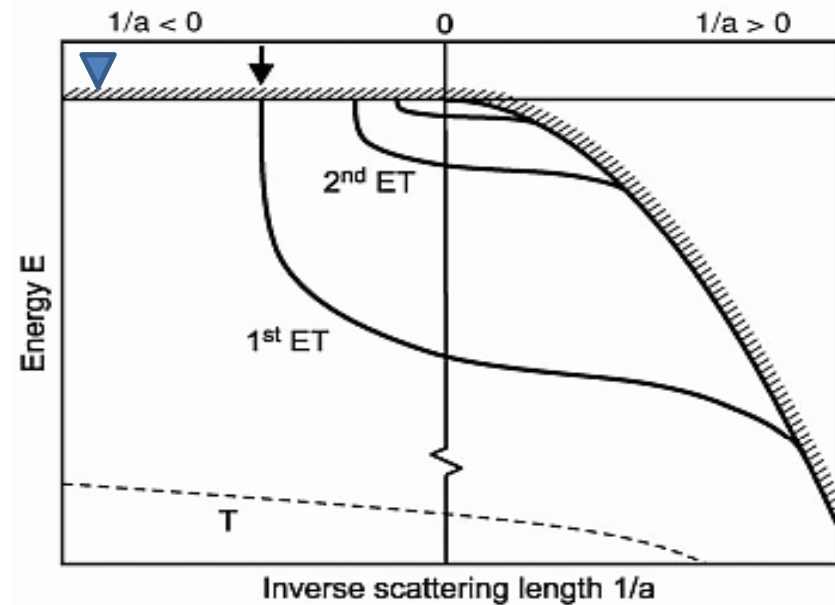
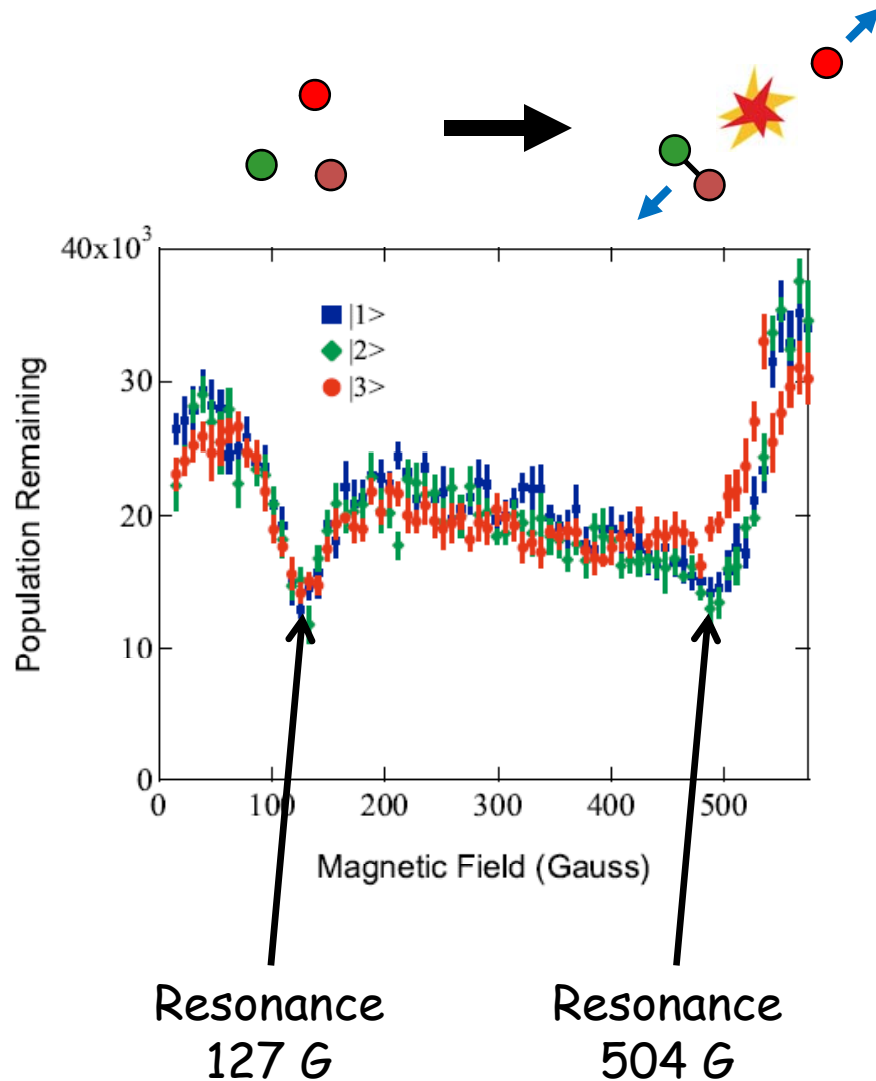
T. Kraemer *et al.* Nature **440** 315 (2006) [Cs]

S. Knoop *et al.* Nature Physics **5** 227 (2009) [Cs]

G. Barontini *et al.* arXiv:0901.4584 [K, Rb]

M. Zaccanti *et al.* arXiv:0904.4453 [K]

Resonant Loss Features

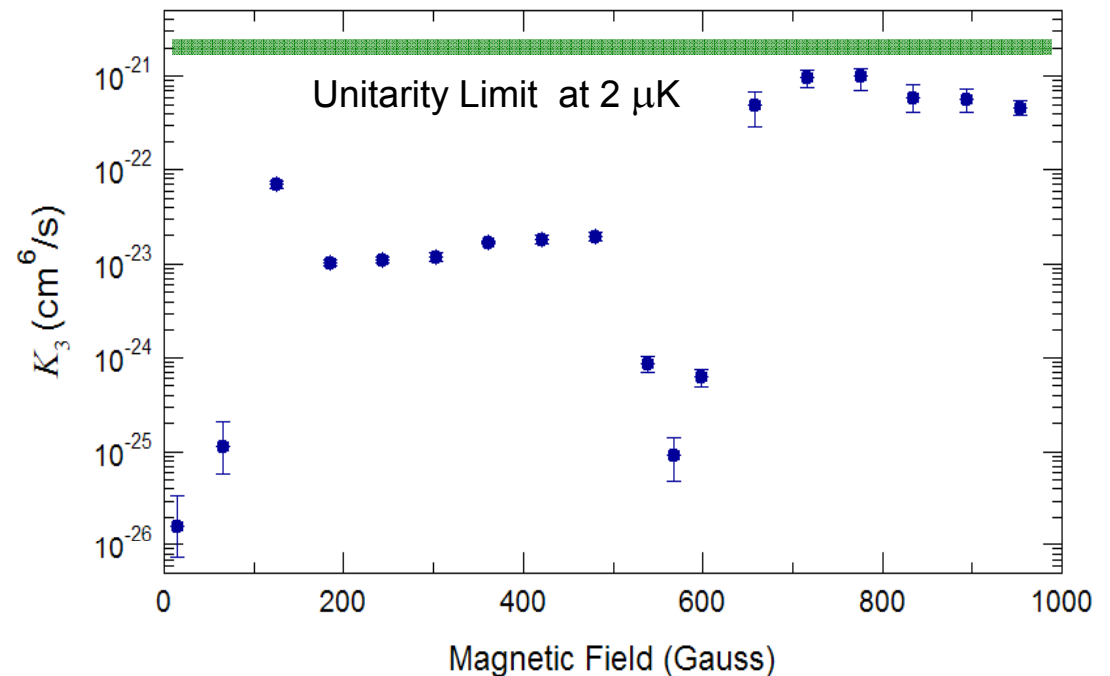
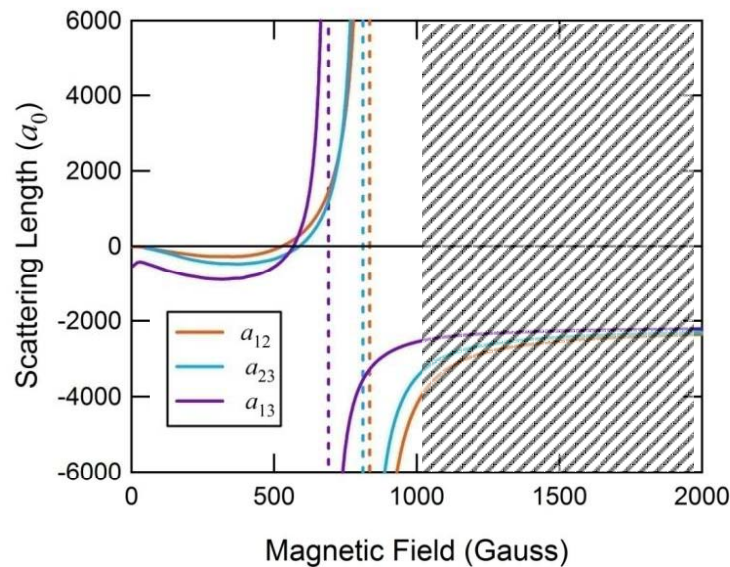


P. Naidon & M. Ueda, arXiv:0811.4086

Resonances in the 3-Body Recombination Rate!

Low-Field 3-Body Loss Data

Low Field Scattering Lengths



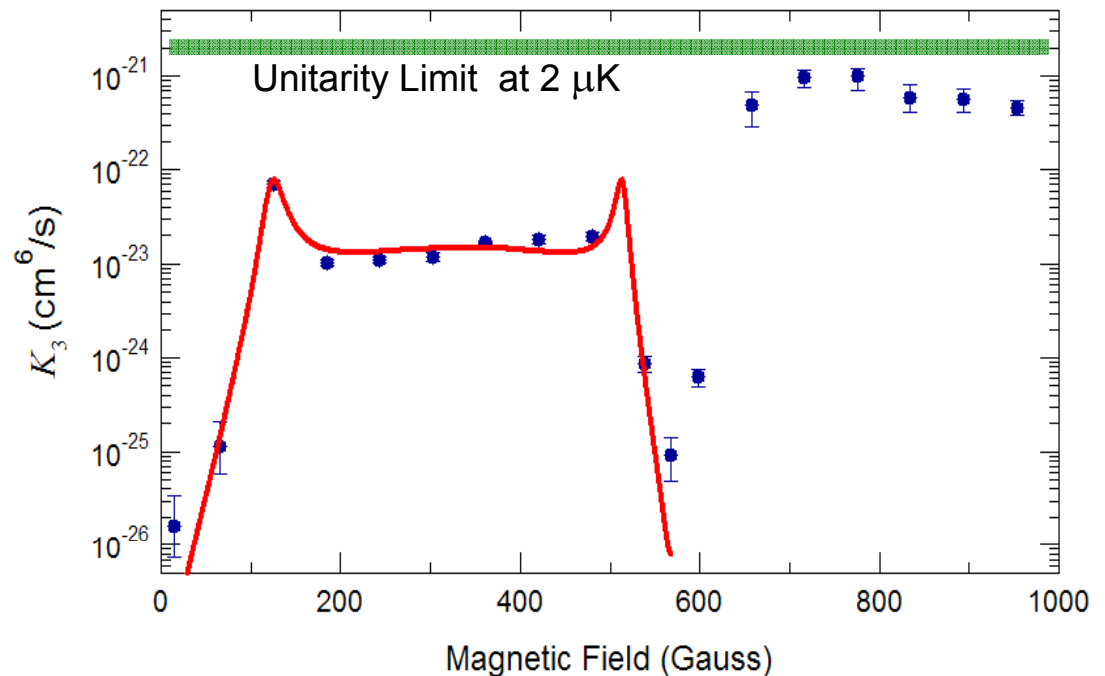
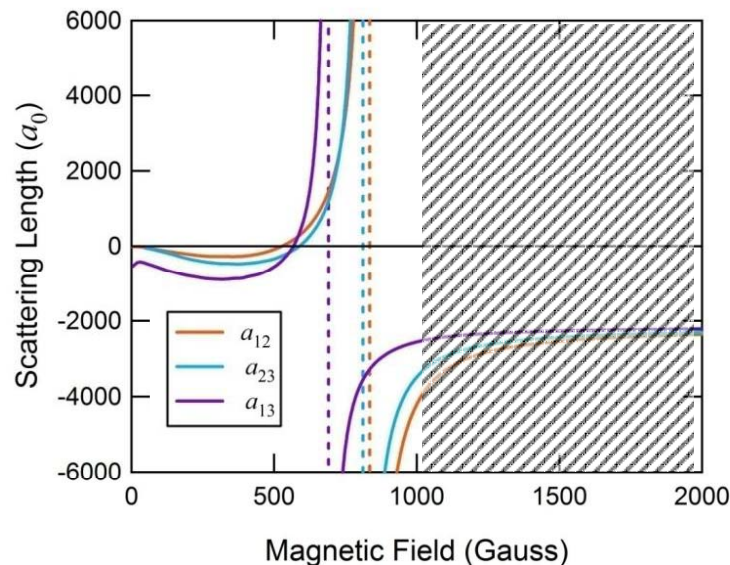
Low field 3-Body loss rate (K_3) over $10\text{ G} < B < 960\text{ G}$

Huckans *et al.*, PRL **102**, 165302 (2009) (Penn State)

Ottenstein *et al.*, PRL **101**, 203202 (2008) (Heidelberg)

Low-Field 3-Body Loss Data

Low Field Scattering Lengths



Fit with 2 free parameters:

κ^* and η^*

$$K_3 = \frac{4677 \sinh(2\eta_*)}{\sin^2[s_0 \ln(0.6642|a|\kappa_*)] + \sinh^2 \eta_*} \frac{\hbar a^4}{m}$$

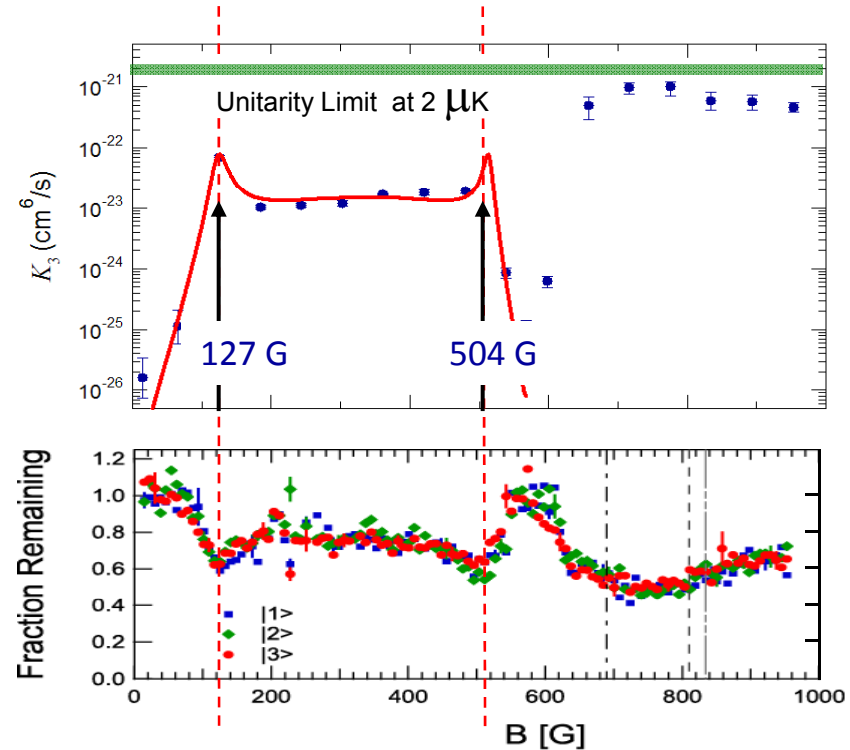
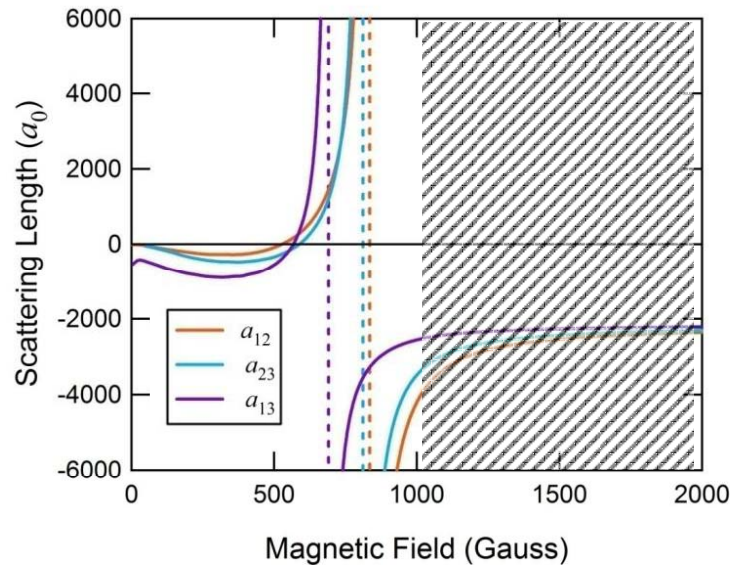
E. Braaten, *et al.*, arXiv:0811.3578

$$a_{\text{eff}} = \sqrt[4]{\frac{1}{3} (a_{12}^2 a_{23}^2 + a_{23}^2 a_{13}^2 + a_{12}^2 a_{13}^2)}$$

Andre Wenz, Diploma Thesis (2009)

Low-Field 3-Body Loss Data

Low Field Scattering Lengths

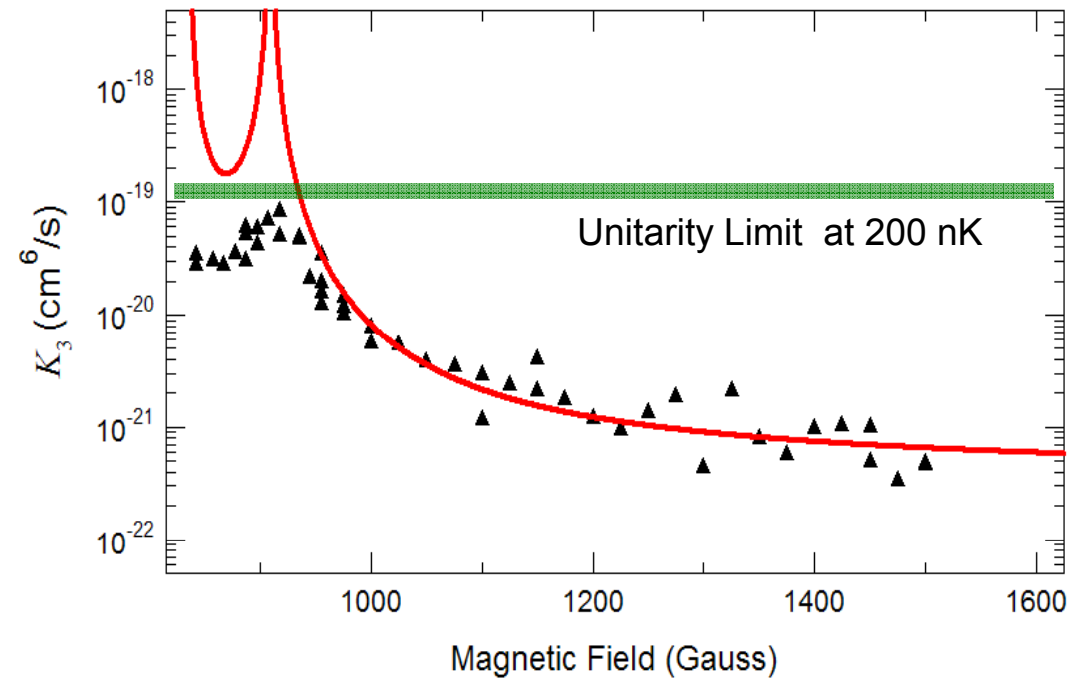
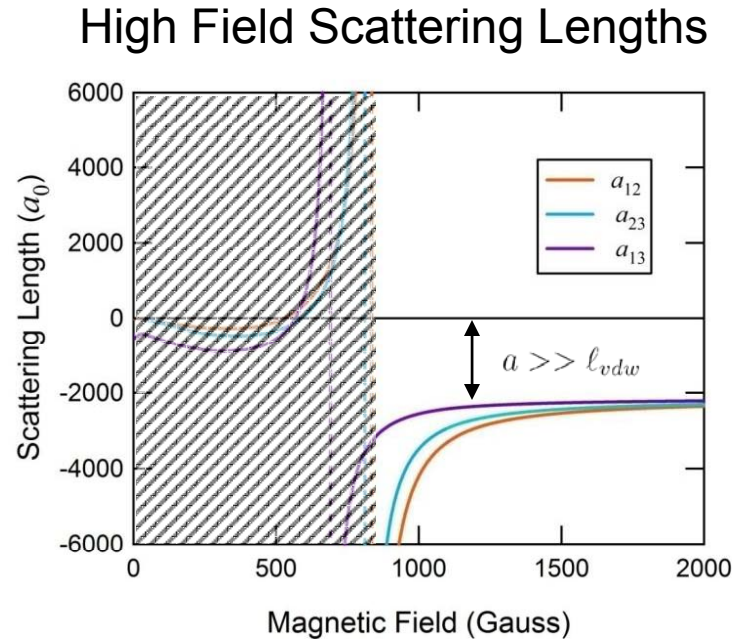


Resonant loss features at 127 and 504 Gauss

$$a_{\text{eff}}(127 \text{ G}) \simeq a_{\text{eff}}(504 \text{ G})$$

$$K_* = 59.9 \text{ } a_0^{-1}, \quad \eta_* = .075$$

High-Field 3-Body Loss Data



High field 3-Body loss rate (K_3) over $840 \text{ G} < B < 1500 \text{ G}$

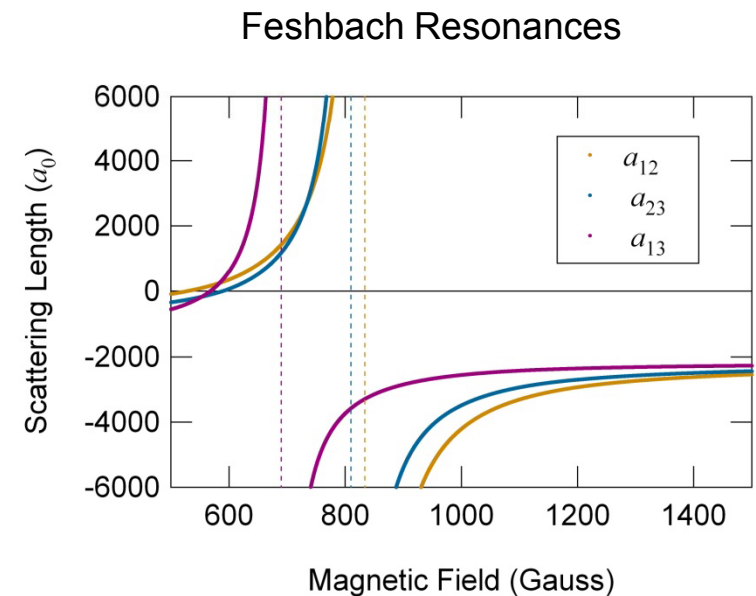
Resonant loss feature at 910 Gauss $a_{ij} > 35 r_0$

$$K_* = 79.2 a_0^{-1}, \quad \eta_* = .01125$$

Future Prospects

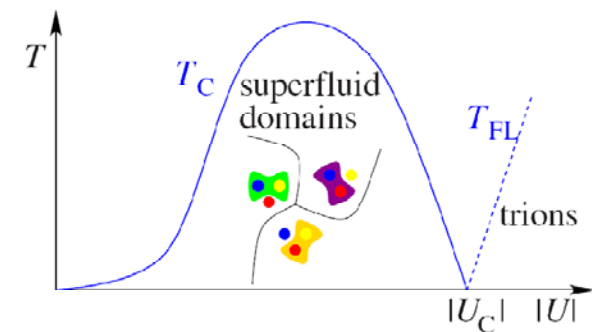
Efimov Physics with Overlapping Feshbach Resonances

- Multiple scattering lengths: (a_{12} , a_{23} , a_{13})
E. Braaten, *et al.*, arXiv:0811.3578 (2008)
J. P. D’Incao and B. D. Esry, arXiv:0905.0772 (2009)
- Mutual attractive and repulsive interactions:
($a_{12}, a_{23} > 0$ & $a_{13} < 0$)



BCS Pairing in a 3-State Fermi Gas

- Pairing competition (attractive interactions)
- Unequal Fermi Surfaces (E_F)
- Superfluid Domain Formation
- Superfluid/Trion phase transition (3D Lattice)
- No condensed matter analog

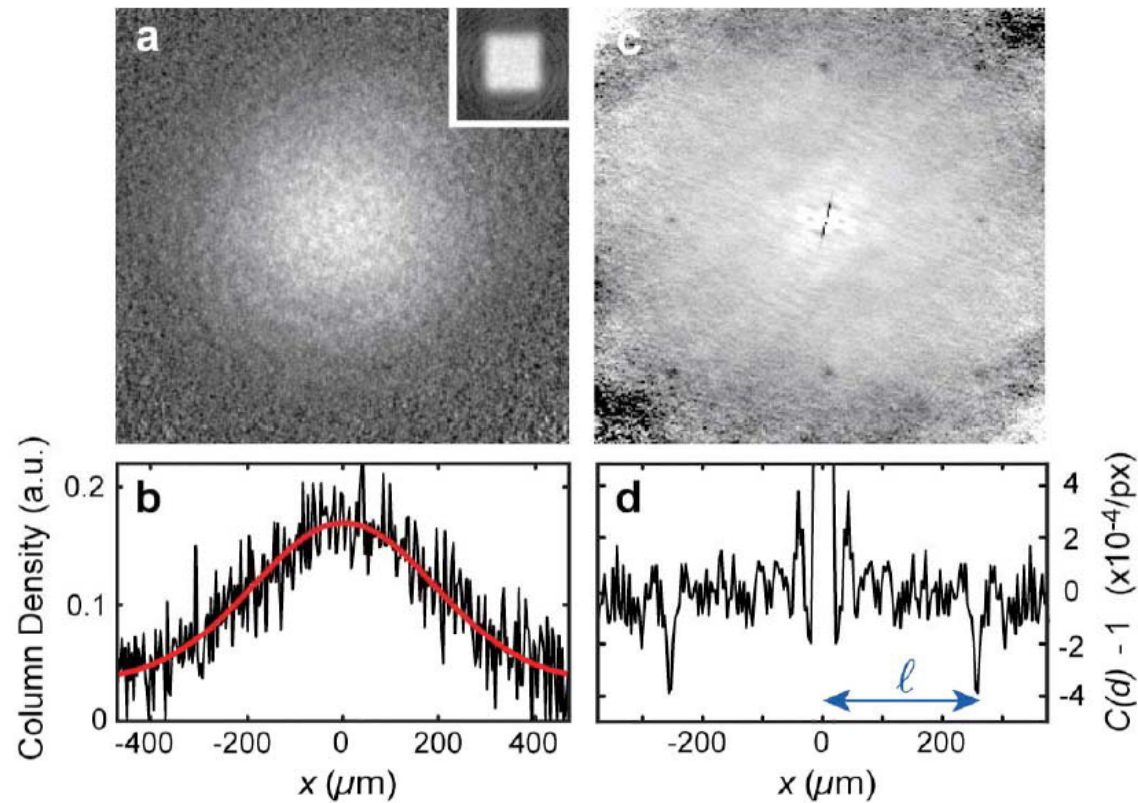


Rapp, Hofstetter & Zaránd,
PRB 77, 144520 (2008)

Summary

- Hubbard model can be realized with ultracold atoms
- Mott insulator state has been observed in Zurich and Mainz
- Observation of AF will require reducing the entropy
- Entropy can be significantly reduced by filtering high entropy atoms
- Three-component Fermi gases in a lattice: Quant. Sim. of QCD
- Observation of Efimov resonances with unequal scattering lengths
 - Two low-field resonances with moderate scattering lengths (127 & 504 G)
 - Third high-field resonance at 910 G (well into universal regime)
- Agreement with Efimov theory $|a| \gg r_0$ and $|a| \sim r_0$
 - Fits yield 3-body parameters for ${}^6\text{Li}$ at negative scattering lengths

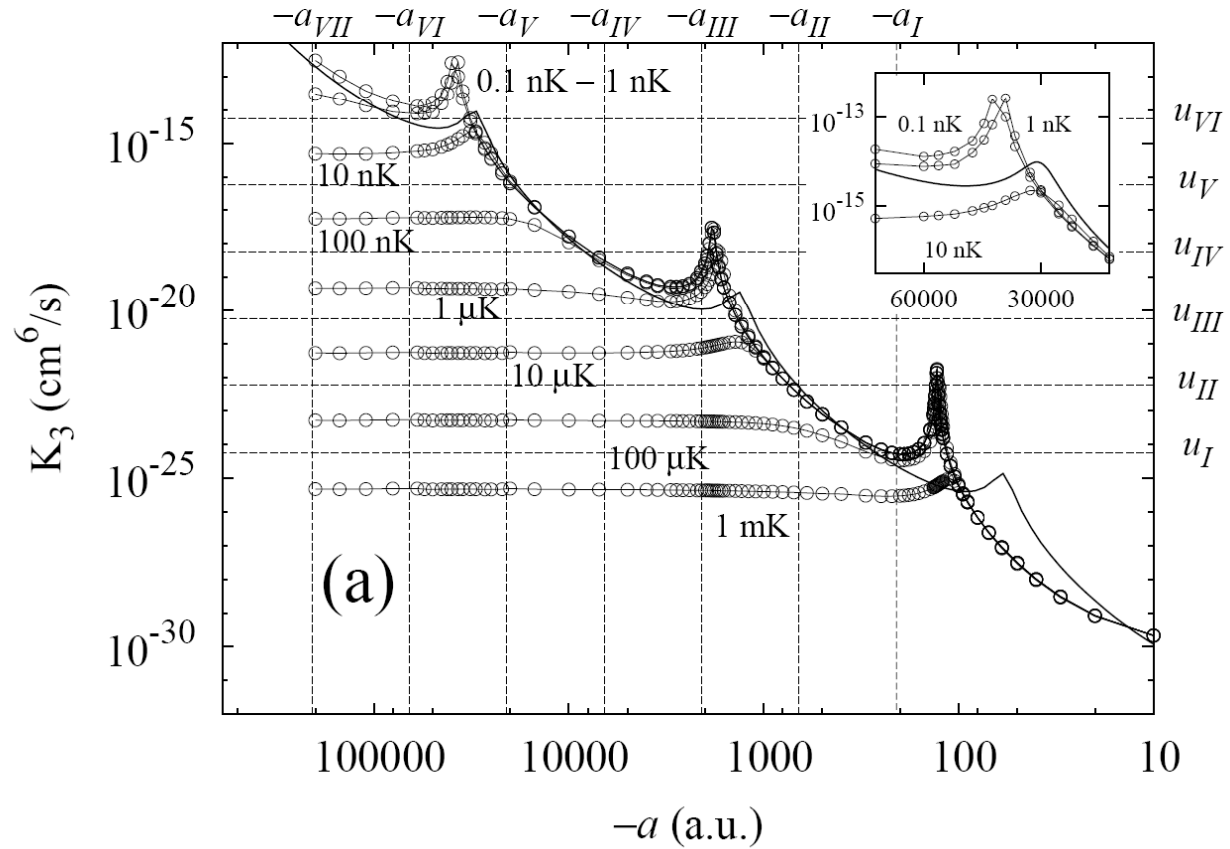
Density-Density Correlations



Nature, Vol 444, 733-736 (2006)

Density-density correlation function

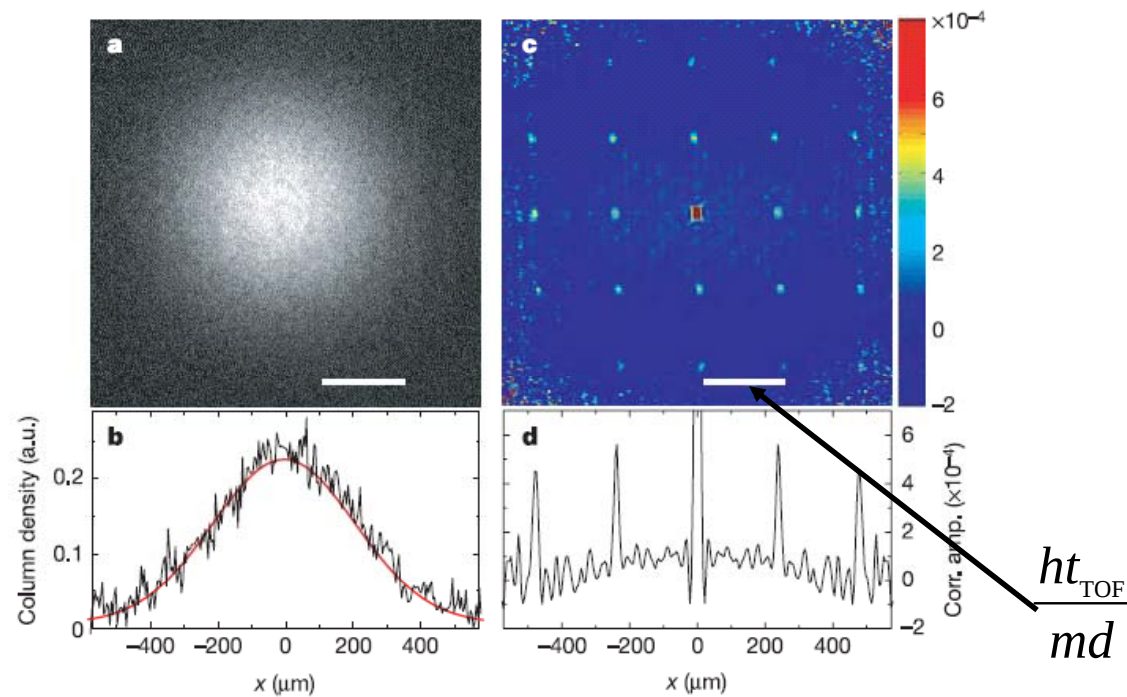
Limits on Universality



J. P. D’Incao, H. Suno & B. Esry, Phys. Rev. Lett. **93**, 123201 (2004)

Density-Density Correlations

- TOF Expansion and Image
- Density-Density Correlations in Expanded Cloud



(I. Bloch, 2005)

Density-Density Correlations

- State Selective Imaging of TOF Expansion
- Density-Density Correlations in Expanded Cloud

