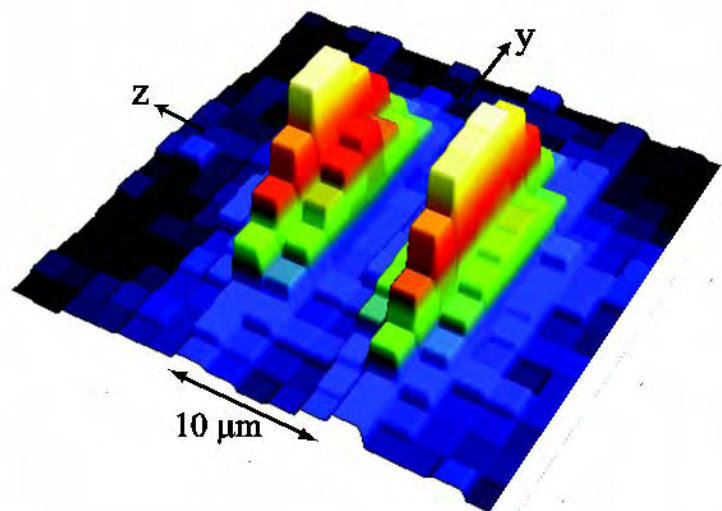


Demonstration of a neutral atom CNOT gate and perspectives for multi-qubit operations

Mark Saffman



two trapped atoms

Neutral atom quantum gates

Quantum gates via Rydberg blockade

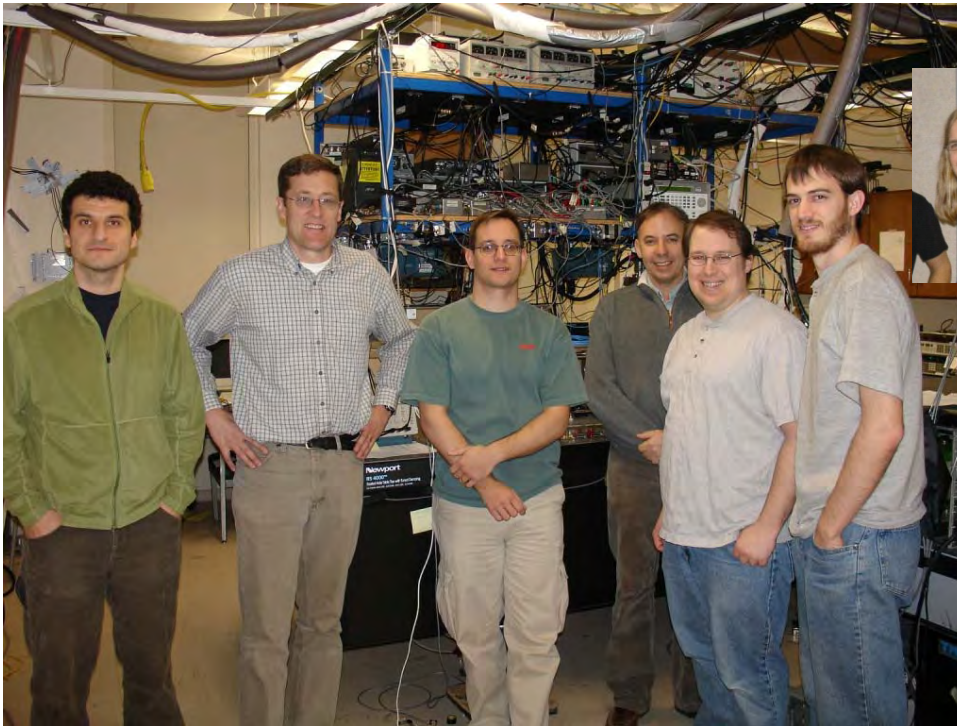
Scalability

Preparation of single atom states

Collective encoding



People



Deniz
Yavuz

Thad
Walker

Erich
Urban

MS

Thomas
Henage

Larry
Isenhower



Todd
Johnson -> Boulder

Alumni

Marie Delaney
Pasad Kulatunga
Deniz Yavuz
Todd Johnson

Precision Photonics, Boulder
Hobart & William Smith Colleges, faculty
UW Madison, faculty
NIST, Boulder

Collaborations

Kelvin Wagner
Jungsang Kim

UC Boulder – qubit addressing optics
Duke – qubit addressing optics

Klaus Mølmer

Aarhus - theory

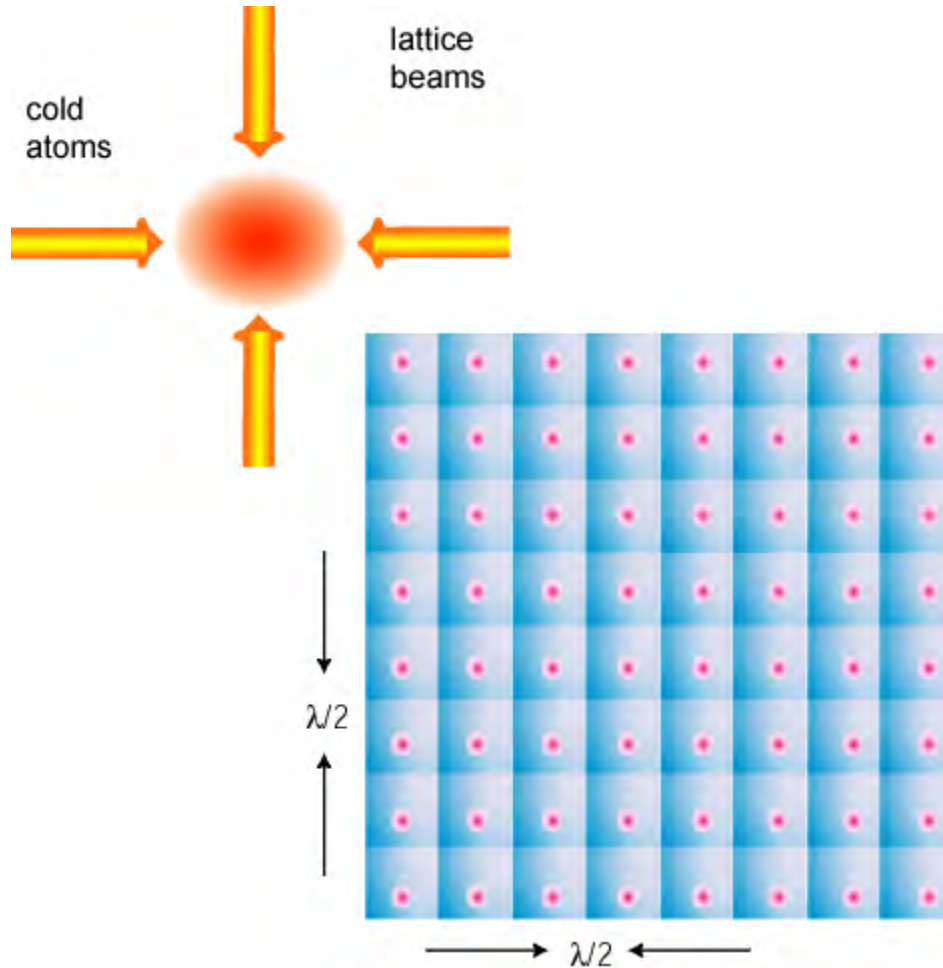
Also

Alex Gill (PhD)

Jon Sedlacek, Jake Covey, Adam Beardsley (undergrads)



Neutral atoms in optical lattices



Say $\lambda = 1 \mu\text{m}$.

4 million lattice sites per mm^2

8 billion lattice sites per mm^3

Critical issues:

- loading
- addressability
- interconnects

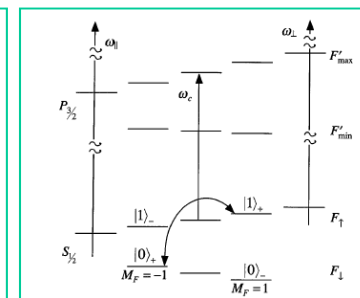
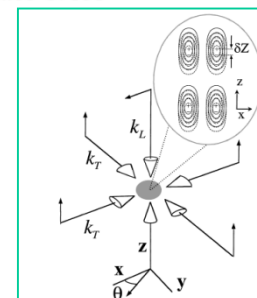
Quantum Logic Gates in Optical Lattices

Gavin K. Brennen,¹ Carlton M. Caves,¹ Poul S. Jessen,² and Ivan H. Deutsch¹

¹*Center for Advanced Studies, Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico 87131*

²*Optical Sciences Center, University of Arizona, Tucson, Arizona 85721*

(Received 29 May 1998)



Entanglement of Atoms via Cold Controlled Collisions

D. Jaksch,¹ H.-J. Briegel,¹ J. I. Cirac,¹ C. W. Gardiner,² and P. Zoller¹

¹*Institut für Theoretische Physik, Universität Innsbruck, Technikerstrasse 25, A-6020 Innsbruck, Austria*

²*School of Chemical and Physical Sciences, Victoria University, Wellington, New Zealand*

(Received 28 October 1998)

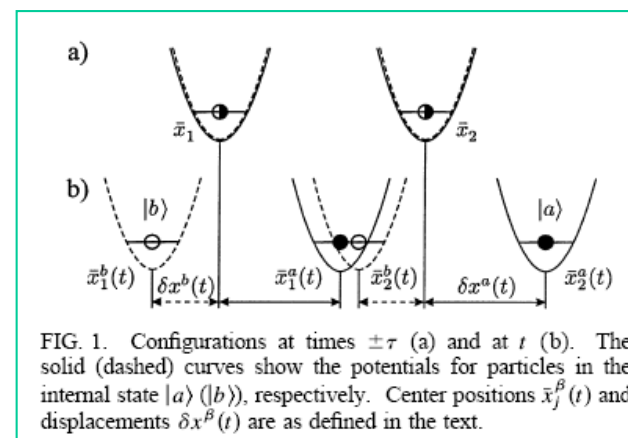


FIG. 1. Configurations at times $\pm\tau$ (a) and at t (b). The solid (dashed) curves show the potentials for particles in the internal state $|a\rangle$ ($|b\rangle$), respectively. Center positions $\bar{x}_j^\beta(t)$ and displacements $\delta x^\beta(t)$ are as defined in the text.

Many body collisional entanglement

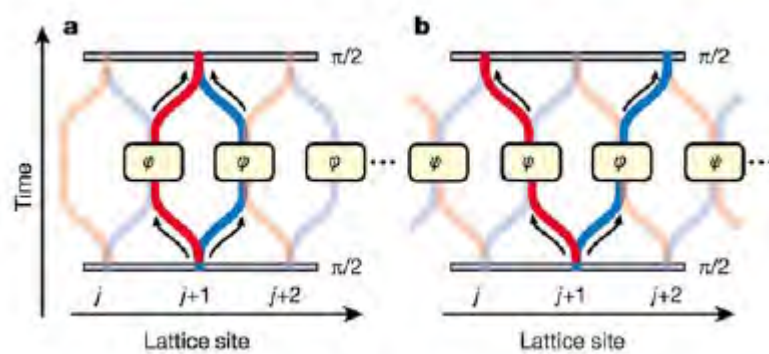
Signatures of entanglement have been demonstrated in lattices, but not at the two-atom level:

NATURE | VOL 425 | 30 OCTOBER 2003 **letters to nature**

.....

Controlled collisions for multi-particle entanglement of optically trapped atoms

Olaf Mandel, Markus Greiner, Artur Widera, Tim Rom, Theodor W. Hänsch & Immanuel Bloch

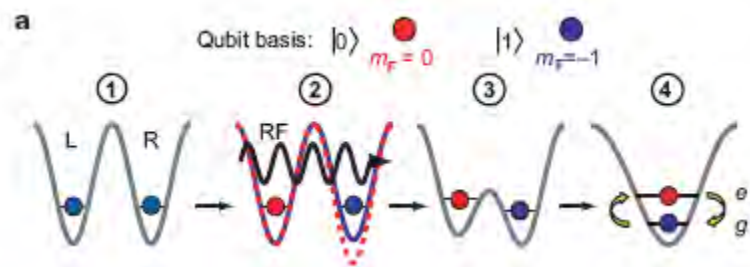


nature Vol 448 | 26 July 2007 | doi:10.1038/nature06011

LETTERS

Controlled exchange interaction between pairs of neutral atoms in an optical lattice

Marco Anderlini¹†, Patricia J. Lee¹, Benjamin L. Brown¹, Jennifer Sebby-Strabley¹†, William D. Phillips¹ & J. V. Porto¹



Fast Quantum Gates for Neutral Atoms

D. Jaksch, J.I. Cirac, and P. Zoller

Institut für Theoretische Physik, Universität Innsbruck, Technikerstrasse 25, A-6020 Innsbruck, Austria

S.L. Rolston

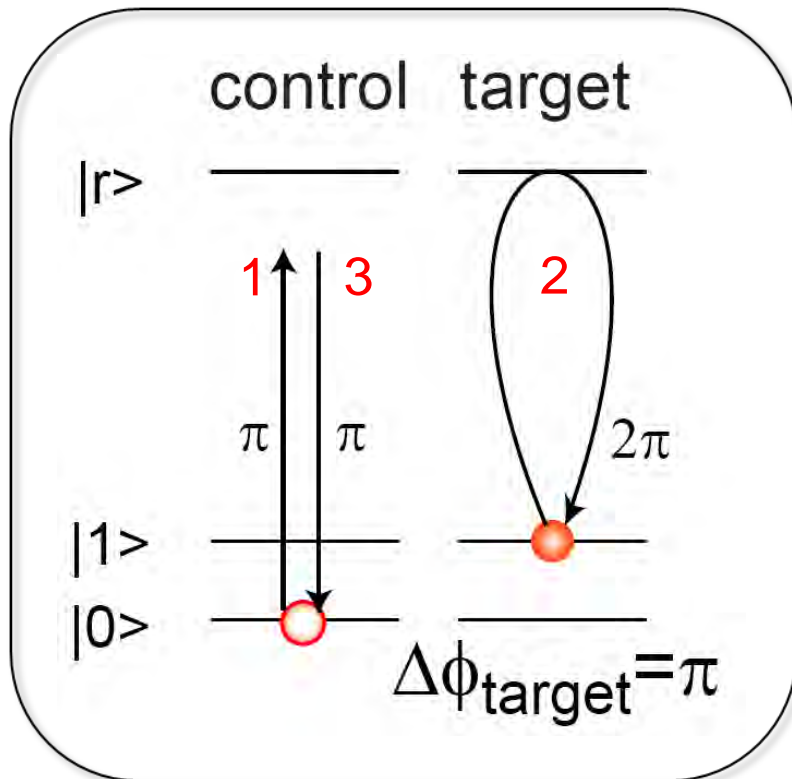
National Institute of Standards and Technology, Gaithersburg, Maryland 20899

R. Côté¹ and M.D. Lukin²

¹*Physics Department, University of Connecticut, 2152 Hillside Road, Storrs, Connecticut 06269-3046*

²*ITAMP, Harvard-Smithsonian Center for Astrophysics, Cambridge, Massachusetts 02138*

(Received 7 April 2000)



Fast Quantum Gates for Neutral Atoms

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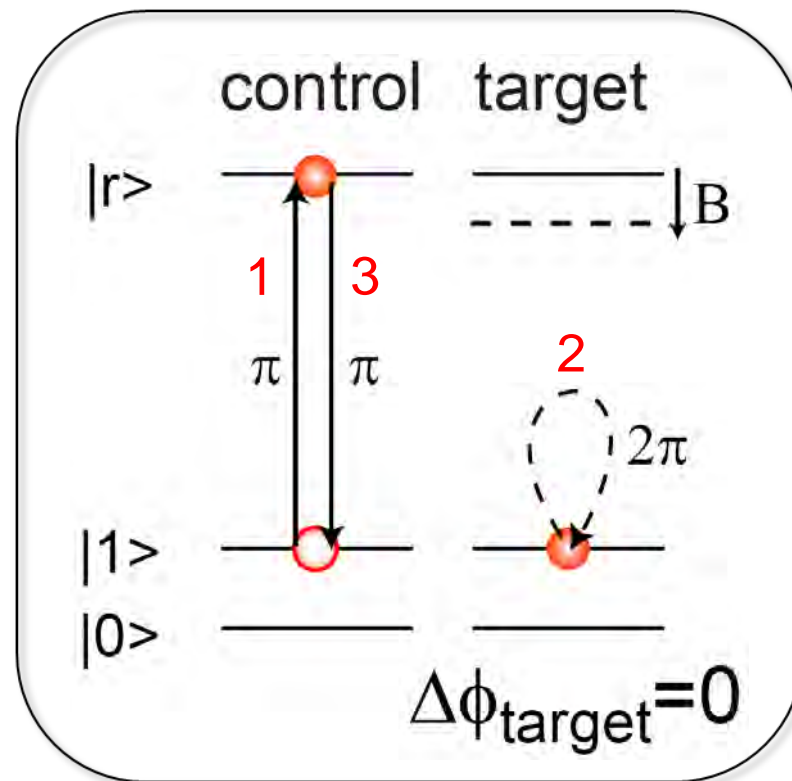
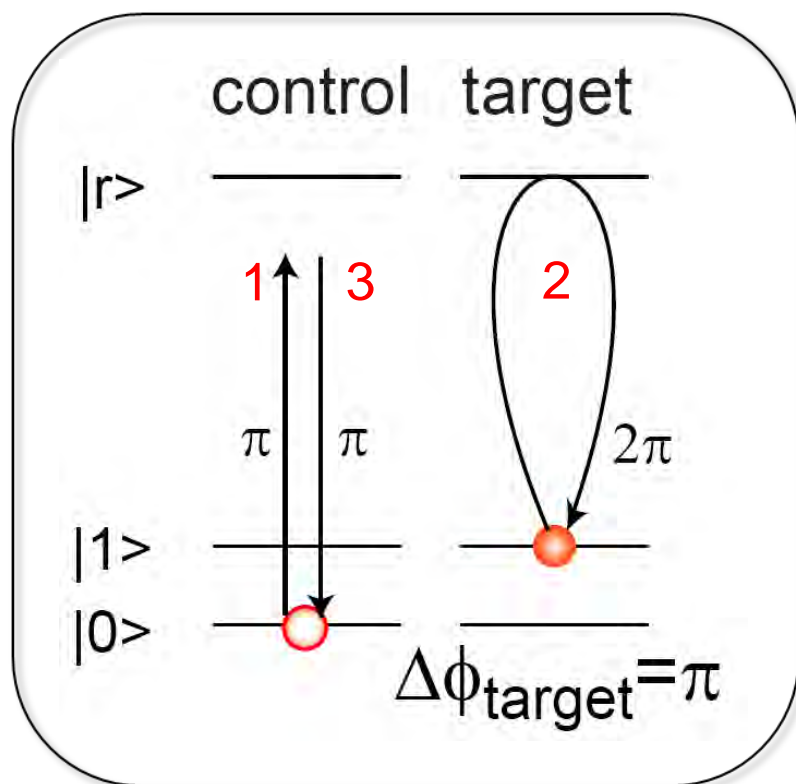
National Institute of Standards and Technology, Gaithersburg, Maryland 20899

R. Côté¹ and M.D. Lukin²

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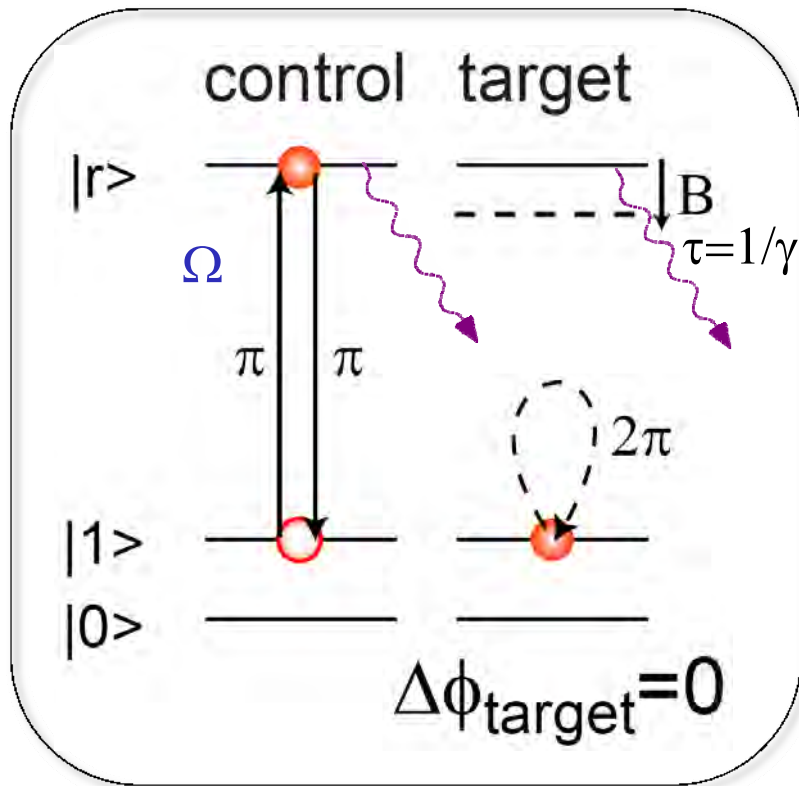
²*ITAMP, Harvard-Smithsonian Center for Astrophysics, Cambridge, Massachusetts 02138*

(Received 7 April 2000)



$$\text{C-phase gate} = U_{\pi}$$

Error scaling of Rydberg gate



residual excitation error $\sim \Omega^2/B^2$

Rydberg lifetime τ gives error $\sim 1/\Omega\tau$

optimum Rabi frequency $\Omega \sim B^{2/3}/\tau^{1/3}$

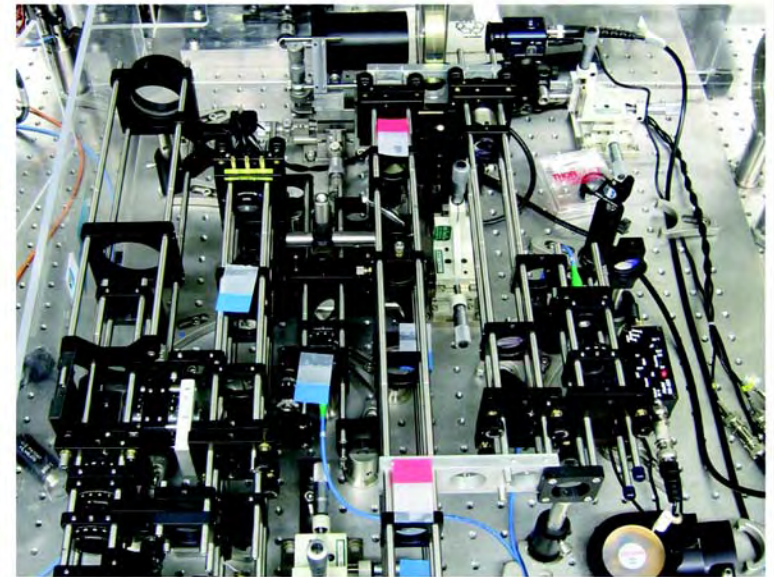
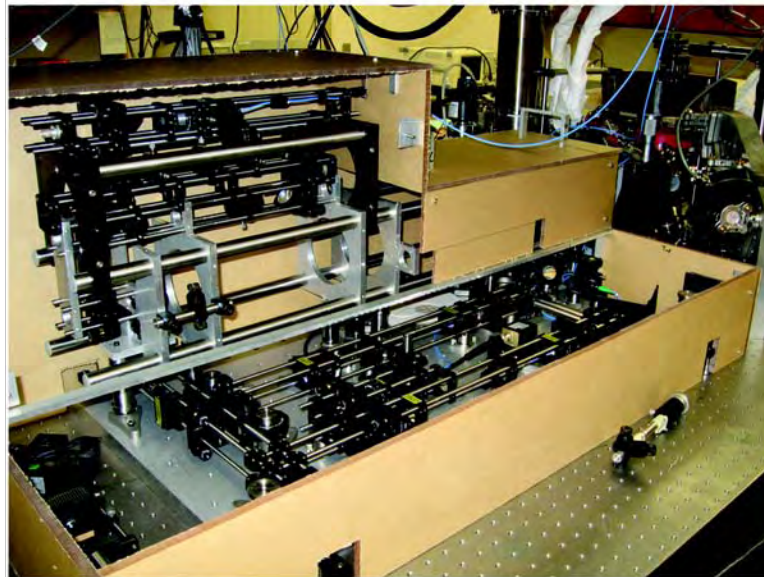
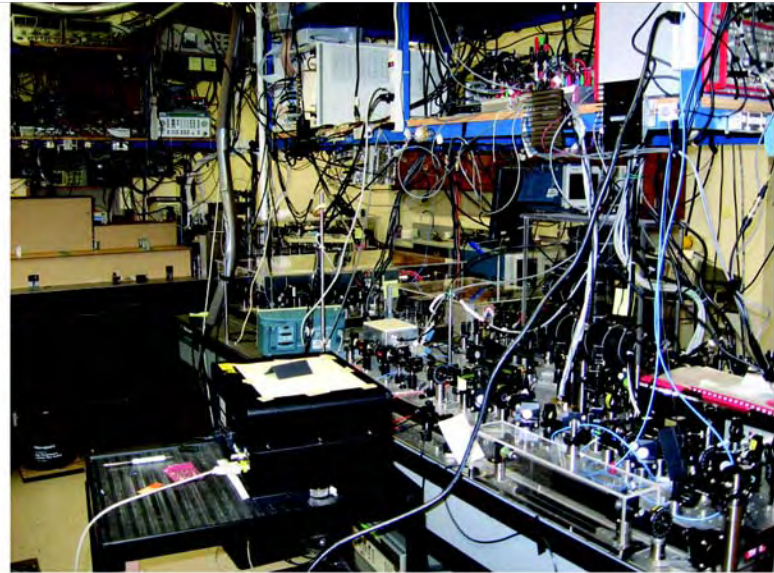
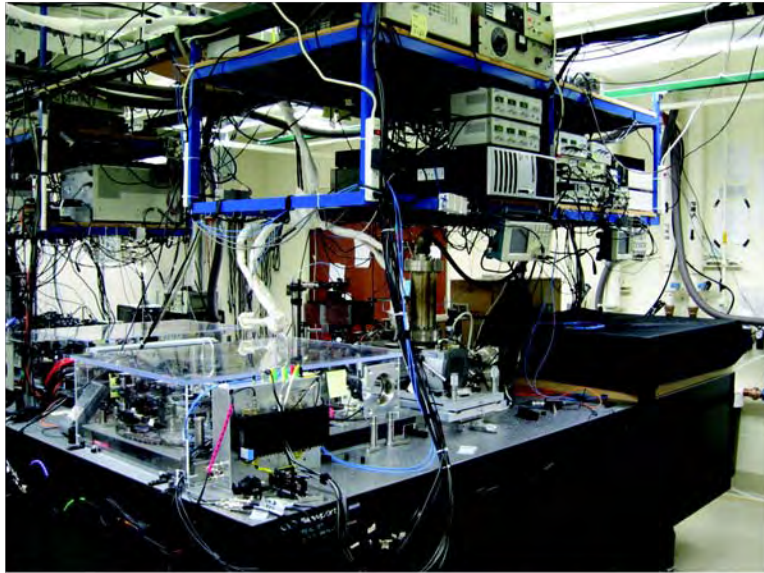
minimum error $\sim 1/(B\tau)^{2/3}$

Rb atoms, $n=100$ and $R=10\text{ }\mu\text{m}$
 $B \sim 25\text{ MHz}$, $\tau \sim 300\text{ }\mu\text{s}$

Error ~ 0.003

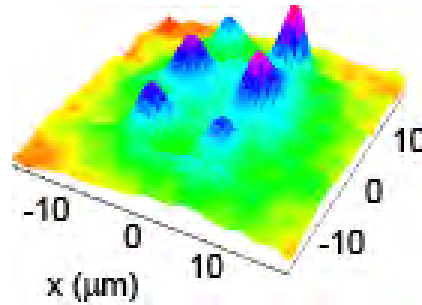
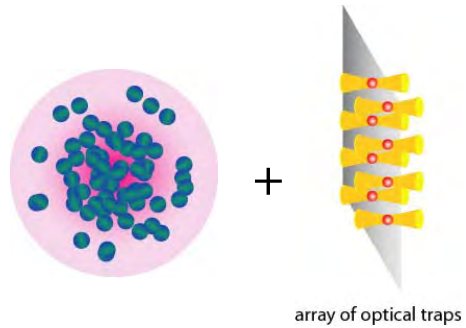
Saffman & Walker PRA (2005)

Experiment

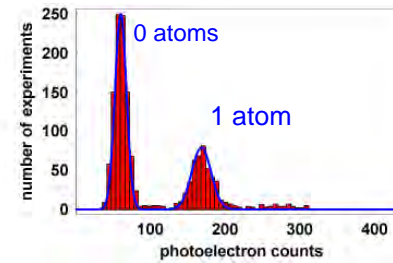


Experimental approach

MOT + Optical traps
(FORT)

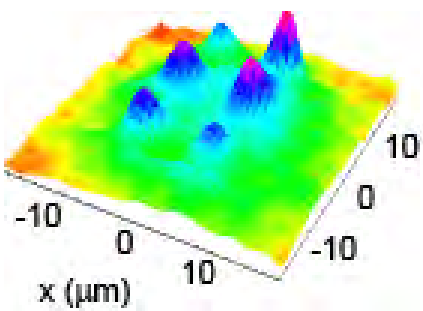
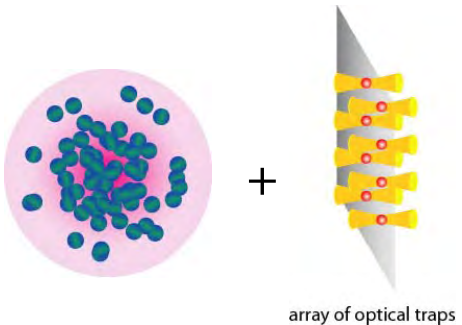


Single atom detection

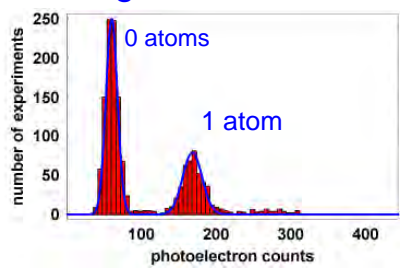


Experimental approach

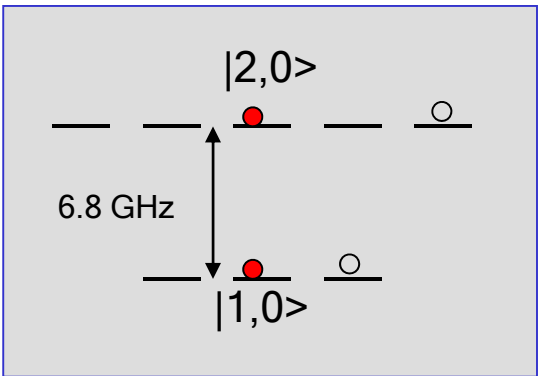
MOT + Optical traps
(FORT)



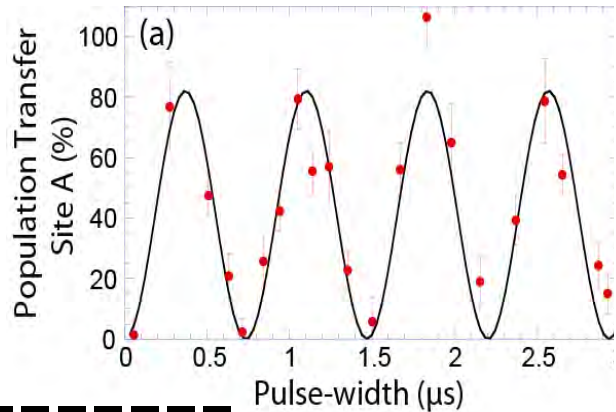
Single atom detection



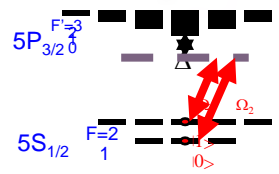
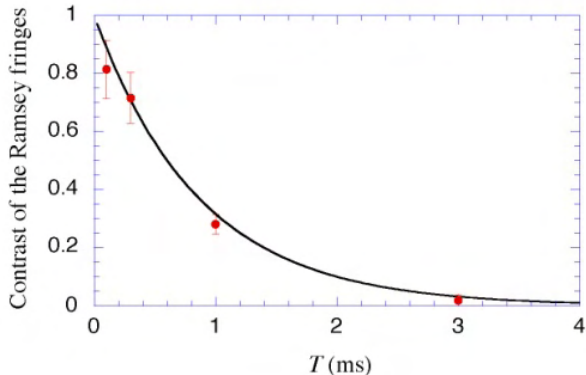
^{87}Rb hyperfine qubit



Ground Rabi flopping
 $\Omega/2\pi=1.4$ MHz



Ramsey measurement
 $T_2 \sim 1$. ms

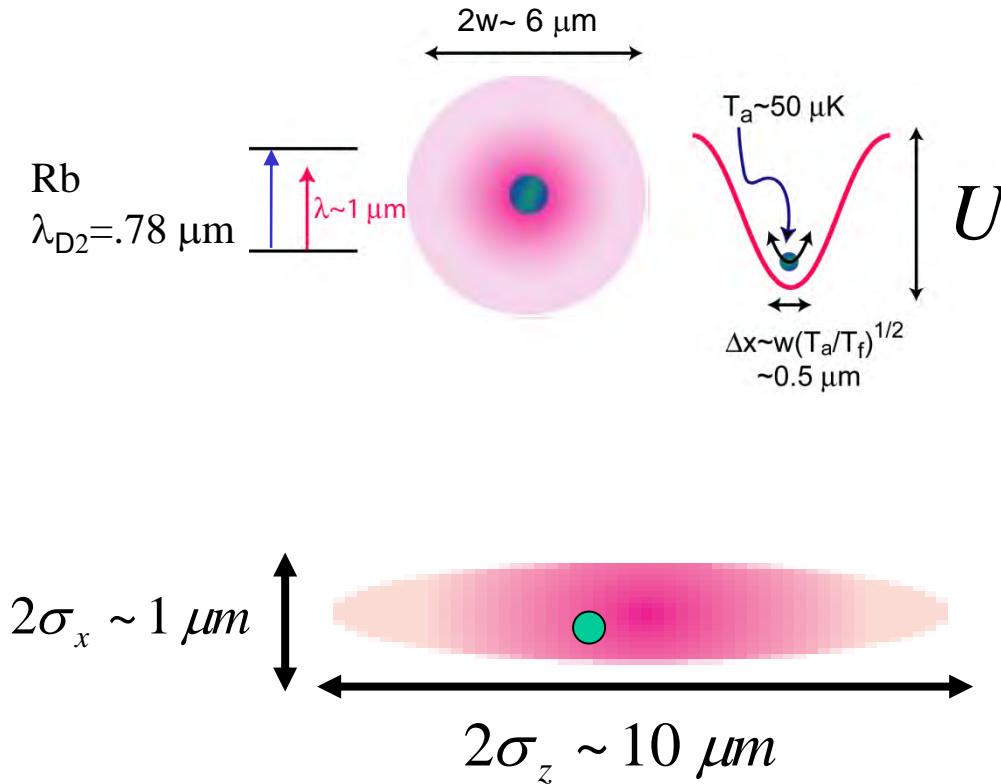


Yavuz, et al. PRL (2006)

Microscopic optical trap

Single beam dipole trap $-U \sim \text{Intensity}$

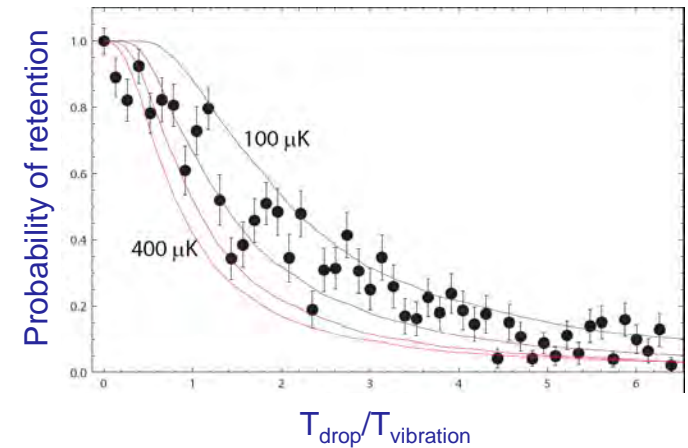
Pioneering experiments by
Grangier,
Meschede, groups



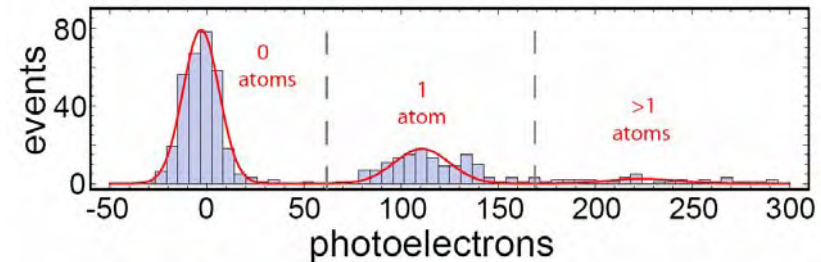
$$w = 3.0 \mu\text{m}, \quad \lambda = 1.06 \mu\text{m}$$

$$T_a = 0.1 - 0.5 \text{ mK}, \quad U_f / k_B = 5 \text{ mK}$$

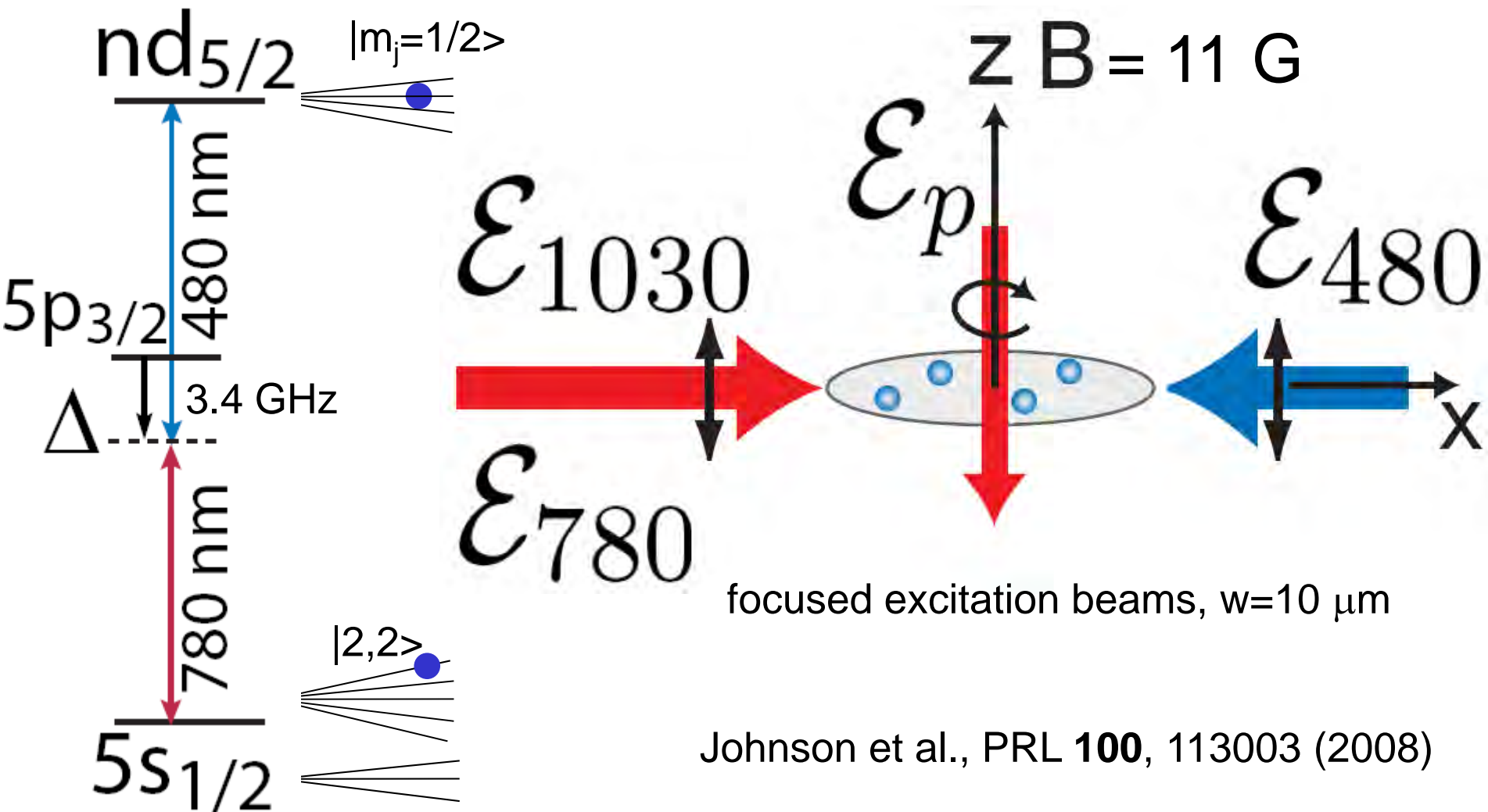
Atom temperature measured by drop and recapture. Typically $T \sim 200 \mu\text{K}$



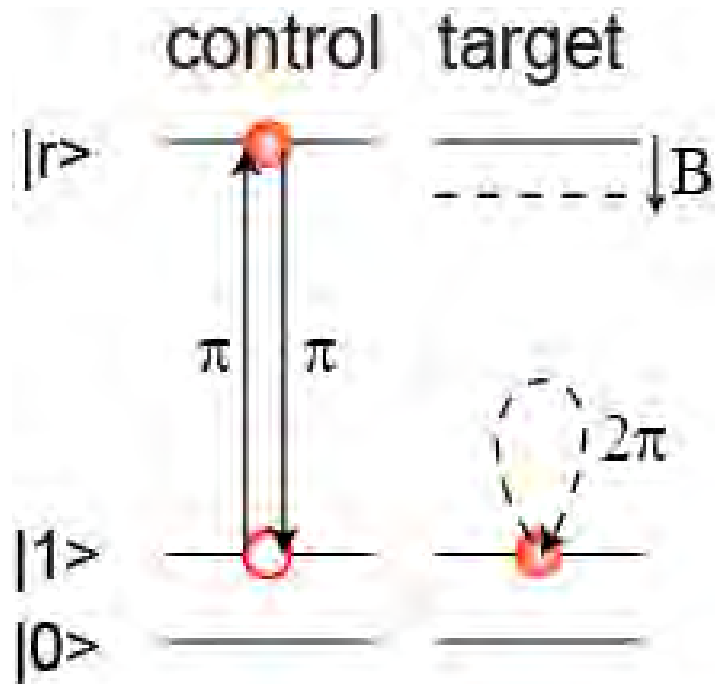
Single atom readout with 3D molasses



Rydberg excitation geometry

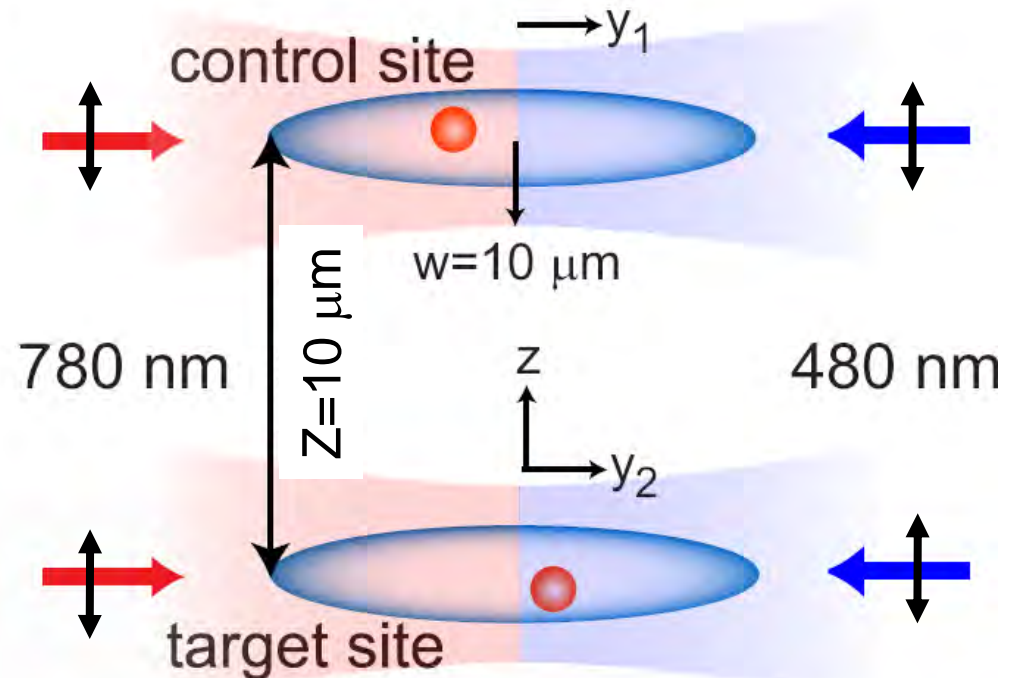
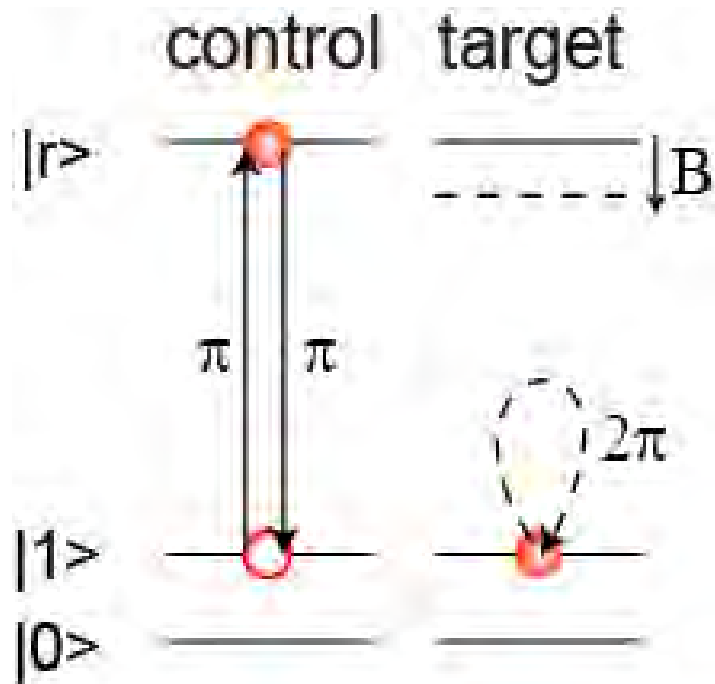


Rydberg blockade between two atoms



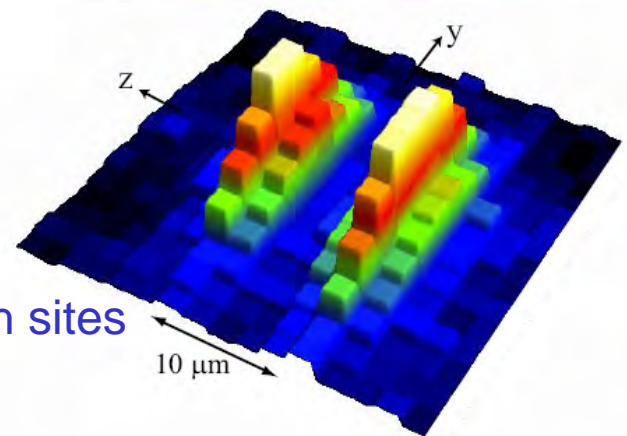
Jaksch, et al. 2000

Rydberg blockade between two atoms



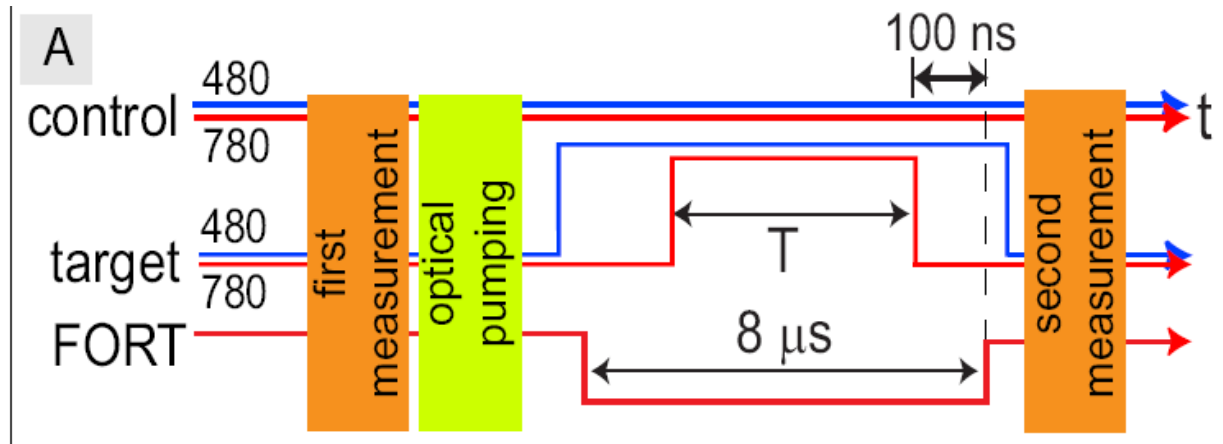
Jaksch, et al. 2000

Average of 146 shots.
Probability of one atom in both sites
about 5-10%.



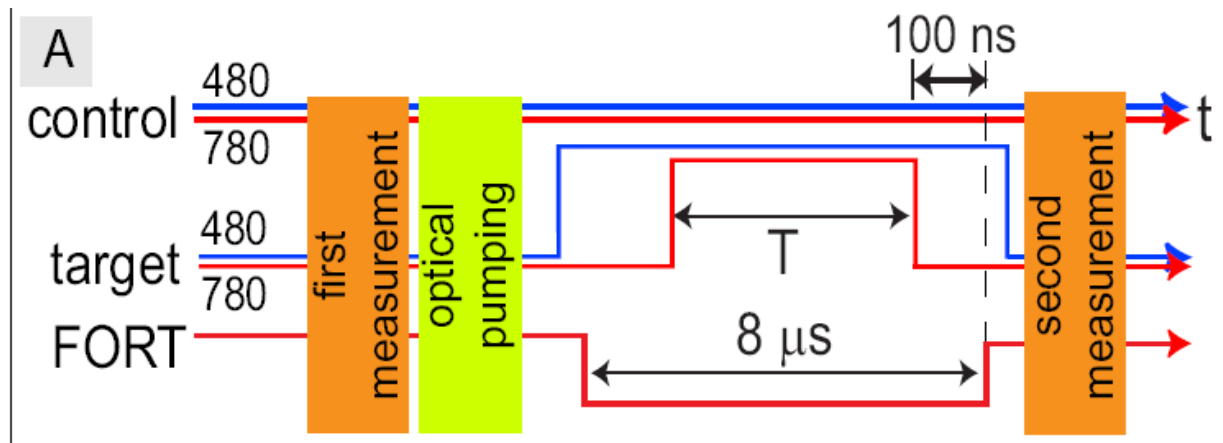
Rabi oscillations

$79\text{d}_{5/2}$

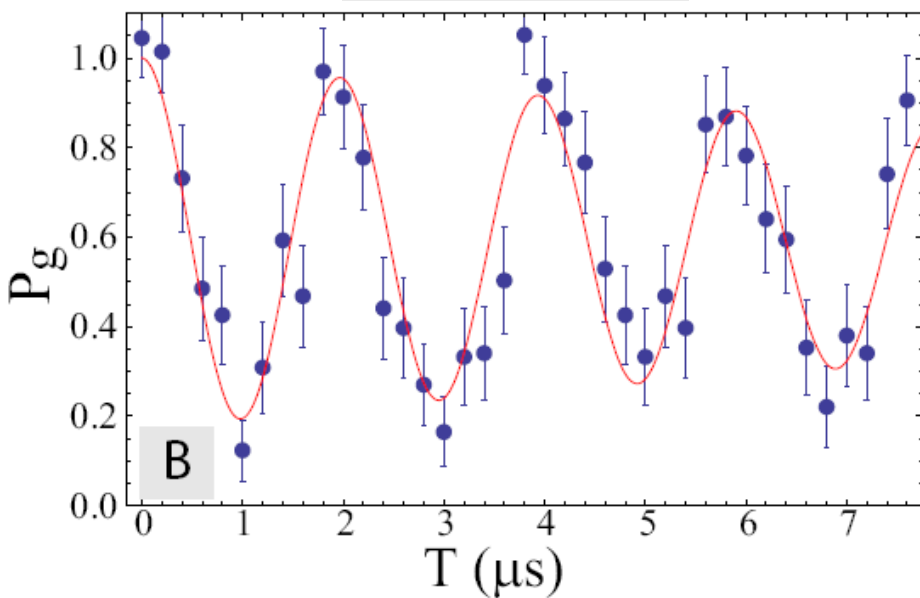


Rabi oscillations

$79d_{5/2}$



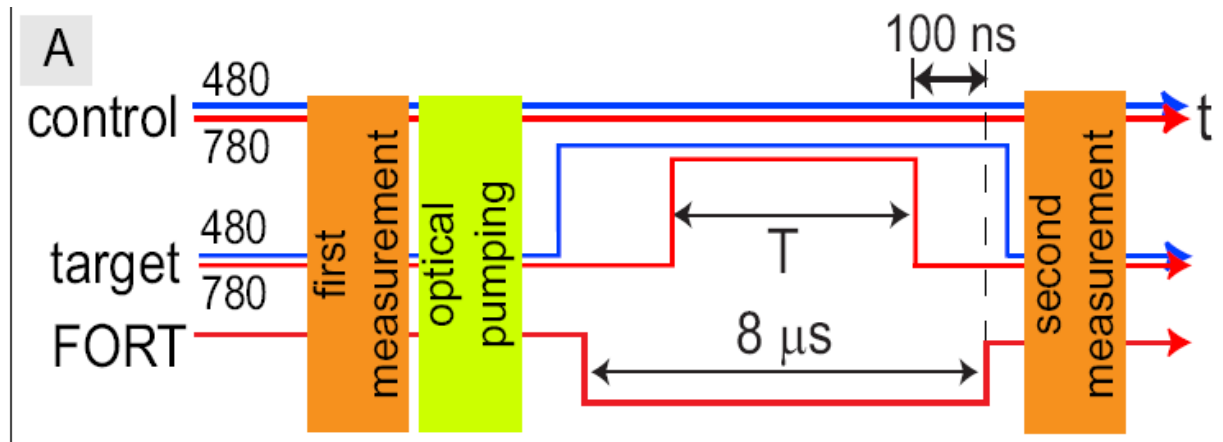
Targeted site



$\Omega/2\pi = 0.51$ MHz (exp.), 0.59 MHz (theory)

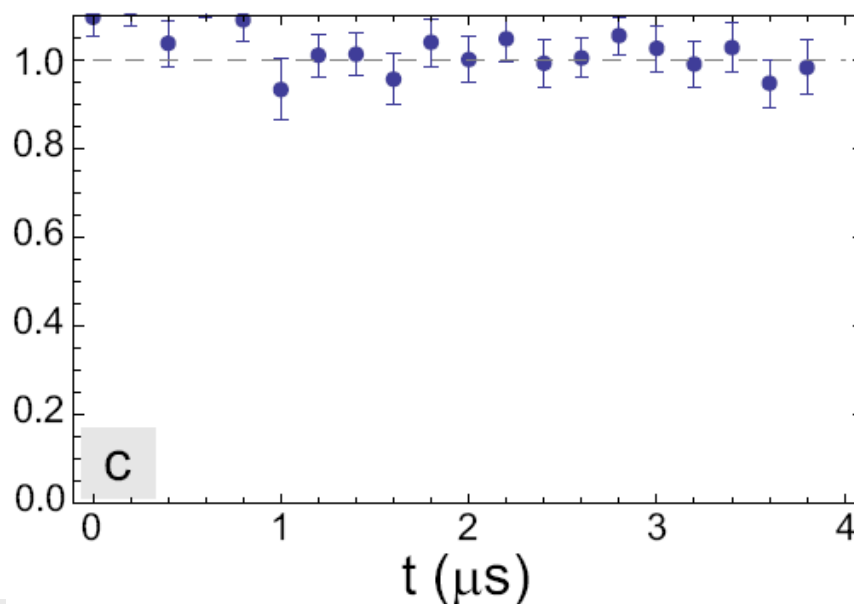
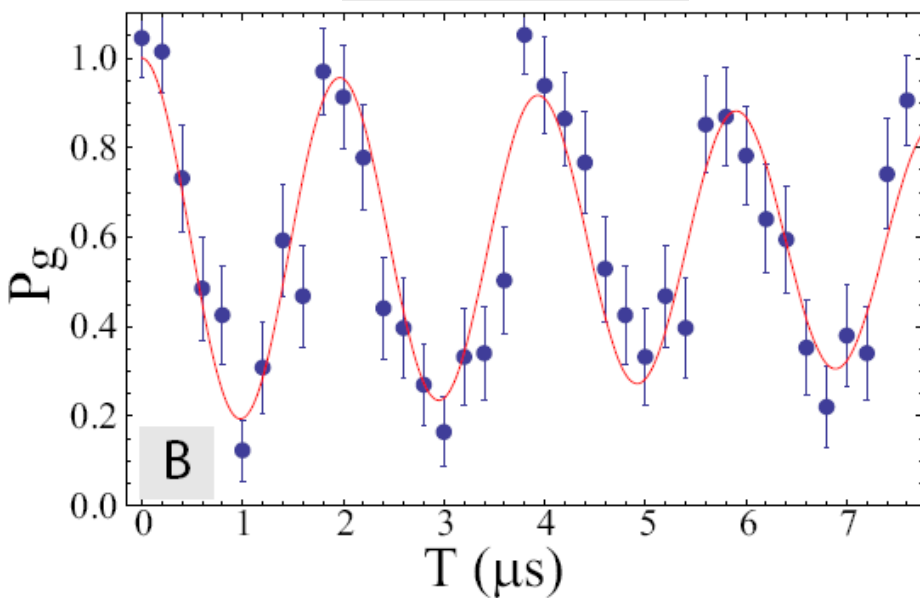
Rabi oscillations

$79d_{5/2}$



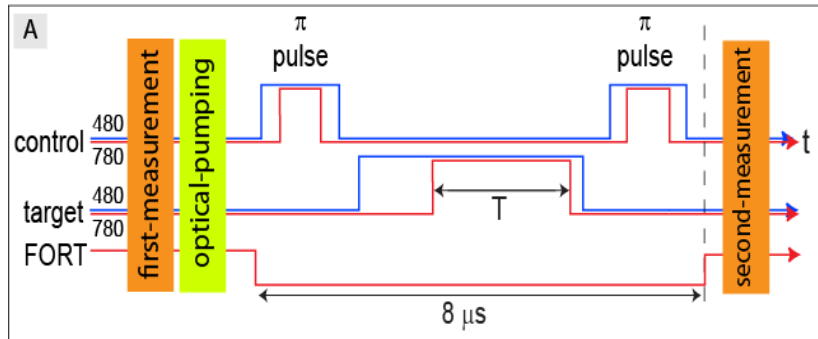
Targeted site

crosstalk at the other site



$\Omega/2\pi = 0.51$ MHz (exp.), 0.59 MHz (theory)

Rydberg Blockade

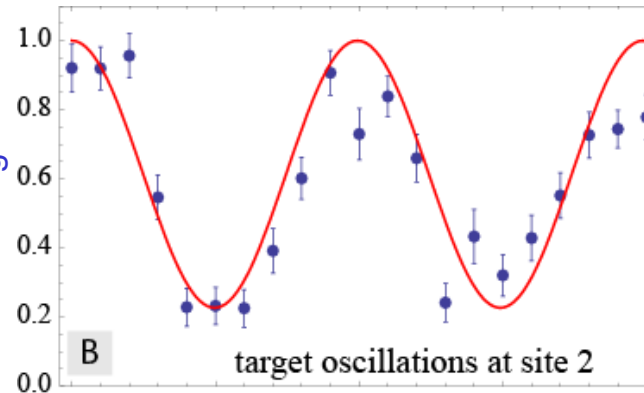


$79d_{5/2}$

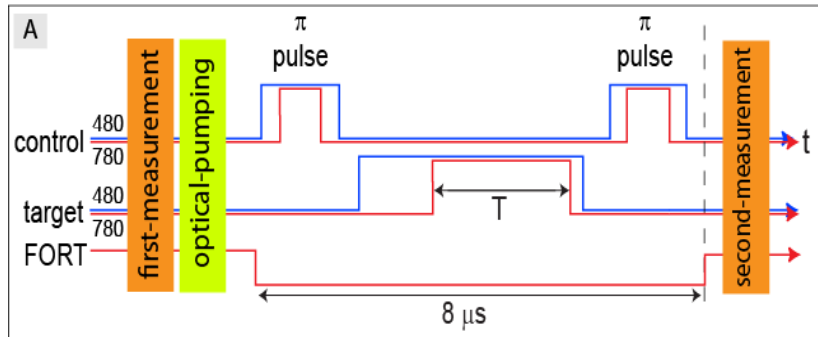
no control atom
target oscillations

site
1 \rightarrow 2

P_g



Rydberg Blockade

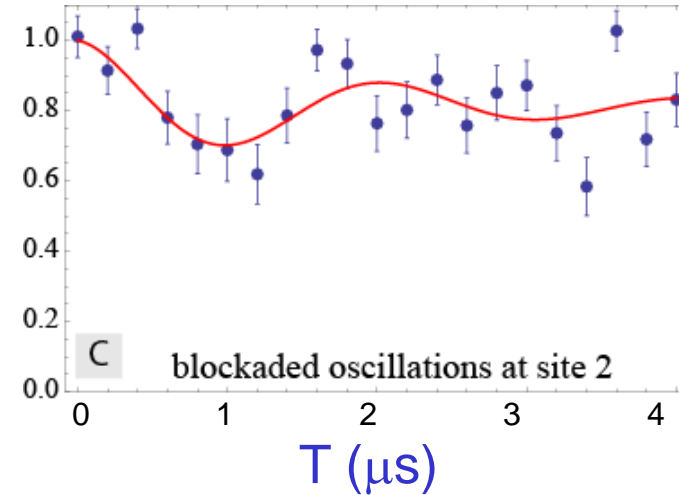
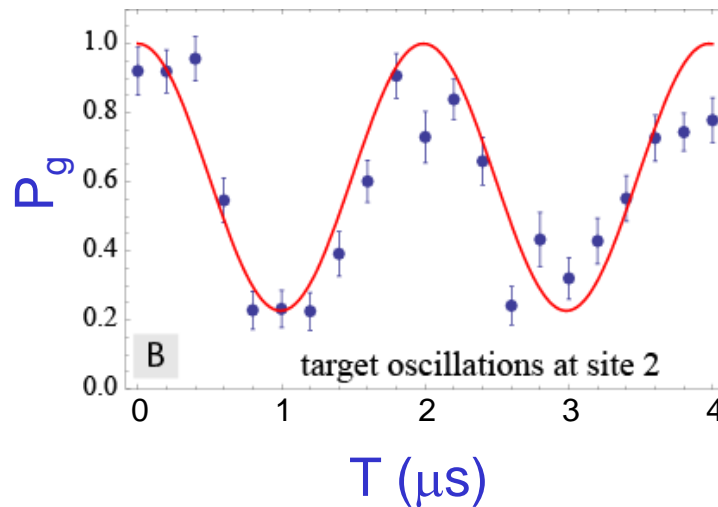


$79d_{5/2}$

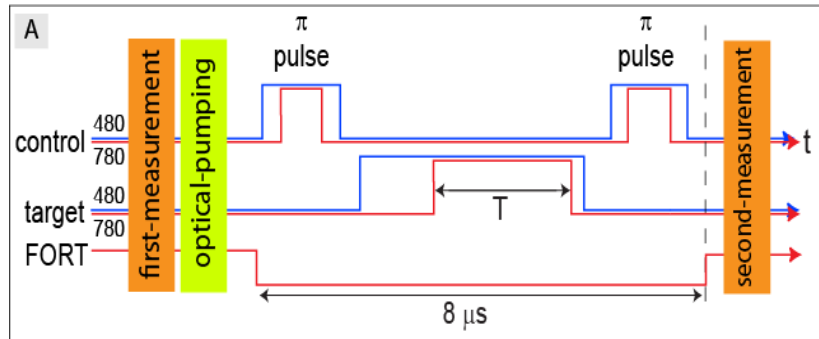
no control atom
target oscillations

with control atom
target blocked

site
1 \rightarrow 2



Rydberg Blockade



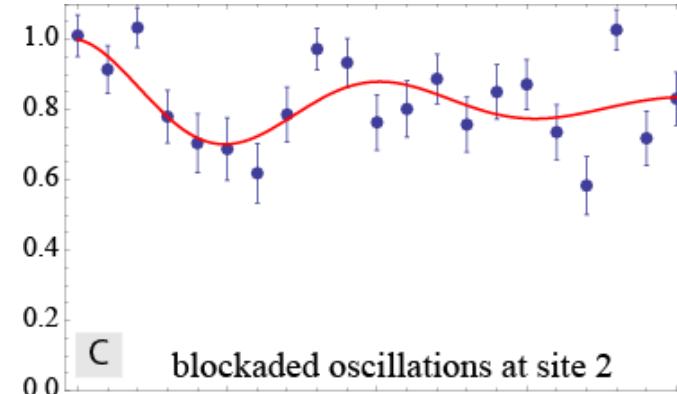
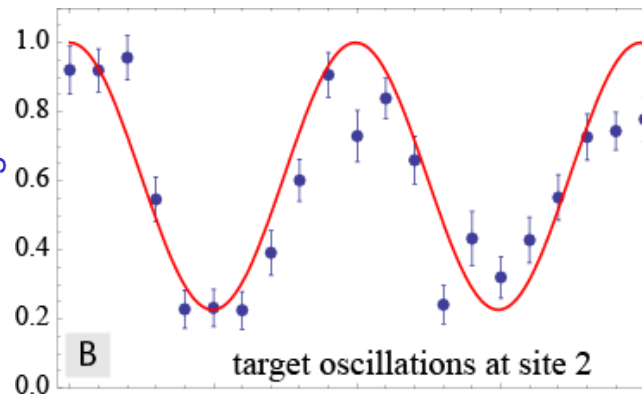
$79d_{5/2}$

no control atom
target oscillations

with control atom
target blocked

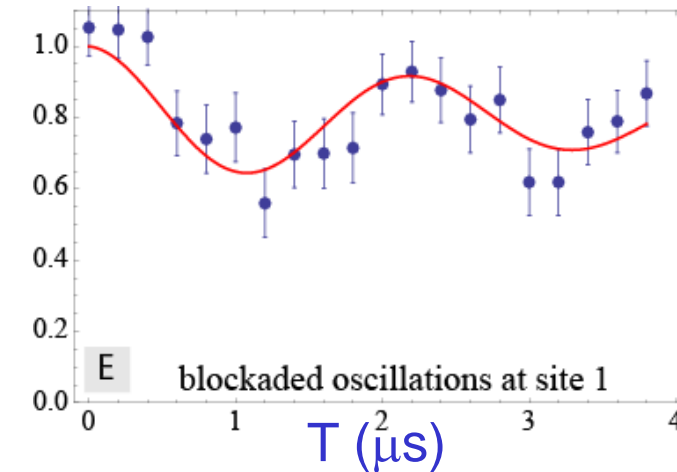
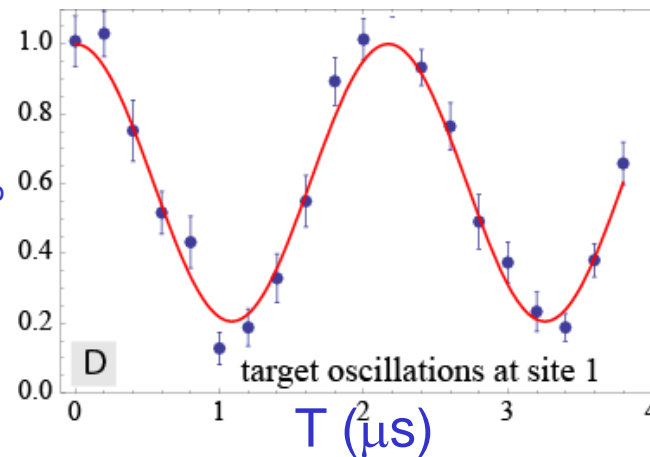
site
1 \rightarrow 2

P_g



site
2 \rightarrow 1

P_g

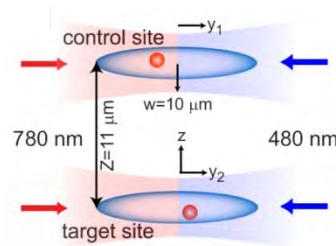


Higher fidelity blockade

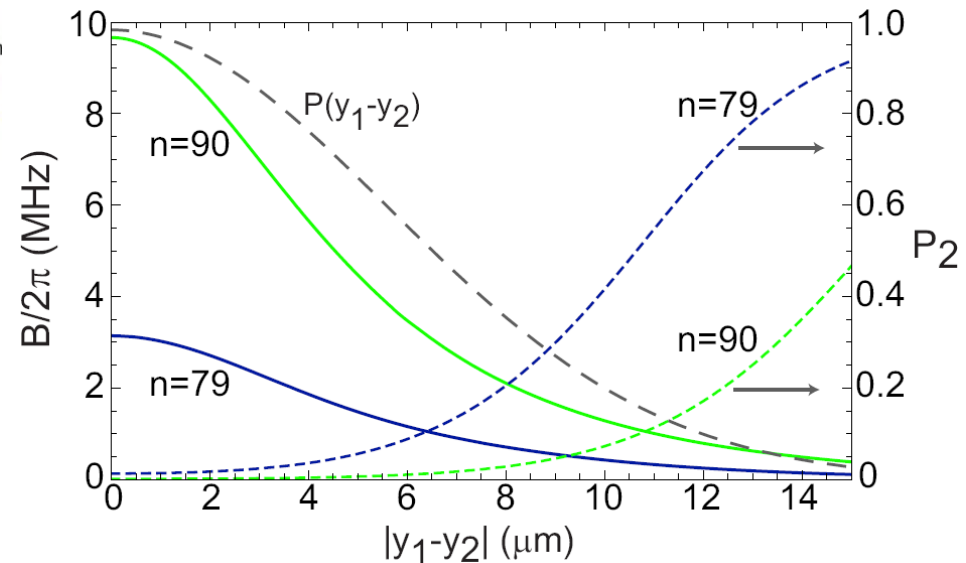
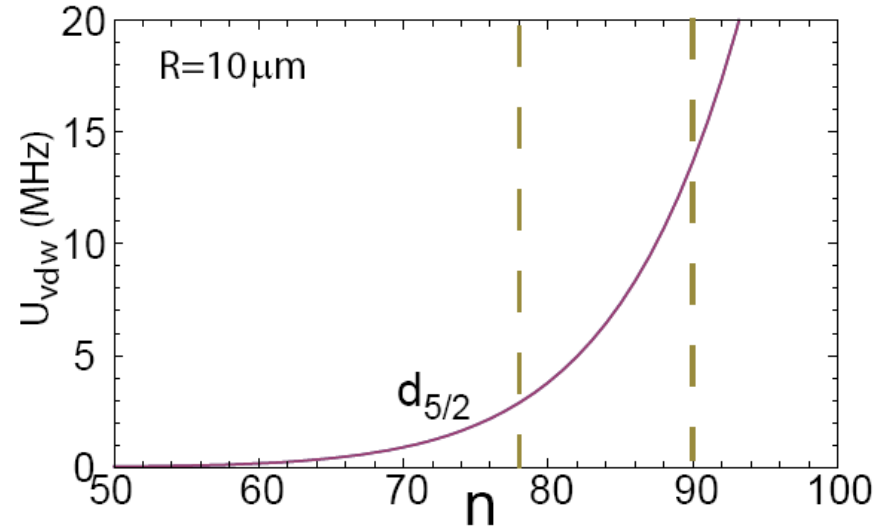
We can do better by increasing the interaction strength

Interaction strength vs. n

$$B \sim n^{11} - n^{12}$$

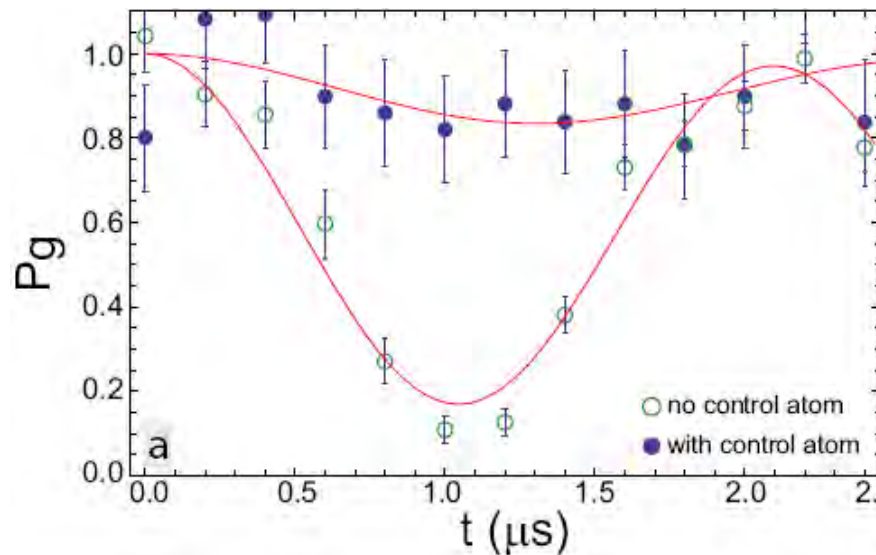


Predicted blockade and P_2
for $n=90$

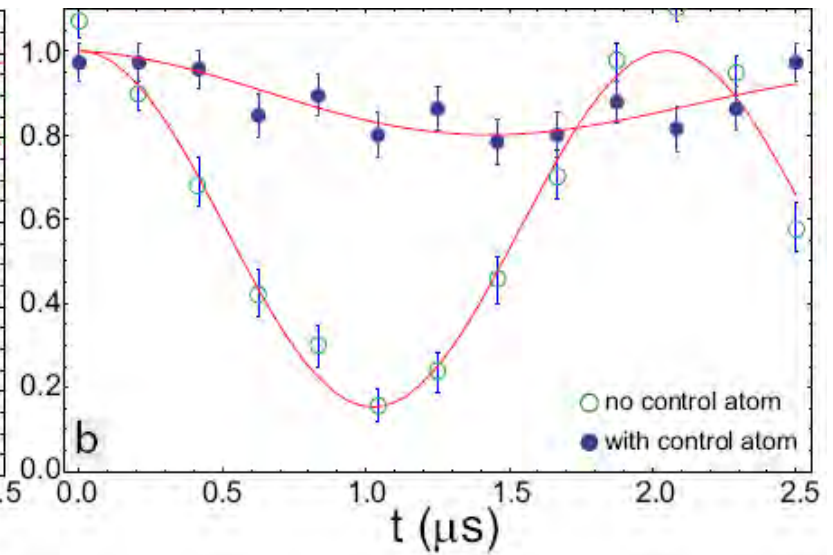


Blockade experiment n=90

Experiment



Monte Carlo Simulation



The data are fit to
$$P_g = (1 - a) + a \cos(2\pi f t) e^{-t/\tau}$$

With a control atom present we find

$a = 0.09$ (experiment)

$a = 0.11$ (simulation)

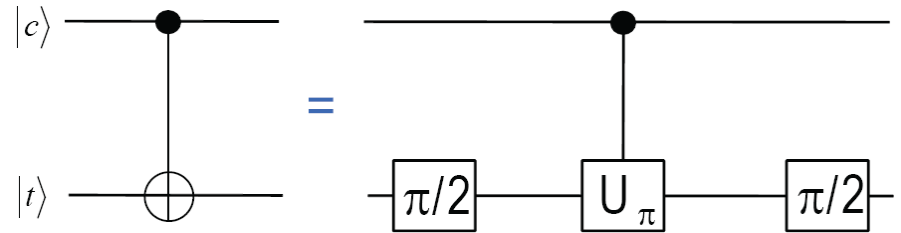
Residual blockade leakage is almost entirely due to state preparation and measurement errors.

Urban, et al. Nature Physics Feb. (2009)

Gaetan, et al. Nature Physics Feb. (2009)

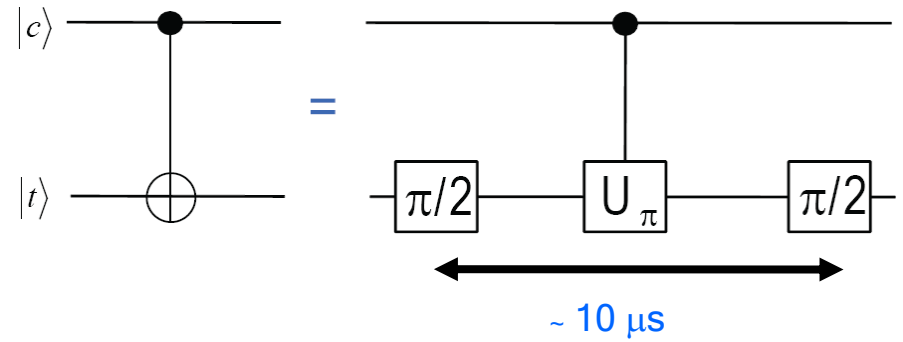
From blockade to quantum gates

The CNOT gate can be realized
with a controlled phase plus Hadamards

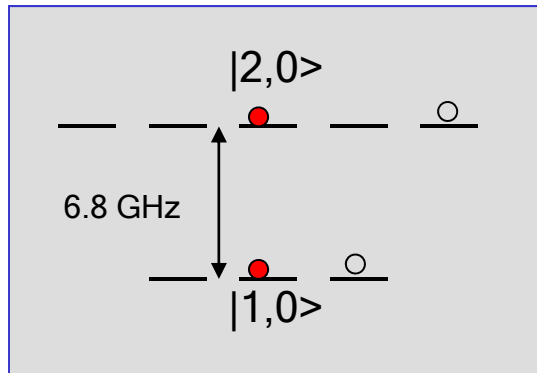


From blockade to quantum gates

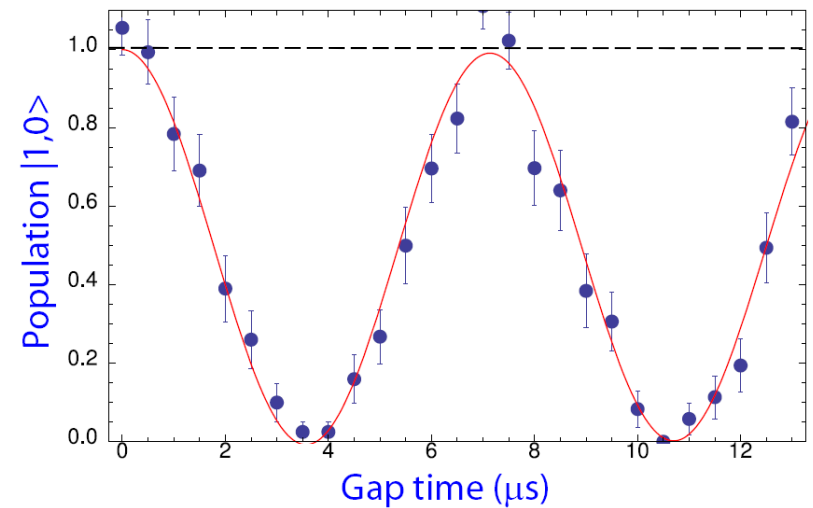
The CNOT gate can be realized with a controlled phase plus Hadamards



Better ground state coherence with $m=0$ states at B~4 G.



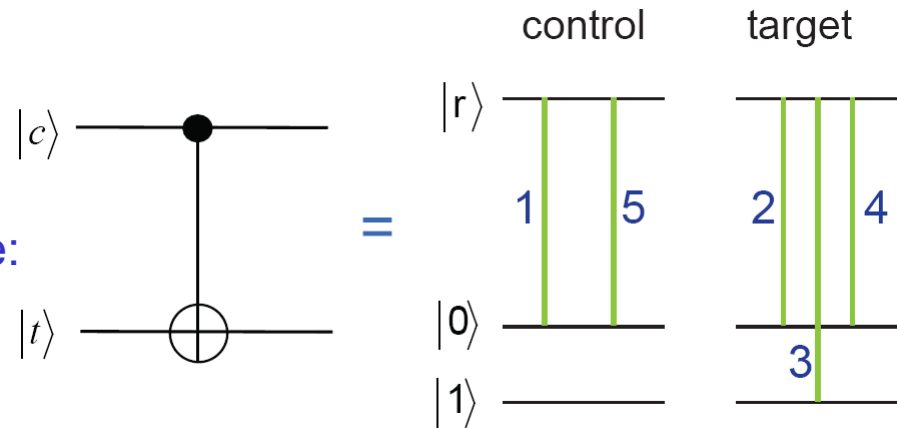
Ground state Ramsey spectroscopy



Other CNOT protocols

It is also possible to perform a CNOT gate using only π pulses and blockade:

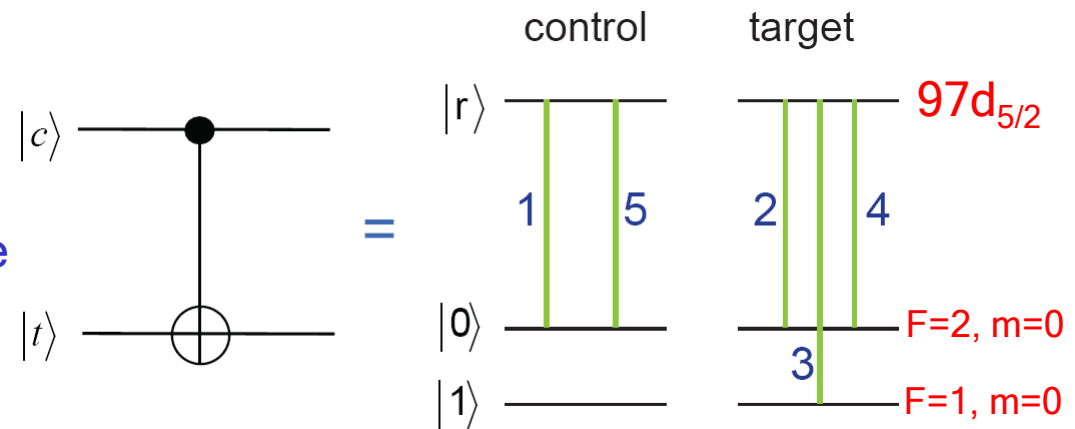
Ohlsson, et al. Opt. Commun. (2002)



Other CNOT protocols

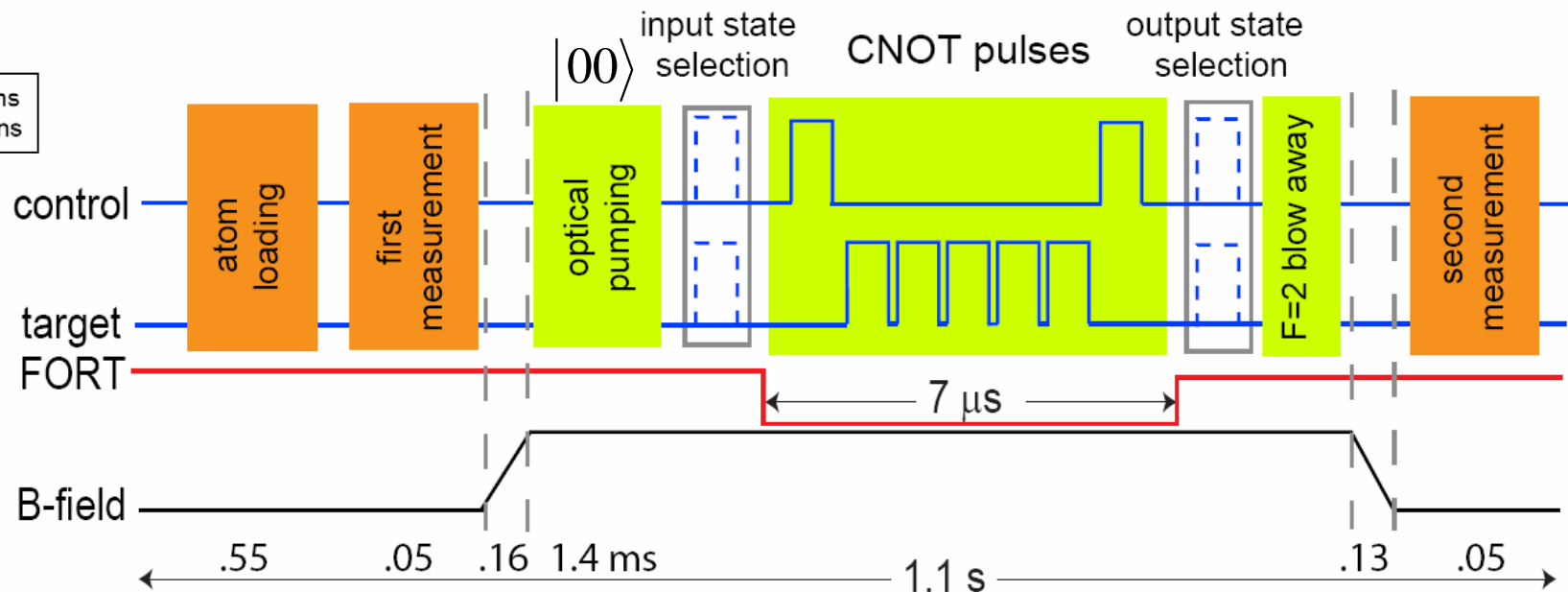
It is also possible to perform a CNOT gate using only π pulses and blockade

Ohlsson, et al. Opt. Commun. (2002)



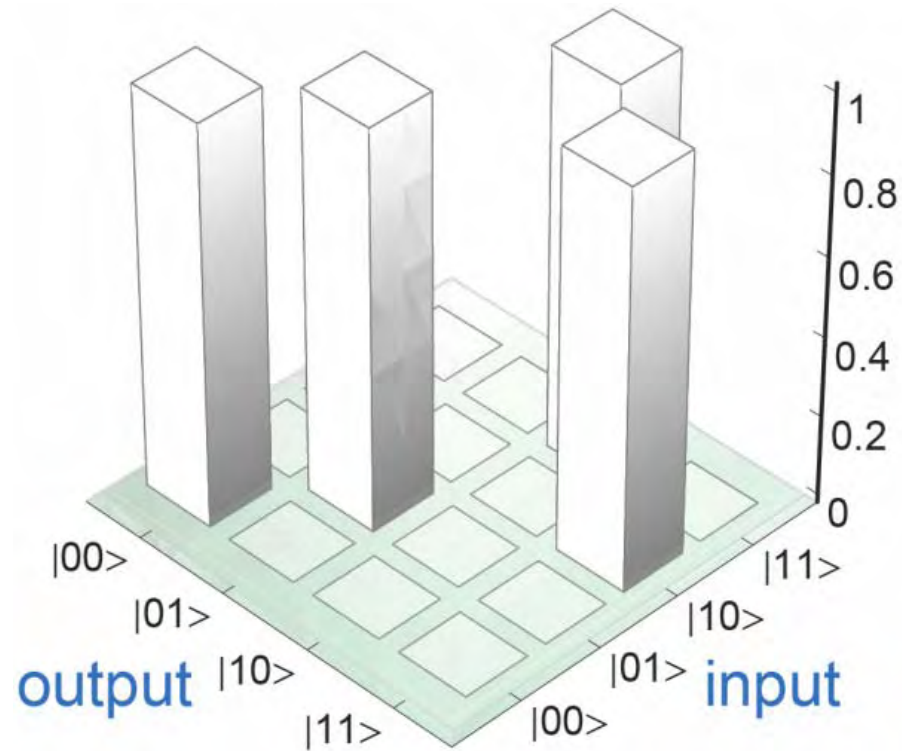
Experimental sequence

ground $t_\pi = 550$ ns
Rydberg $t_\pi = 750$ ns



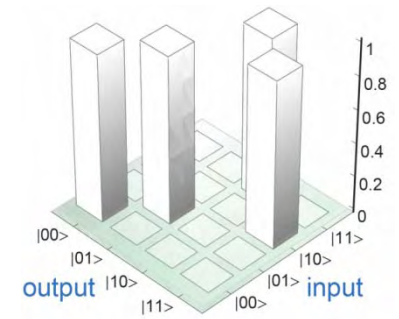
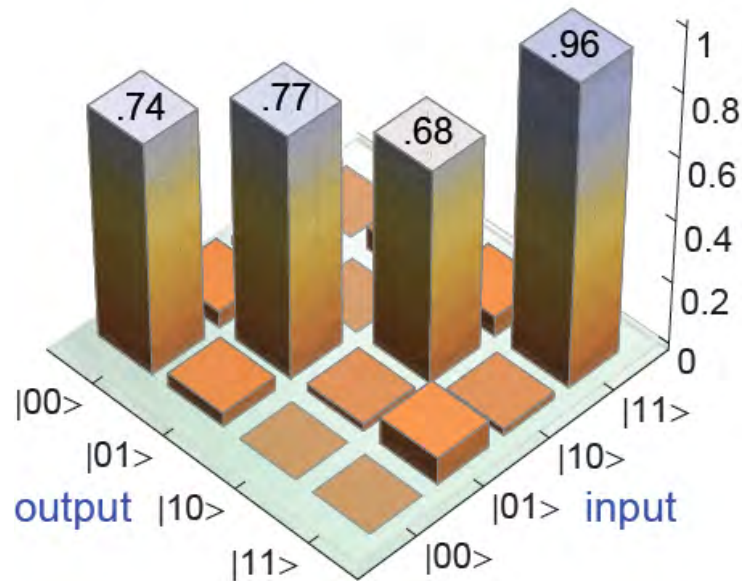
CNOT experiment

ideal
truth table



CNOT experiment

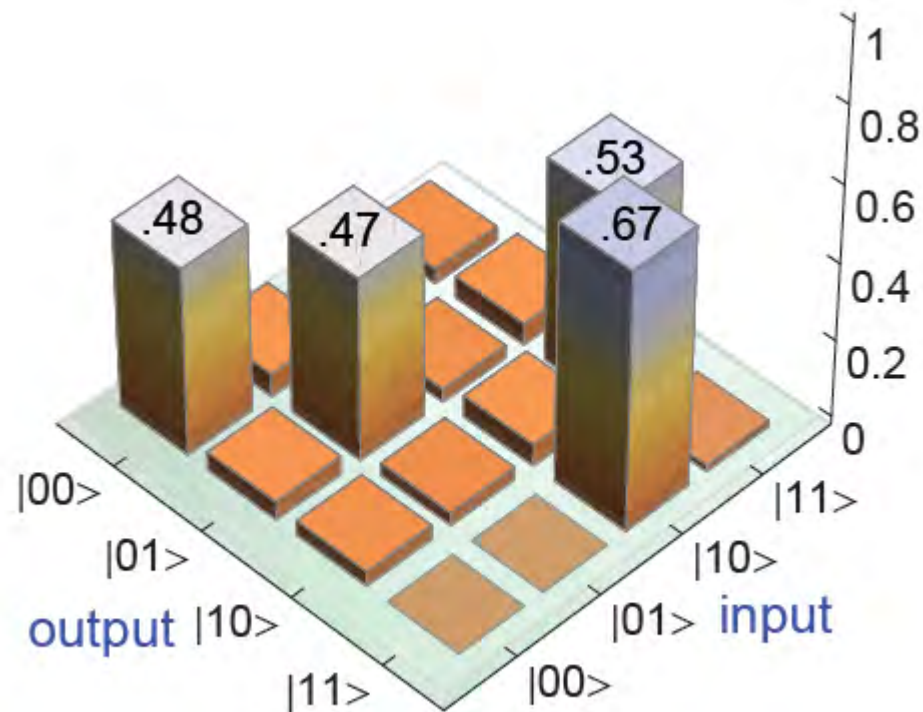
measured
input states



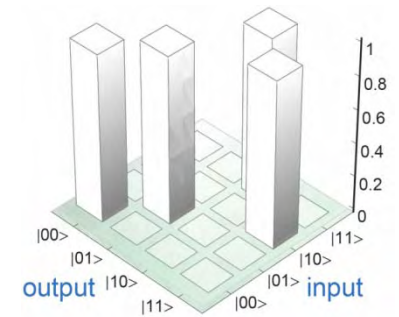
Preliminary data May 2009

CNOT experiment

measured
truth table



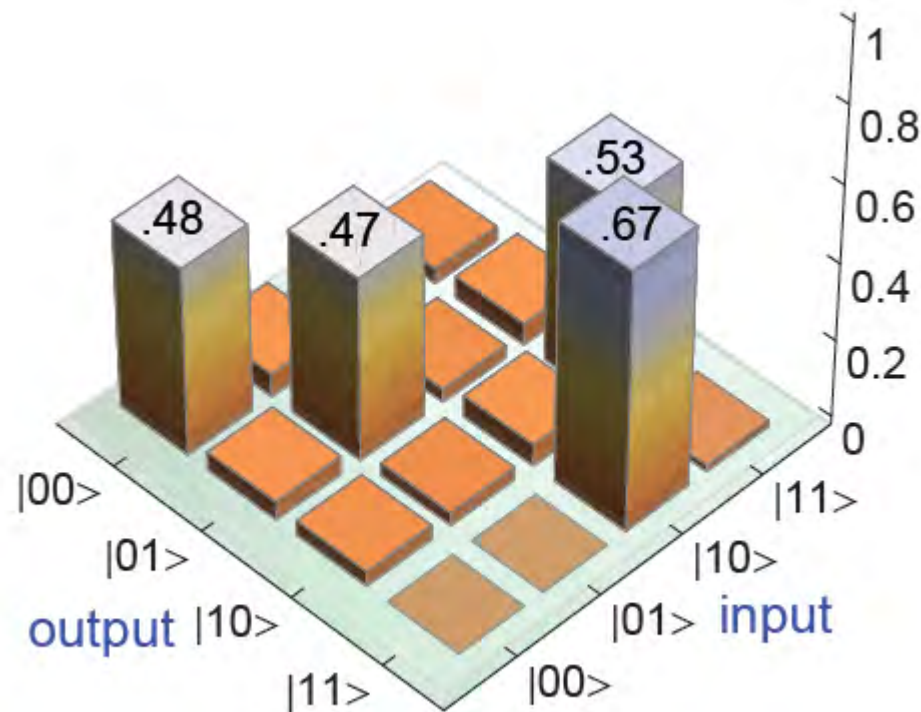
about 75 two-atom trials/ matrix element



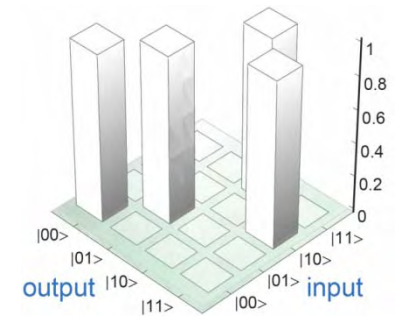
Preliminary data May 2009

CNOT experiment

measured
truth table



about 75 two-atom trials/ matrix element



Probabilities

$$\begin{pmatrix} 0.48 & 0.06 & 0 & 0.04 \\ 0.05 & 0.47 & 0.03 & 0.06 \\ 0.04 & 0.03 & 0.06 & 0.53 \\ 0 & 0 & 0.67 & 0.02 \end{pmatrix}$$

average high/low

$$.54 / .03 = 17 / 1$$

fidelity

$$(1/4) \text{Tr}(U_{ideal}^T U_{exp}) = .54$$

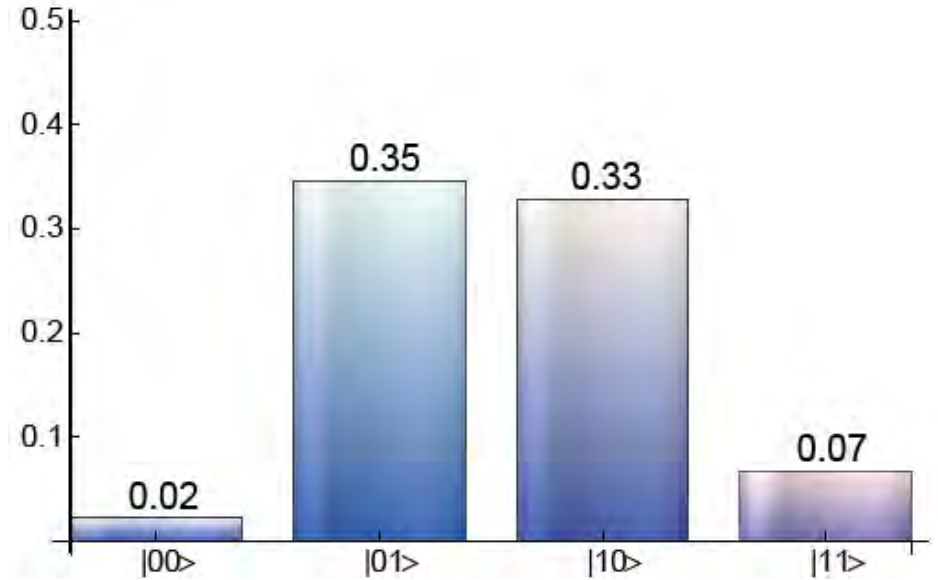
Preliminary data May 2009

Entanglement on demand

$$\frac{\overset{\text{control}}{|0\rangle} + \overset{\text{target}}{|1\rangle}}{\sqrt{2}} |1\rangle \Rightarrow \frac{|01\rangle + |10\rangle}{\sqrt{2}}$$

input
data

ideal
output



Entanglement on demand

$$\frac{\text{control } |0\rangle + |1\rangle}{\sqrt{2}} \text{ target } |1\rangle \Rightarrow \frac{|01\rangle + |10\rangle}{\sqrt{2}}$$

input
data

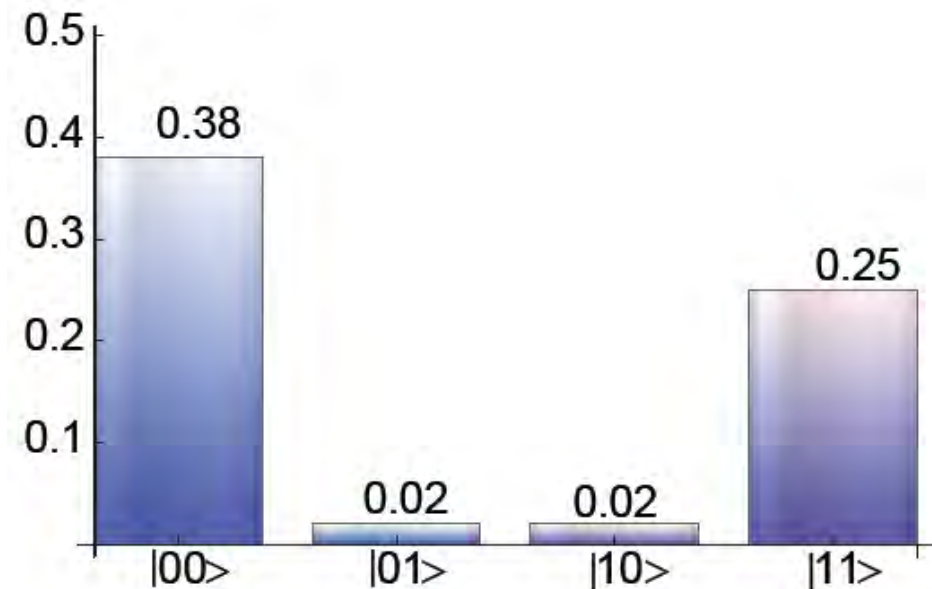
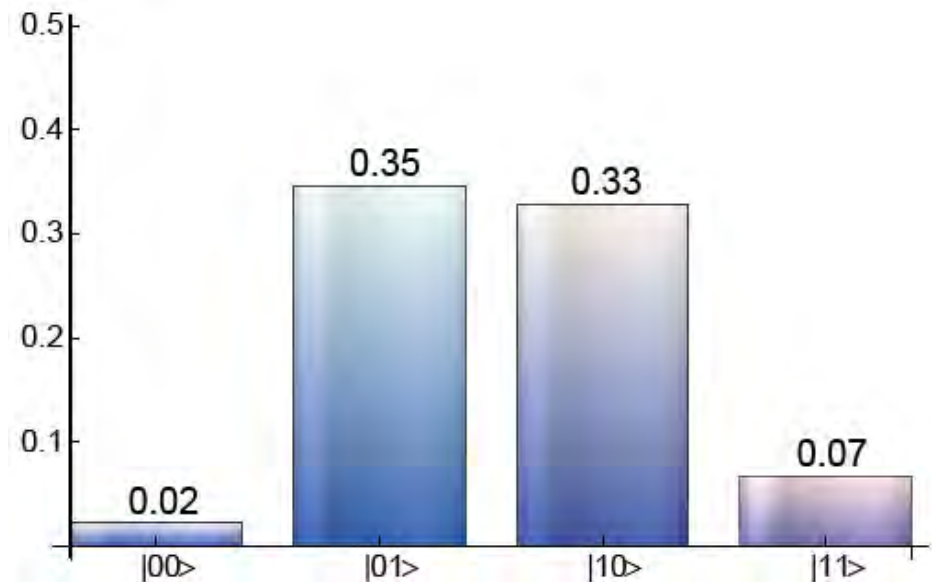
ideal
output

$$\frac{\text{control } |0\rangle + |1\rangle}{\sqrt{2}} \text{ target } |0\rangle \Rightarrow \frac{|00\rangle + |11\rangle}{\sqrt{2}}$$

input
data

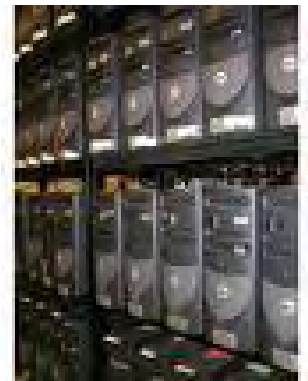
ideal
output

Publication in preparation

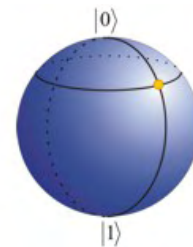
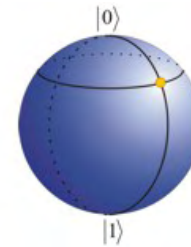
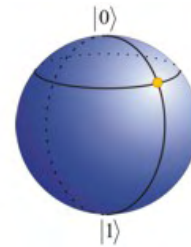
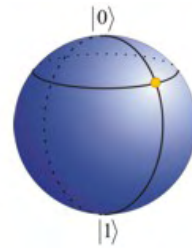
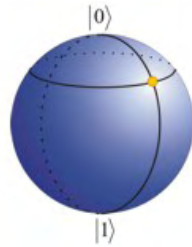
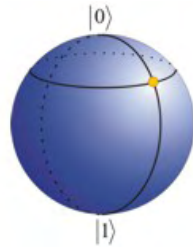
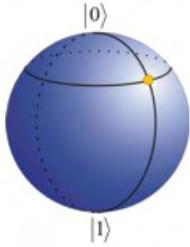
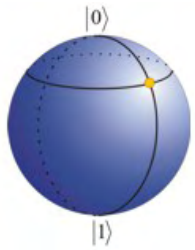


Preliminary data May 2009

Scalability



Scalability



Ion traps $N=8$ qubits (2005)

Neutral atoms

Superconductors

Semiconductor spins

Photons and linear optics

Nitrogen vacancies

.....

Applications:

Entanglement and decoherence

Quantum simulation

Factoring, search, ...

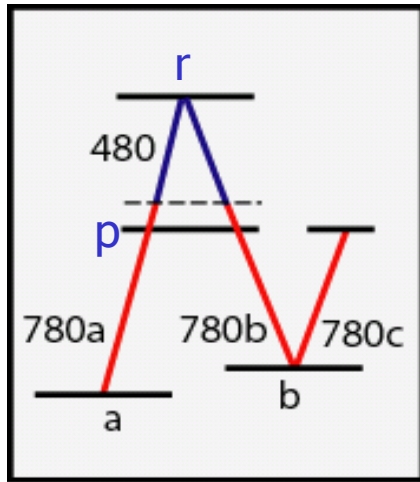
Quantum networks, secure data transmission

$N > 2$

$N > 30$

$N > 10^6$

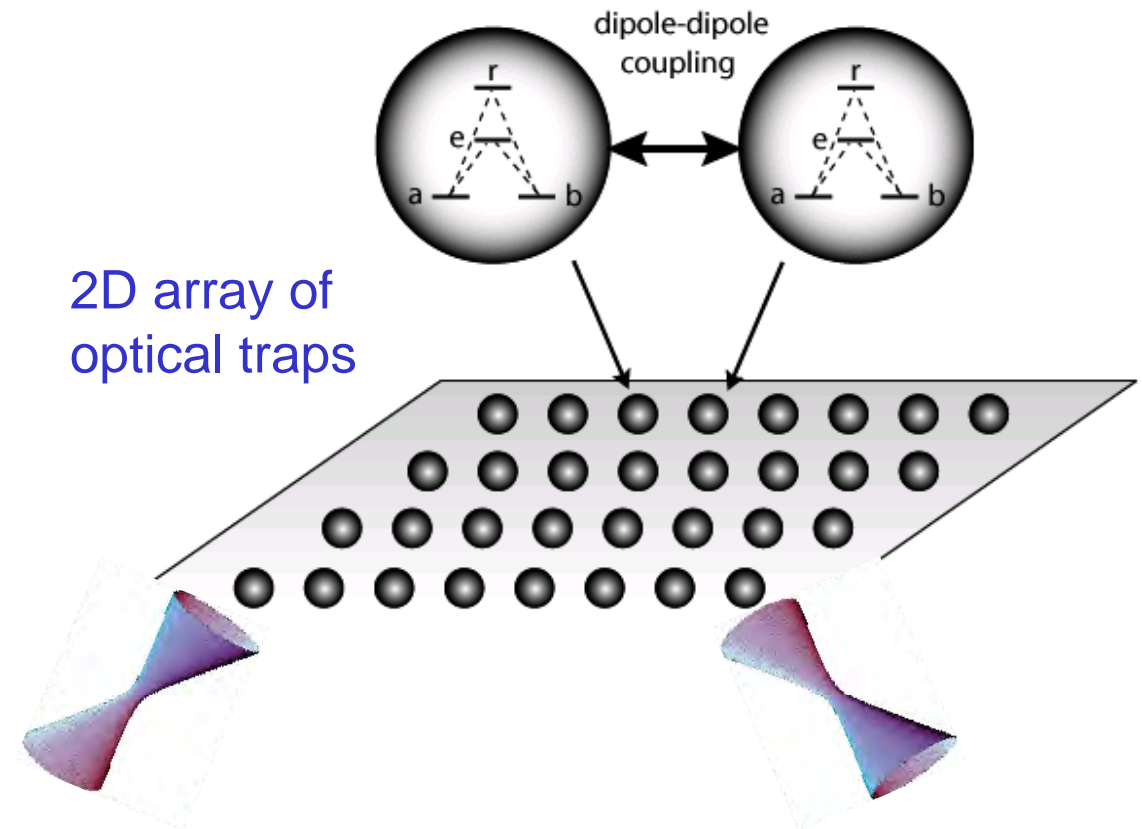
Neutral atom Rydberg gate array



hyperfine qubit

focused laser beams
for state preparation,
single qubit operations,
and measurements

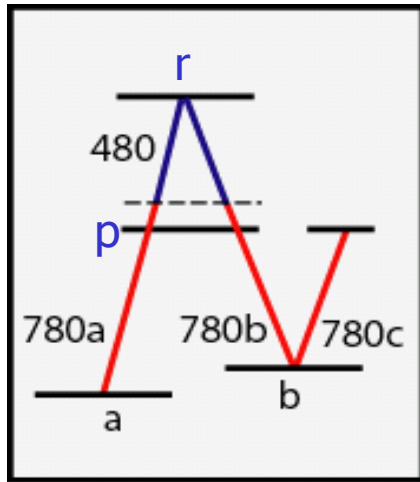
Yavuz, et al. PRL (2006)



focused laser beams
for Rydberg excitation
and two qubit gates

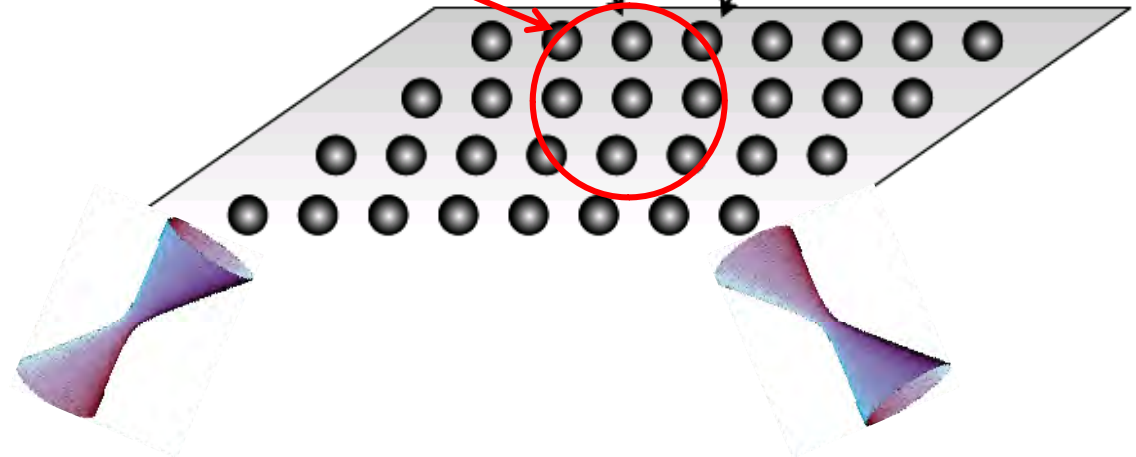
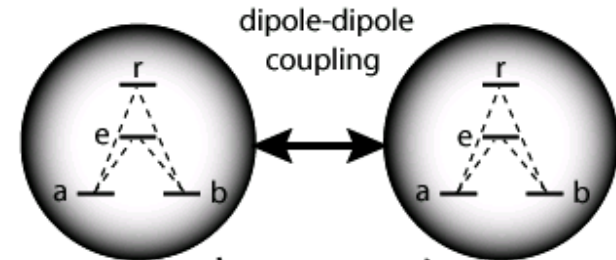
Johnson, et al. PRL (2008)
Urban, et al. NP (2009)

Neutral atom Rydberg gate array concept



hyperfine qubit

long range, not
nearest neighbor
interactions



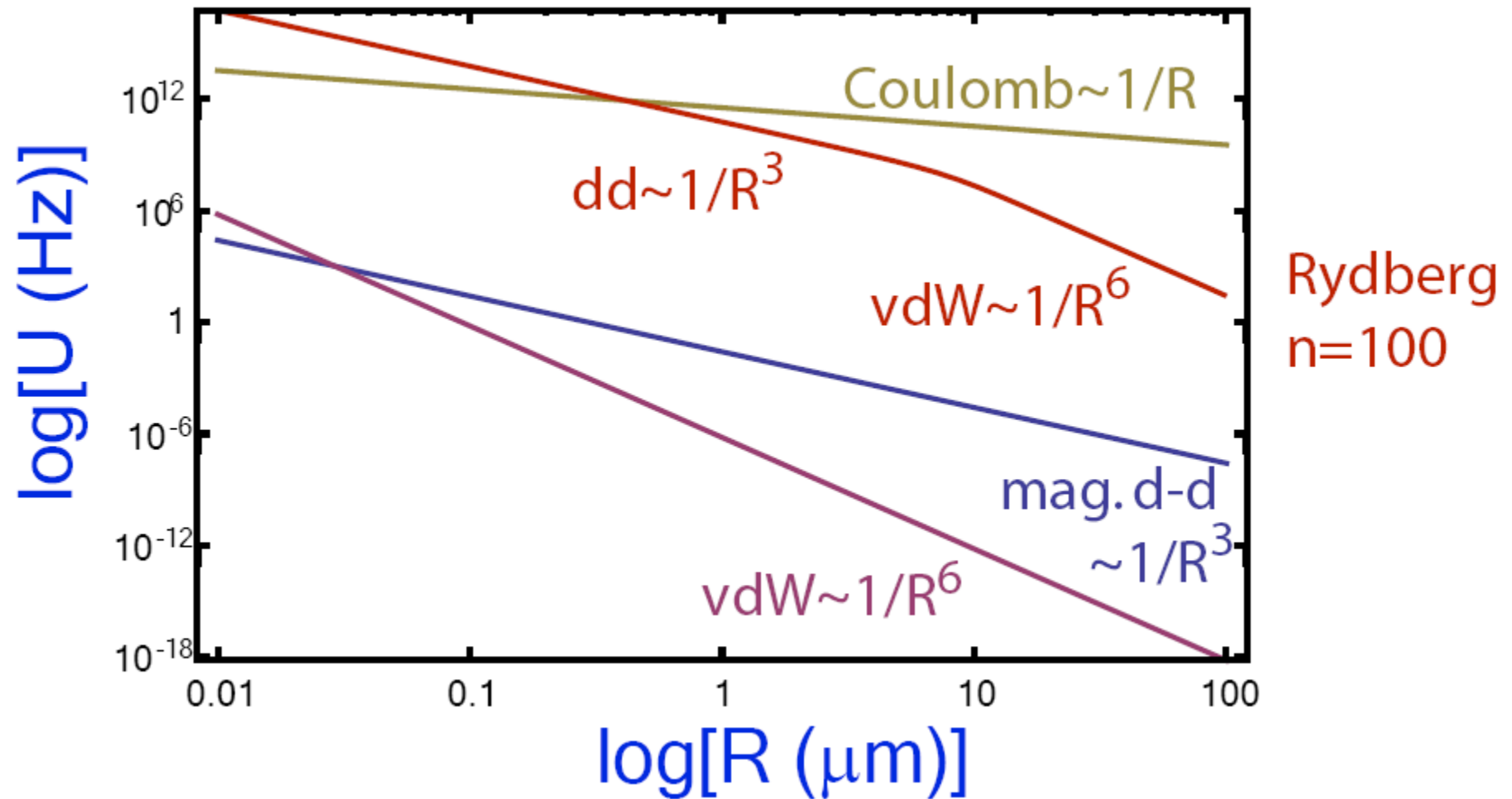
focused laser beams
for state preparation,
single qubit operations,
and measurements

focused laser beams
for Rydberg excitation
and two qubit gates

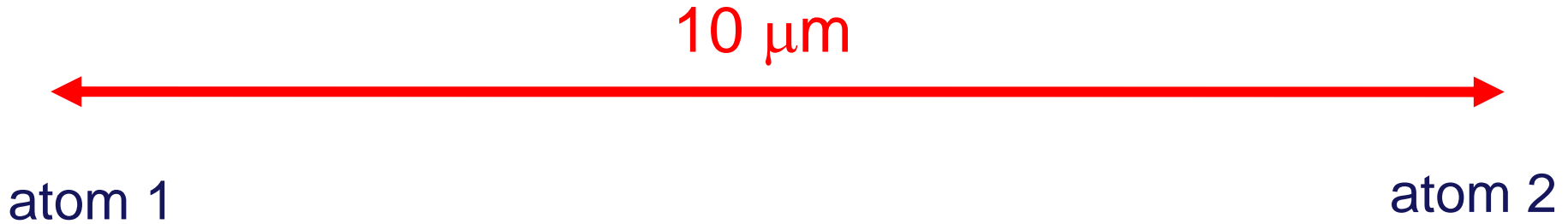
Yavuz, et al. PRL (2006)

Johnson, et al. PRL (2008)
Urban, et al. NP (2009)

Atomic interactions



Long range interactions between ground state atoms

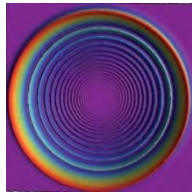
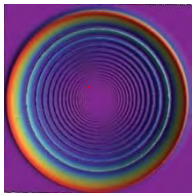


Rb-Rb magnetic dipole interaction

$$\Delta E \sim 25 \mu\text{Hz}$$

Long range interactions between Rydberg atoms

10 μm

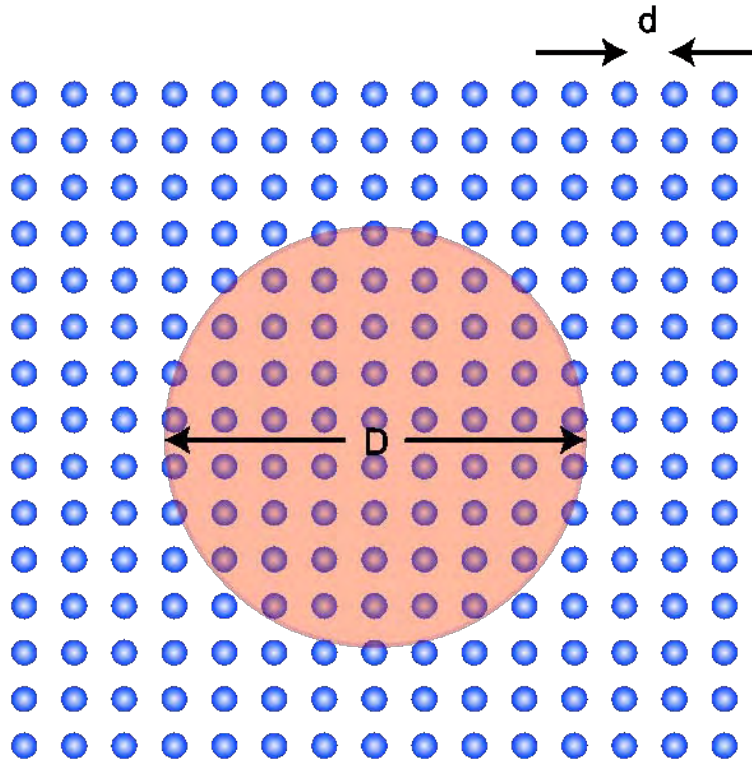


Rydberg $n=100$ van der Waals interaction

$$\Delta E \sim 25 \text{ MHz}$$

We have shown that this interaction can be controlled and used for manipulating quantum states of single atoms.

How many sites can be connected via Rydberg ?



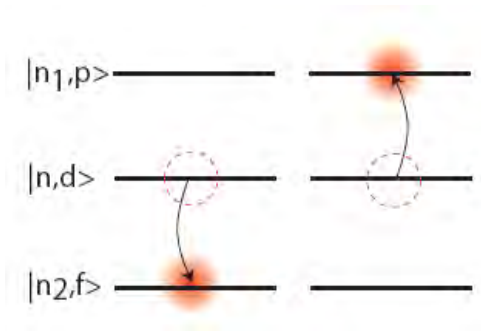
$$N = \frac{\pi}{4} \frac{D^2}{d^2}$$

How large can D be and
how small can d be?

Interaction strength

How large can D be?

$$U_{vdw} = \frac{C_6}{D^6}, \quad C_6 \sim \frac{\langle d \rangle^2 \langle d \rangle^2}{\delta} \sim \frac{n^4 n^4}{1/n^3} \sim n^{11}$$



Numerical calculations for Rb, Cs in the range $50 < n < 100$ show that

$$\delta \sim 1/n^3 \rightarrow \delta \sim k_\delta / n^4$$

due to $|\nu_s - \nu_p| \sim 0.5$. Also we can tune δ with external fields.

So

$$C_6 \sim n^{12}$$

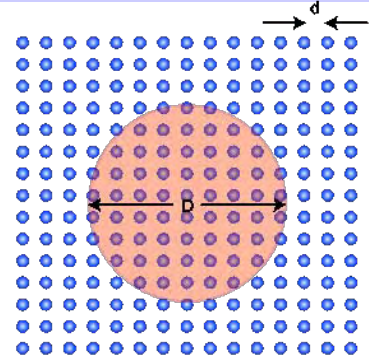
At constant gate error E , and Rydberg lifetime τ $\tau = \tau_0 n^2 \sim n^2$

$$D_{\max} \sim (C_6 \tau E)^{1/6} \sim n^{7/3}$$

Maximum number of connected sites

The minimum separation is set by the requirement that a Rydberg atom does not collide with a ground state neighbor. So

$$d_{\min} = k_1 a_0 n^2 \sim n^2$$



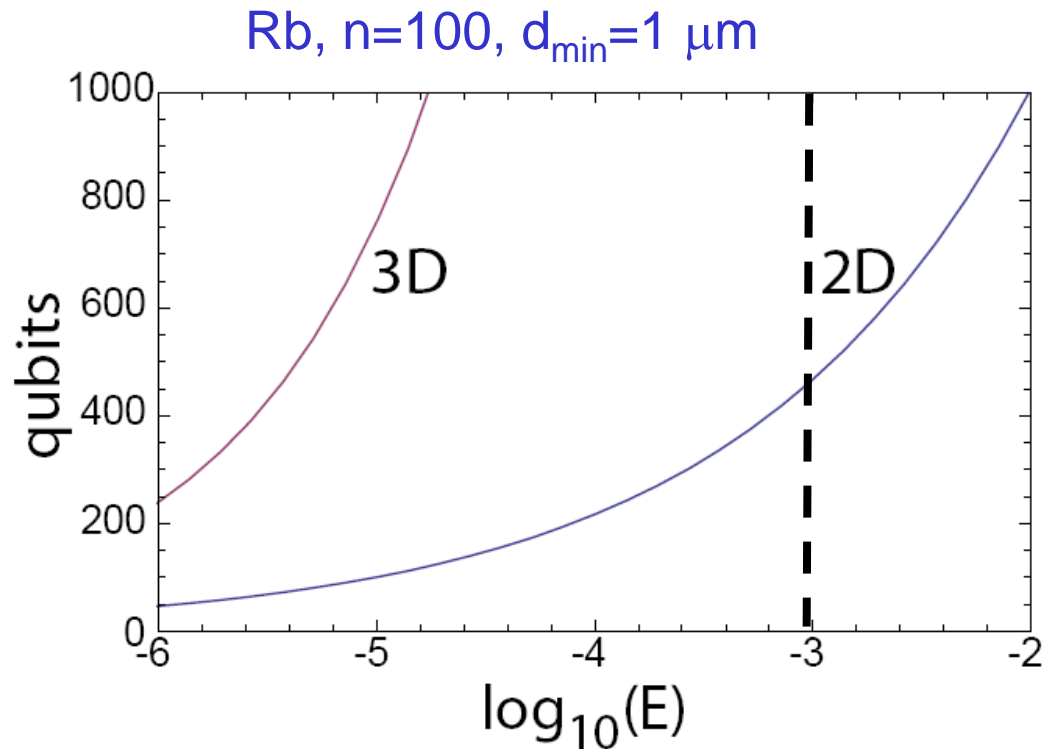
Combining the above we find

$$N_{\max}^{(2D)} = \frac{3\pi^{2/3}}{2^{8/3} k_1^2 k_\delta^{1/3}} \left(\frac{\alpha^2 m c^2 \tau_0}{\hbar} \right)^{1/3} E^{1/3} n^{2/3} \sim n^{2/3}$$

$$N_{\max}^{(3D)} = \frac{3^{1/2} \pi^{1/2}}{4 k_1^3 k_\delta^{1/2}} \left(\frac{\alpha^2 m c^2 \tau_0}{\hbar} \right)^{1/2} E^{1/2} n \sim n$$

Saffman and Mølmer, PRA (2008)

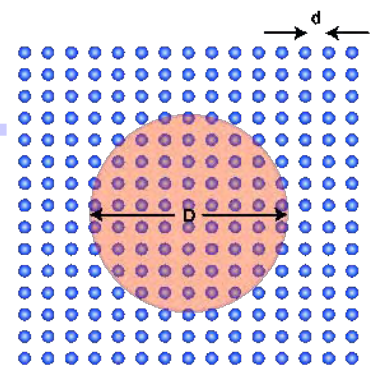
Connected qubits vs. gate error



At $E=0.001$

$$N_{\max}^{(2D)} = 470$$

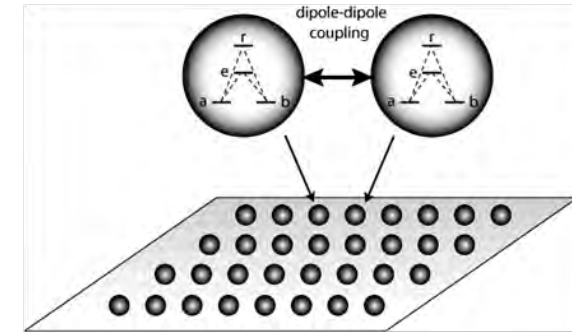
$$N_{\max}^{(3D)} = 7600$$



Challenges: optics, single atom loading,...

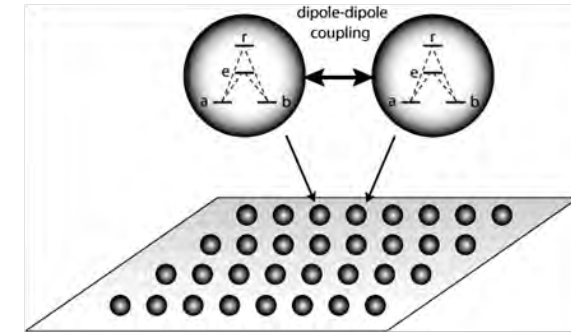
Single atom loading

A primary challenge is preparation of an array of singly occupied sites.



Single atom loading

A primary challenge is preparation of an array of singly occupied sites.



BEC - MI transition + transfer to long period lattice (Phillips, Porto) adds complexity

Stochastic loading: $P_1 \leq 1/e \sim 0.37$ does not use all sites

Collisional blockade: $P_1 \leq .5$ (Grangier) does not use all sites

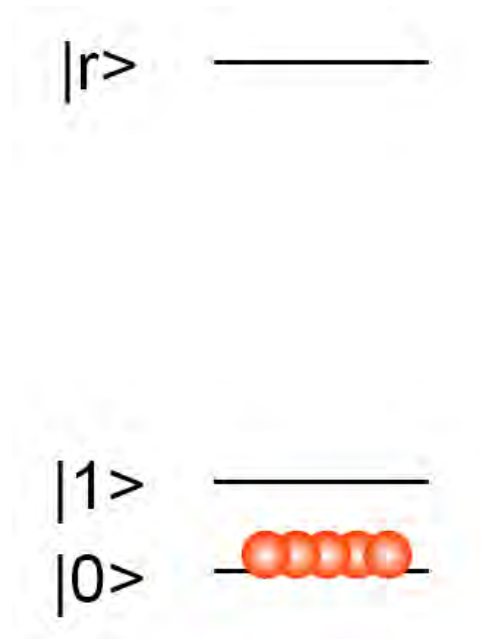
Light assisted collisions in 3D lattice $P_1 \leq .5$ (D. Weiss) does not use all sites

The stochastic loading methods will enable small arrays provided one is willing to throw away half the sites.

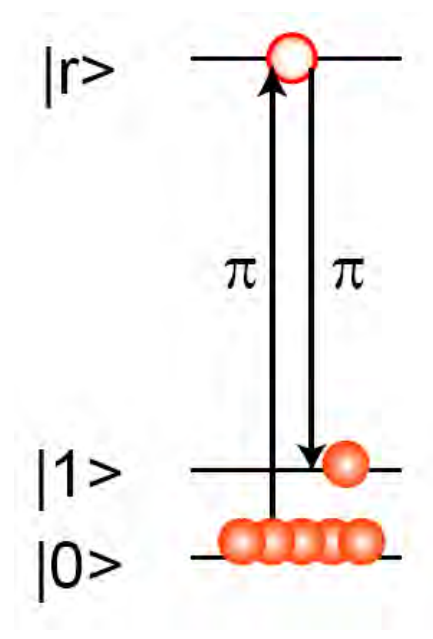
Deterministic loading enabled by entanglement: $P_1 \sim 1$ (Saffman, Walker (2002))

Ensemble qubits

Qubits can be represented by single atoms or by ensembles using Rydberg blockade.



Logical 0 $|\bar{0}\rangle = |0_1 \dots 0_N\rangle$



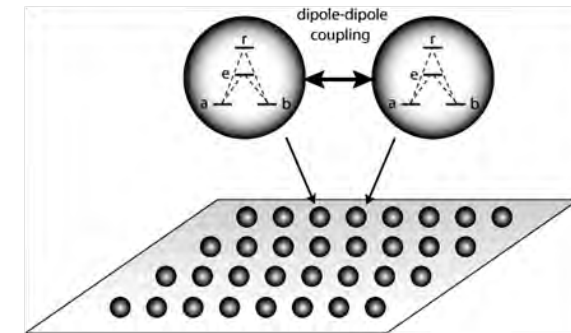
Logical 1 $|\bar{1}\rangle = \frac{1}{\sqrt{N}} \sum_j |0_1 \dots 1_j \dots 0_N\rangle$

“W” state with unit excitation
symmetrically shared among all N atoms

Lukin, et al. PRL (2001)

Ensemble qubits and single atom loading

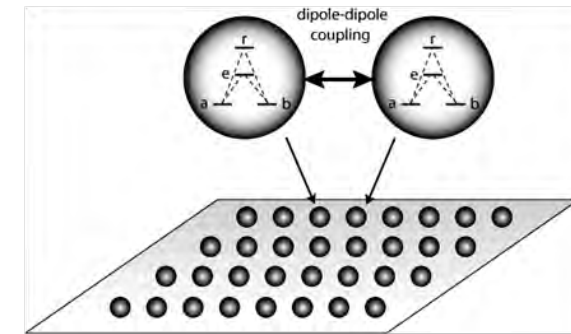
Protocol for deterministic single atom loading.



Ensemble qubits and single atom loading

Protocol for deterministic single atom loading.

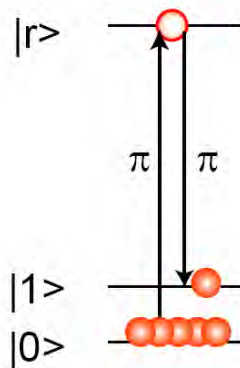
- load N atoms into optical trap
- pump all N atoms to $|0\rangle$
- transfer “1” atom to $|1\rangle$ via Rydberg
- eject $(N-1)$ atoms in $|0\rangle$ using radiation pressure



$$|\psi\rangle = |0_1 \dots 0_N\rangle \quad |\psi\rangle = \frac{1}{\sqrt{N}} \sum_j |0_1 \dots 1_j \dots 0_N\rangle$$

$|r\rangle$ ———

$|1\rangle$ ———
 $|0\rangle$ ———



$|r\rangle$ ———

—————
 $|1\rangle$ ———
 $|0\rangle$ ———

$$|\psi\rangle = |1\rangle$$

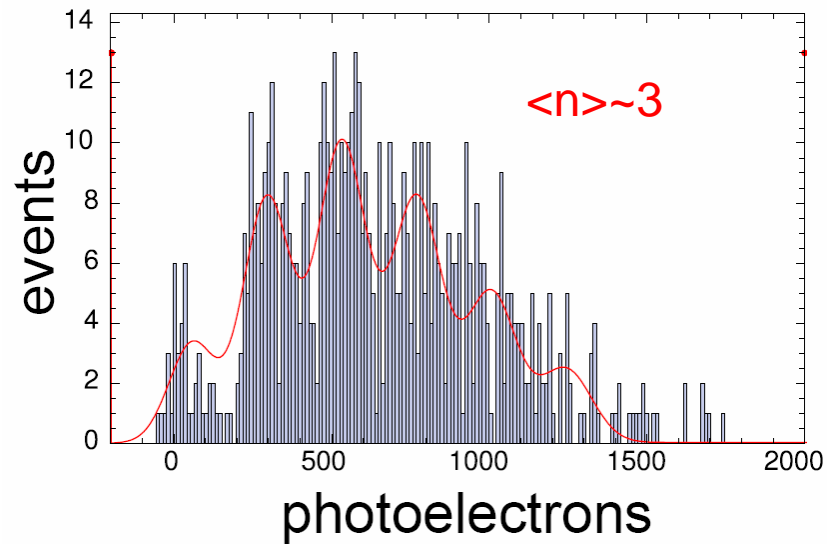
$|r\rangle$ ———

—————
 $|1\rangle$ ———
 $|0\rangle$ ———

Saffman and Walker PRA, (2002)

Ensemble qubits and single atom loading

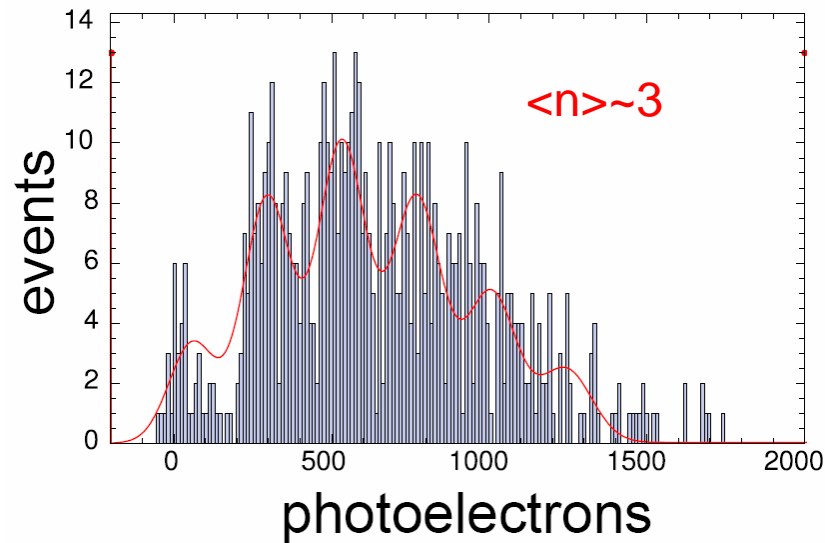
Preliminary results with this loading protocol have produced strongly sub-Poissonian Statistics.



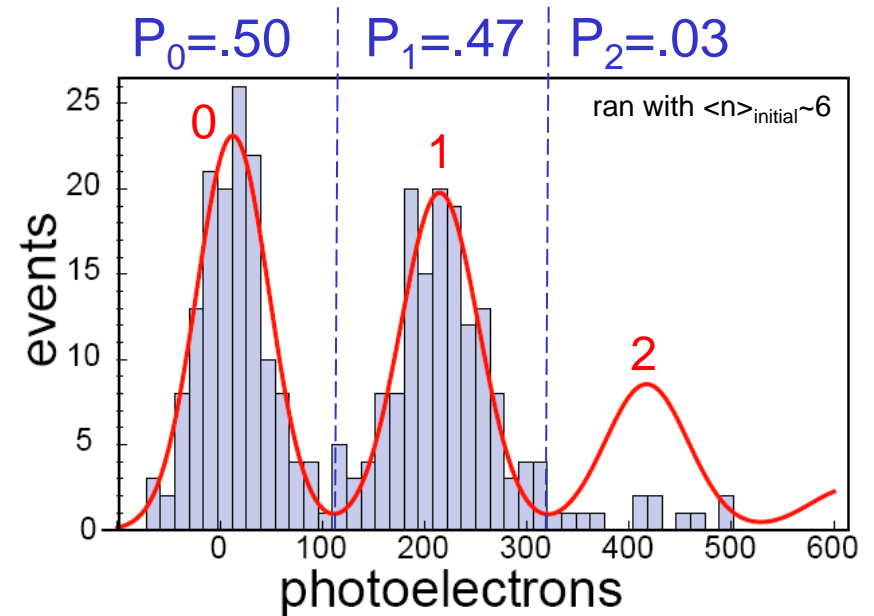
before

Ensemble qubits and single atom loading

Preliminary results with this loading protocol have produced strongly sub-Poissonian Statistics.



before



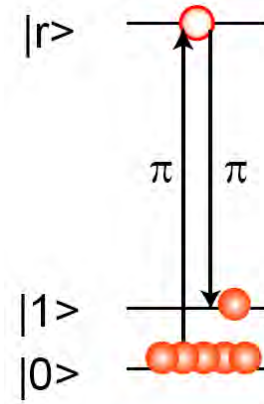
after

Optimization in progress....

N-atom entanglement

The “W” states can be created very naturally with Rydberg blockade

$$|\psi\rangle = \frac{1}{\sqrt{N}} \sum_j |0_1 \dots 1_j \dots 0_N\rangle$$



“W” states do not have the greatest possible degree of entanglement. Measuring one atom does not uniquely determine the state of all others.

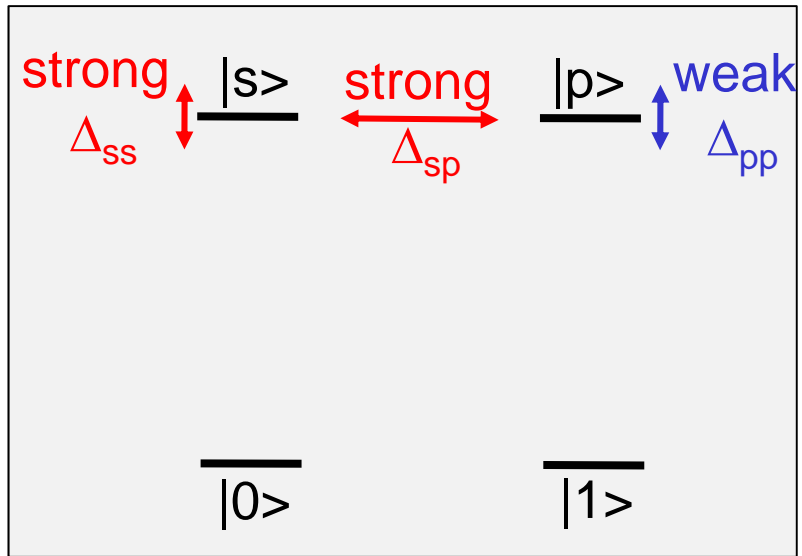
A maximally entangled N atom state is the Schrödinger cat or GHZ state

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|000 \dots 0_N\rangle + |111 \dots 1_N\rangle)$$

This can also be efficiently prepared with Rydberg blockade.

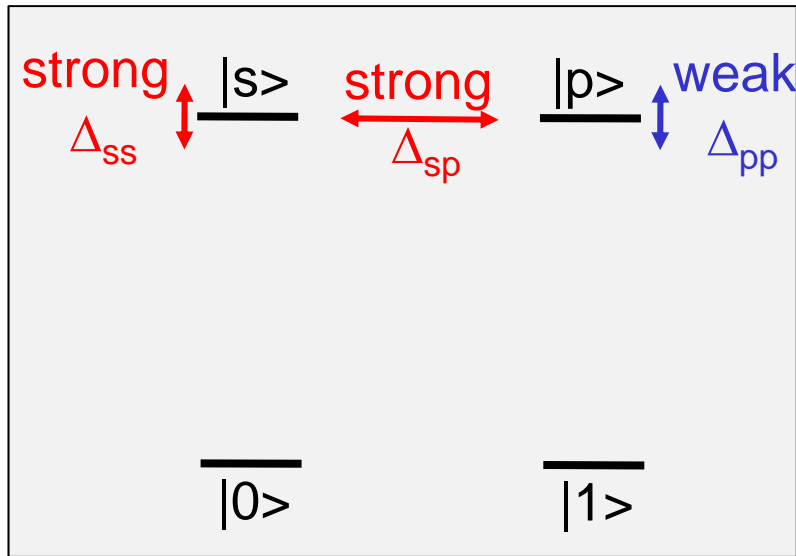
Generation of N-atom Schrödinger Cat states

Basic idea - exploit asymmetric Rydberg interactions:

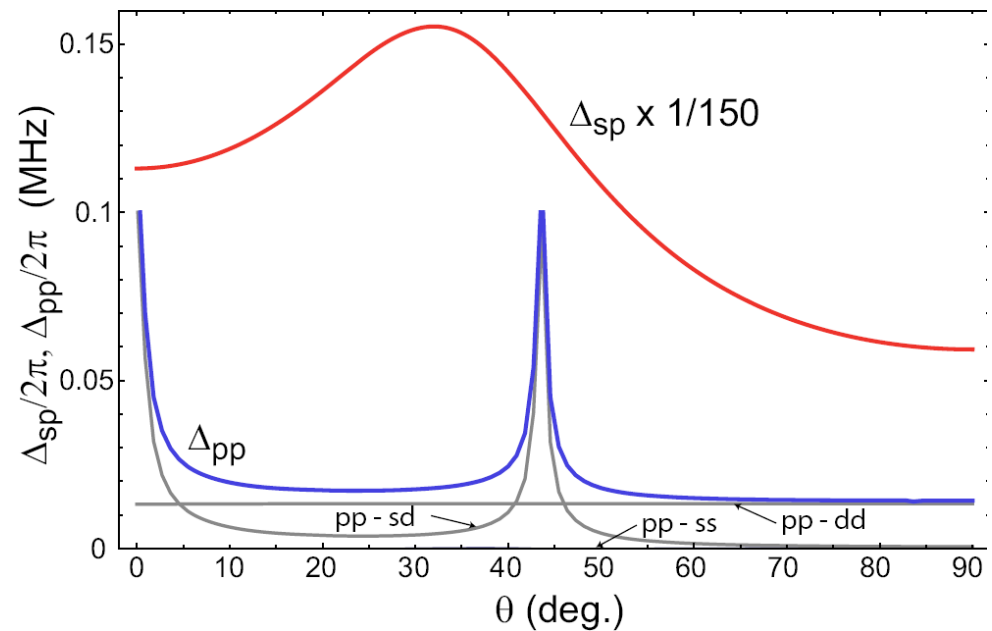


Generation of N-atom Schrödinger Cat states

Basic idea - exploit asymmetric Rydberg interactions:

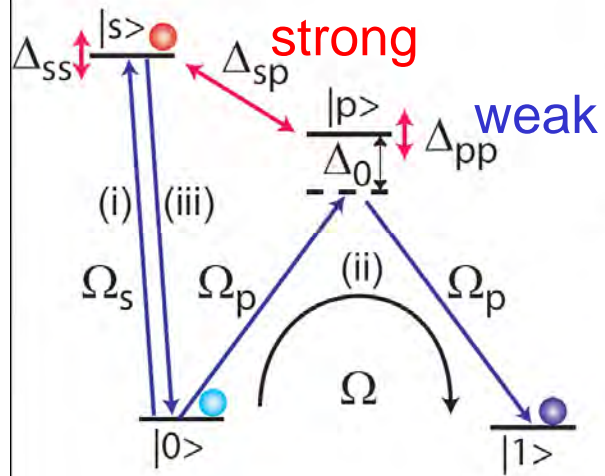


This is possible with a ratio > 150 .

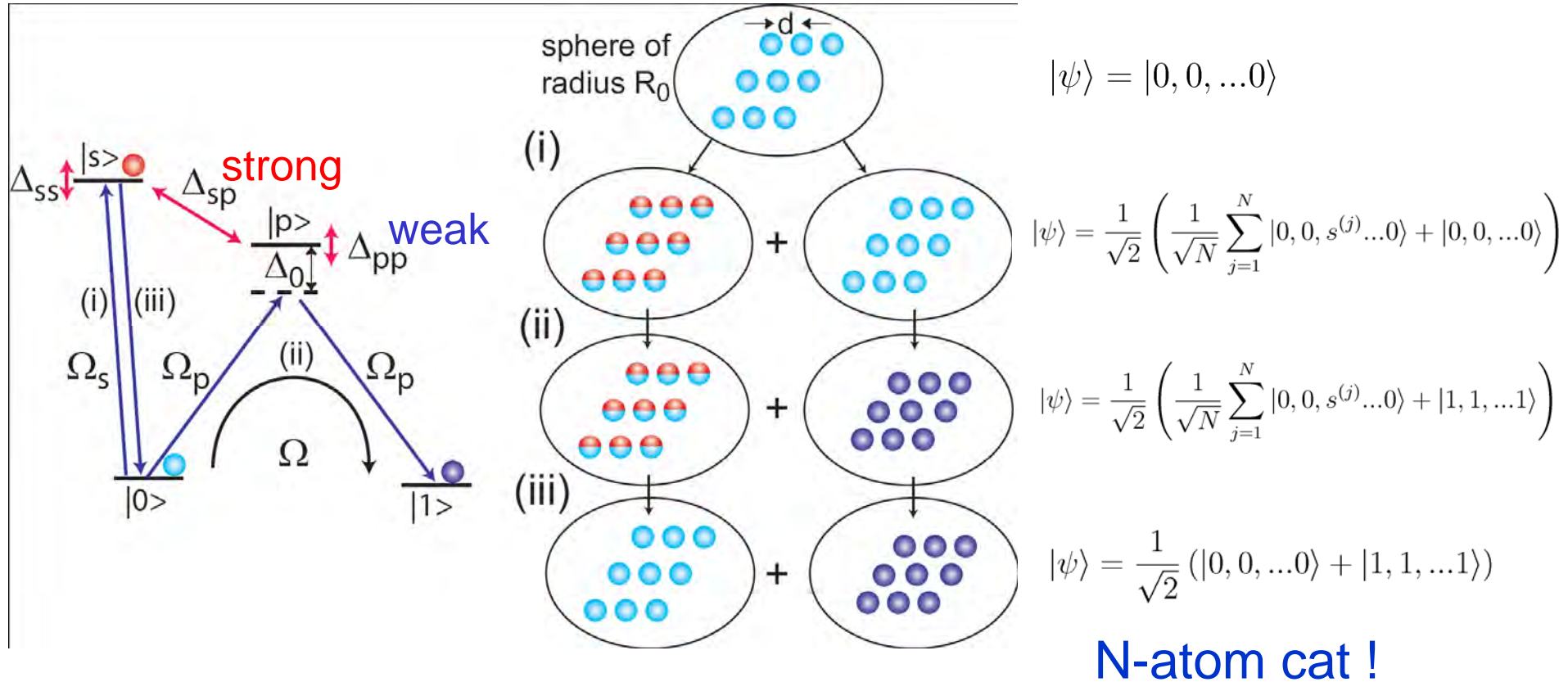


$$\begin{aligned} |s\rangle &= 41s_{1/2} \\ |p\rangle &= 40p_{3/2} \\ R &= 3 \mu\text{m} \end{aligned}$$

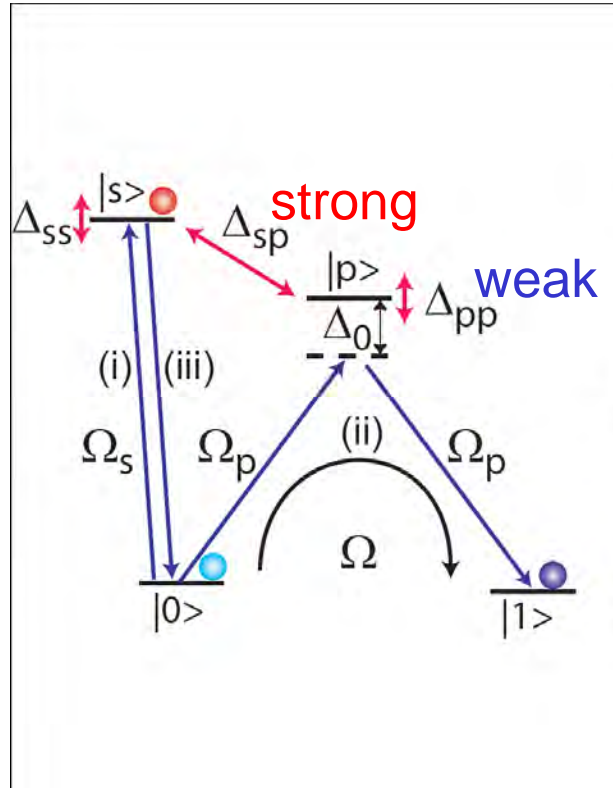
Generation of N-atom Schrödinger Cat states



Generation of N-atom Schrödinger Cat states

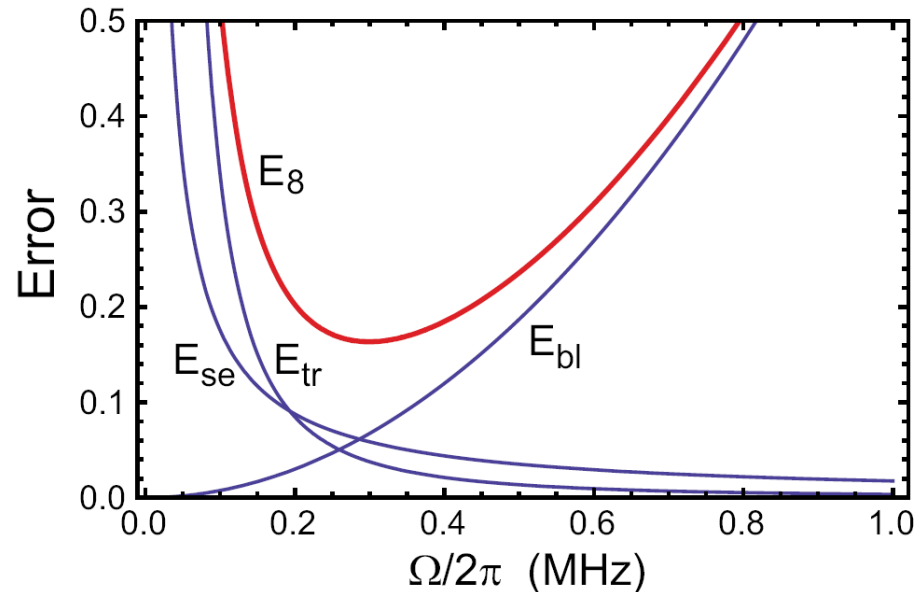


Generation of N atom Schrödinger Cat states



Fidelity of 8 atom
cat state

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|0, 0, \dots, 0\rangle + |1, 1, \dots, 1\rangle)$$

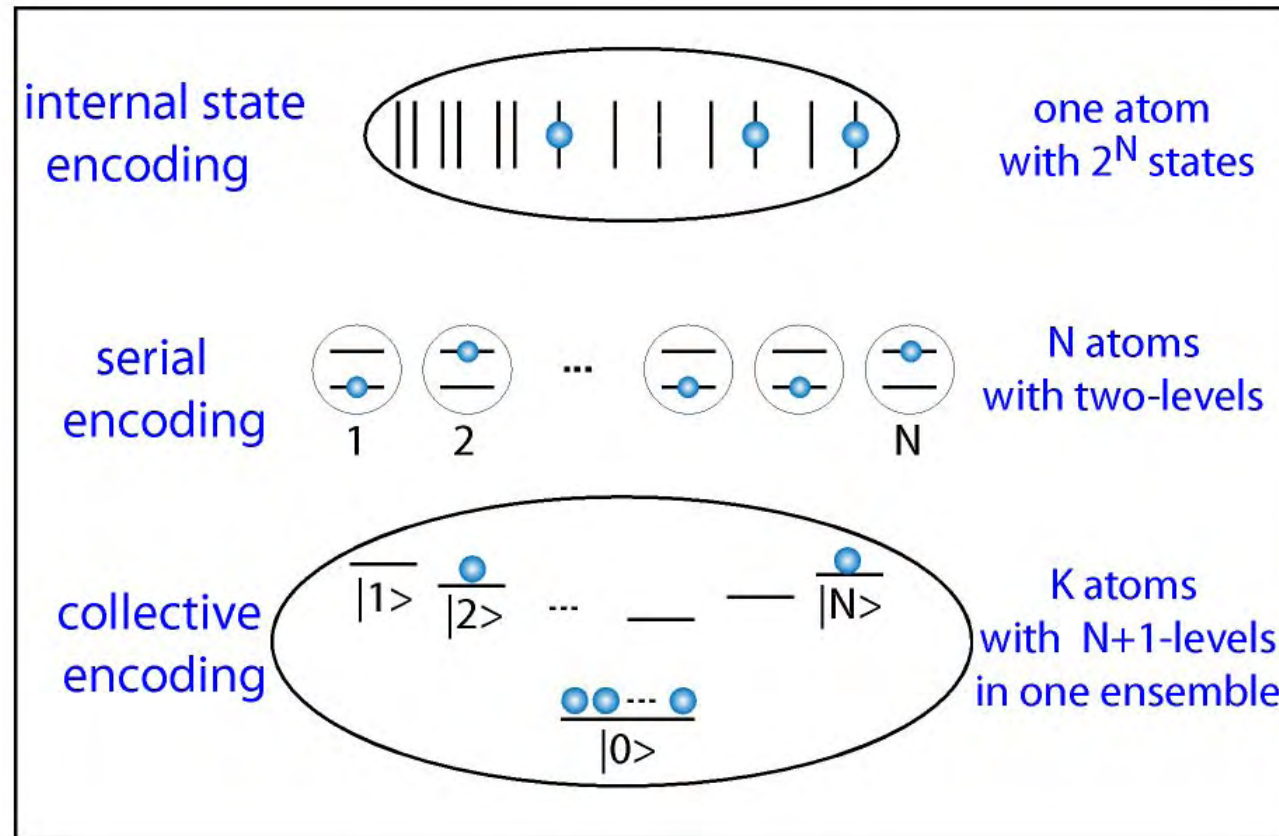


Numerical simulations and perturbation theory
show fidelity of 85% for $N=8$.

Saffman & Mølmer PRL to appear, arXiv:0812.2425

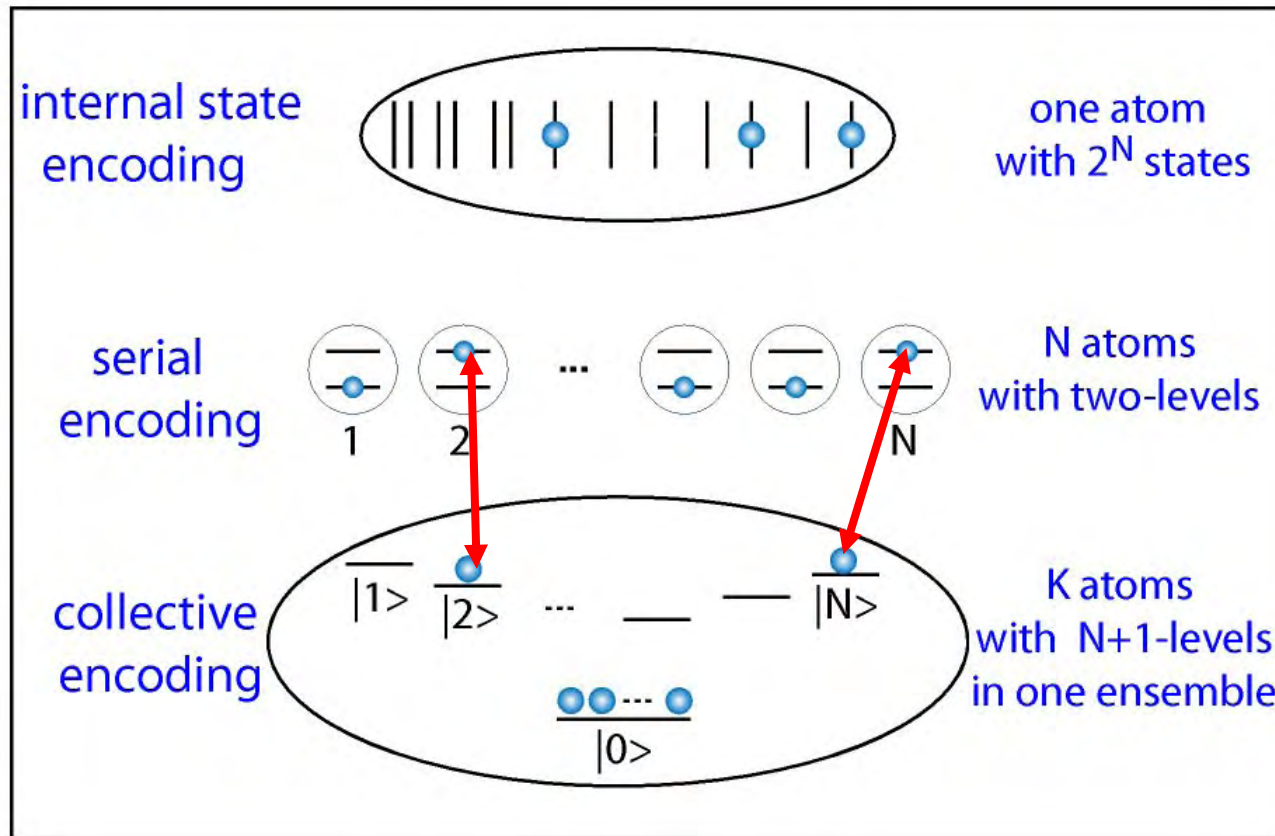
Müller et al., PRL (2009) ($N \sim 3$)

Collective encoding of an N qubit register



$K \geq N$ atoms, each with $N+1$ states all in one site
Less sensitive to atom loss – information is distributed.

Collective encoding of an N qubit register



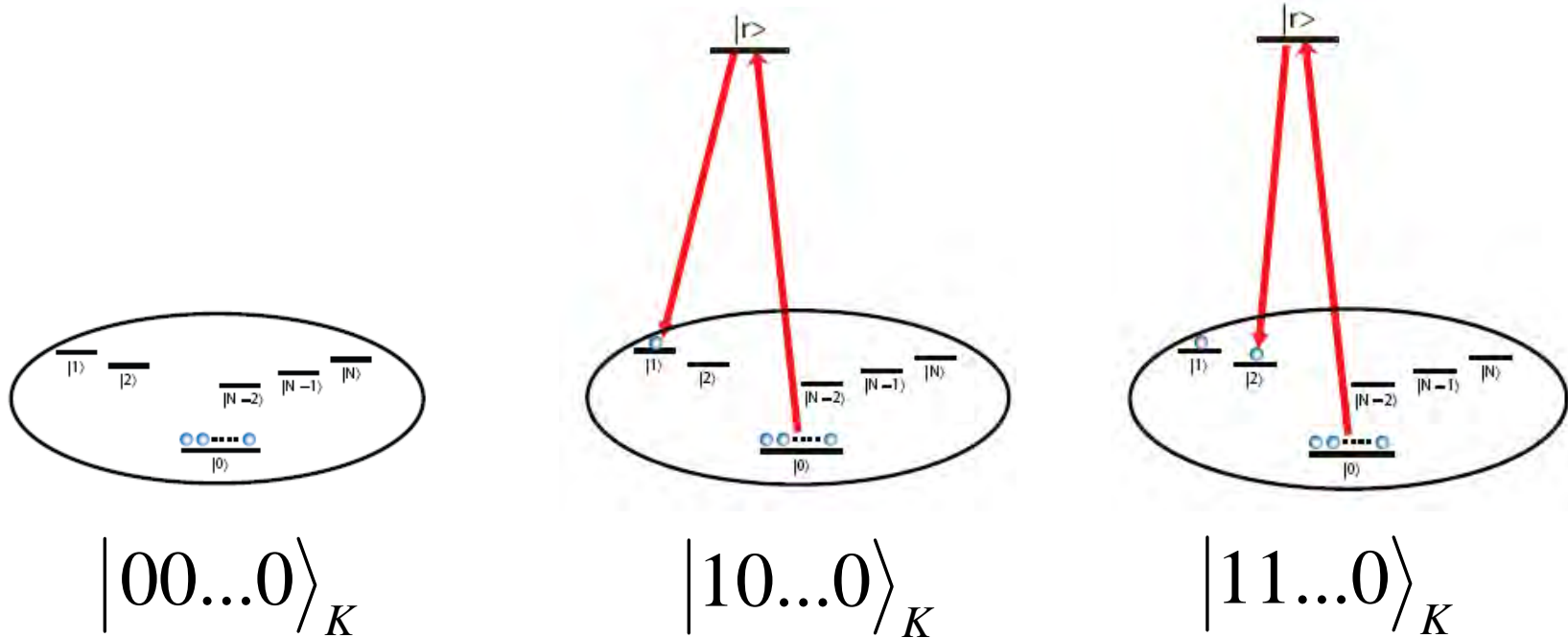
$K \geq N$ atoms, each with $N+1$ states all in one site

Less sensitive to atom loss – information is distributed.

E. Brion, K. Mølmer, and M. Saffman,
Quantum computing with collective ensembles of multi-level systems,
PRL (2007)

Collectively encoded states

Assume K atoms each with $N+1$ long lived ground states $|0\rangle, |1\rangle, |2\rangle, |3\rangle, \dots |N\rangle$



Note fiducial register states are now entangled. For example

$$\begin{aligned}
 |1100\rangle_{K=4} = (1/\sqrt{12})(& |1200\rangle + |1020\rangle + |1002\rangle + |2100\rangle + |0120\rangle + |0102\rangle \\
 & + |2010\rangle + |0210\rangle + |0012\rangle + |2001\rangle + |0201\rangle + |0021\rangle)
 \end{aligned}$$

These states can be prepared starting from $|0\rangle$ using ensemble dipole blockade Lukin, et al. (2001).

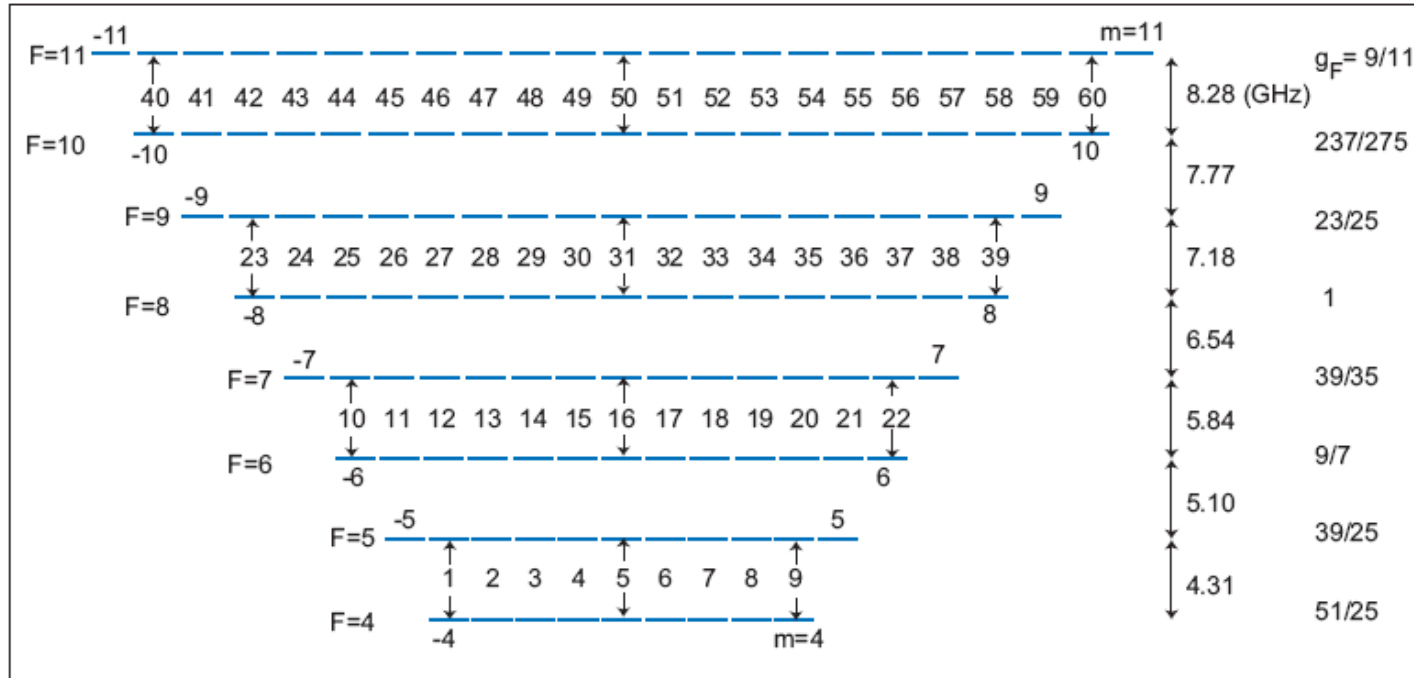
Collective encoding in Ho

With more internal states we can prepare a larger register.

atom	isotope	J	I	ground states	laser cooling demonstrated	BEC demonstrated
Yb	173	0	5/2	6	x	x
Sr	87	0	9/2	10	x	x
Rb	85	1/2	5/2	12	x	x
Al	27	1/2	5/2	12	x	
Cs	133	1/2	7/2	16	x	x
Tm	169	7/2	1/2	16		
Fe	57	4	1/2	18	x	
In	115	1/2	9/2	20	x	
Cr	53	3	3/2	28	x	x
Sc	45	3/2	7/2	32		
Tb	169	15/2	1/2	32		
Ti	49	2	7/2	40		
Mo	95	3	5/2	42		
Pr	141	9/2	5/2	60		
Tb	159	15/2	3/2	64		
Co	59	9/2	7/2	80		
Dy	161	8	5/2	102		
Er	167	6	7/2	104	x	
→ Ho	165	15/2	7/2	128		

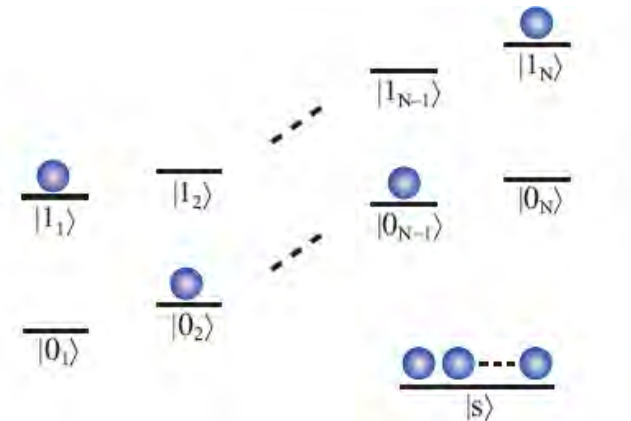
Collective encoding in Ho

Ground state $4f^{11}6s^2$, $J=15/2$ has 128 hyperfine states

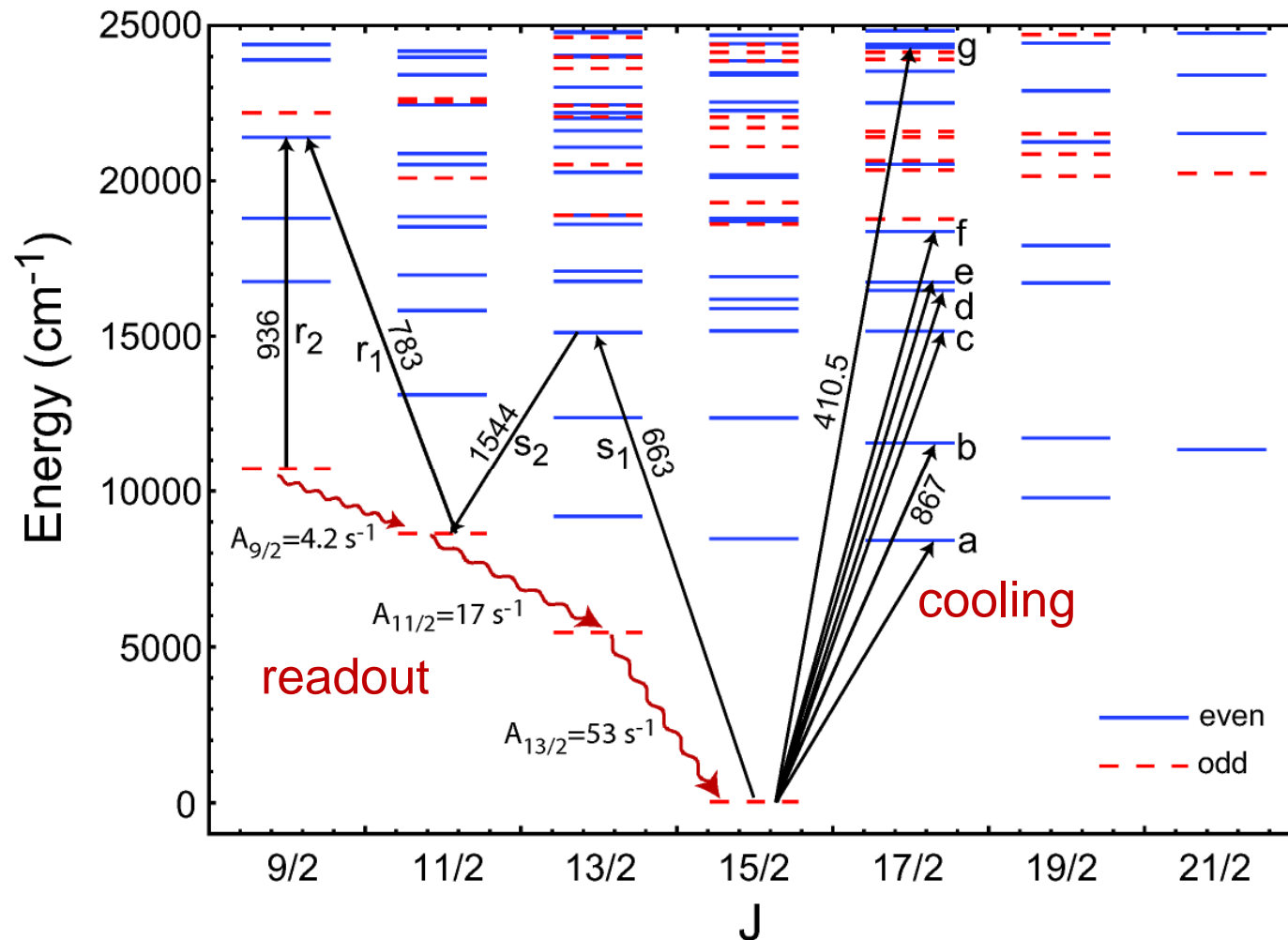


Encoding with two states per qubit gives a 60 qubit register.

E. Brion, K. Mølmer, and M. Saffman, PRL (2008)



Level structure in Ho



a-g cooling and trapping transitions
s1, s2 shelving
r1, r2 readout

Freezing out collisions

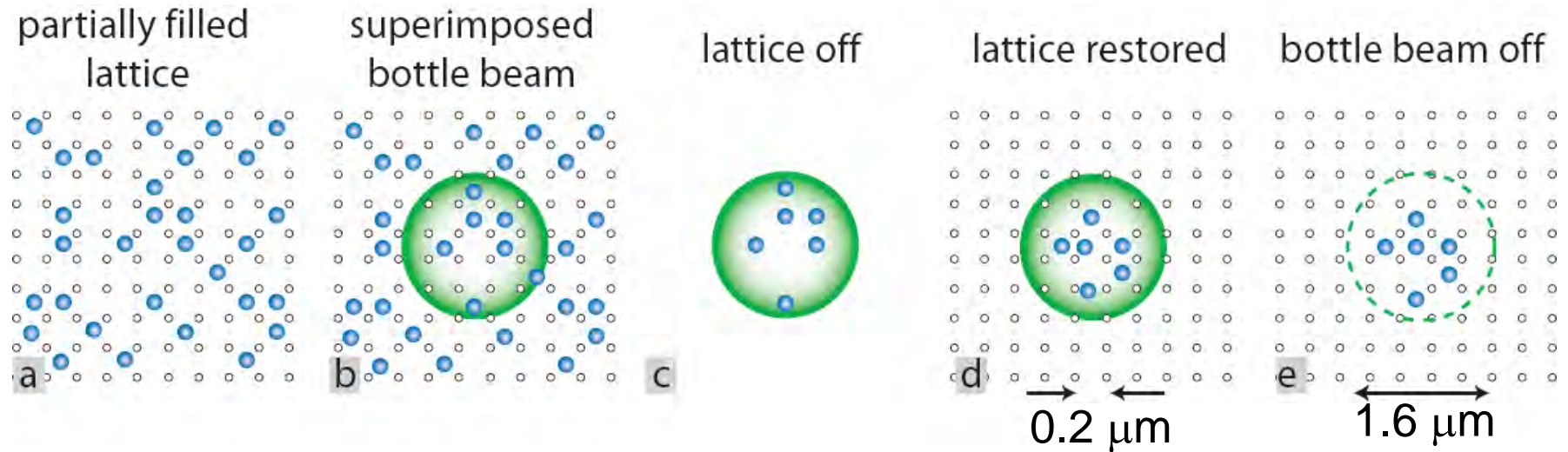


FIG. 6: Protocol for loading small ensembles into a patterned lattice region. The small open circles are the repulsive lattice sites, and the large green circle is the outline of the bottle-beam trap. The lattice is not drawn to scale for clarity. See text for details.

Freezing out collisions

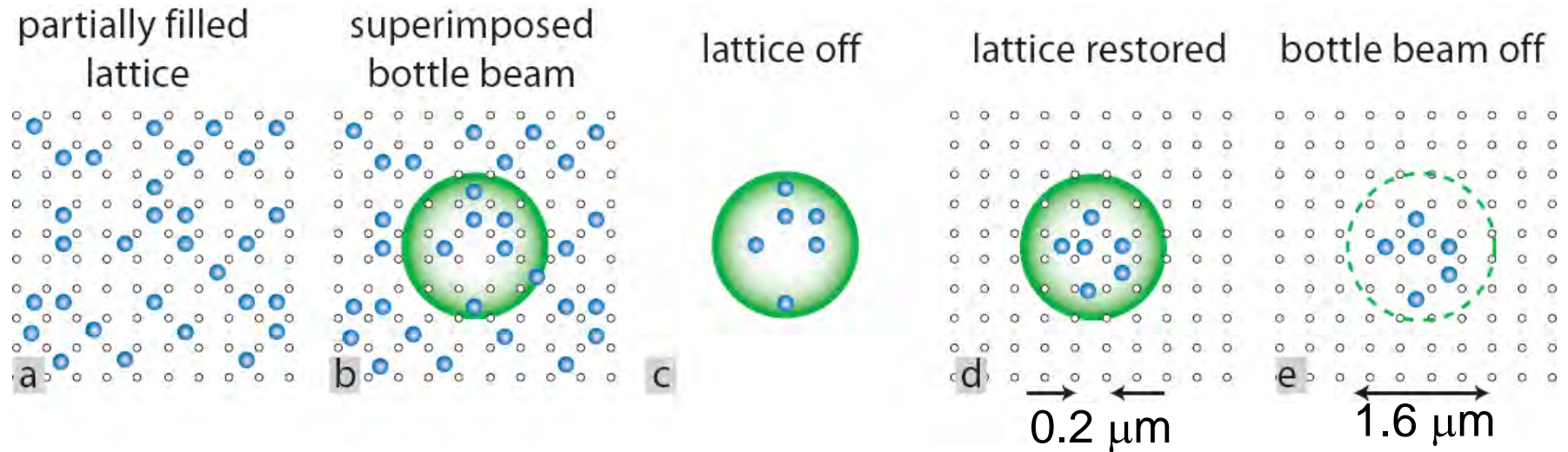
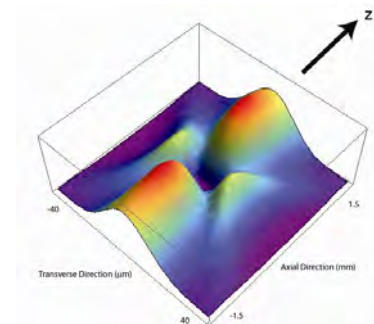
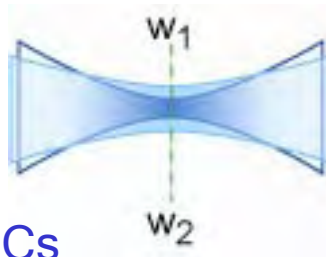


FIG. 6: Protocol for loading small ensembles into a patterned lattice region. The small open circles are the repulsive lattice sites, and the large green circle is the outline of the bottle-beam trap. The lattice is not drawn to scale for clarity. See text for details.

Bottle beam,
demonstrated with Cs

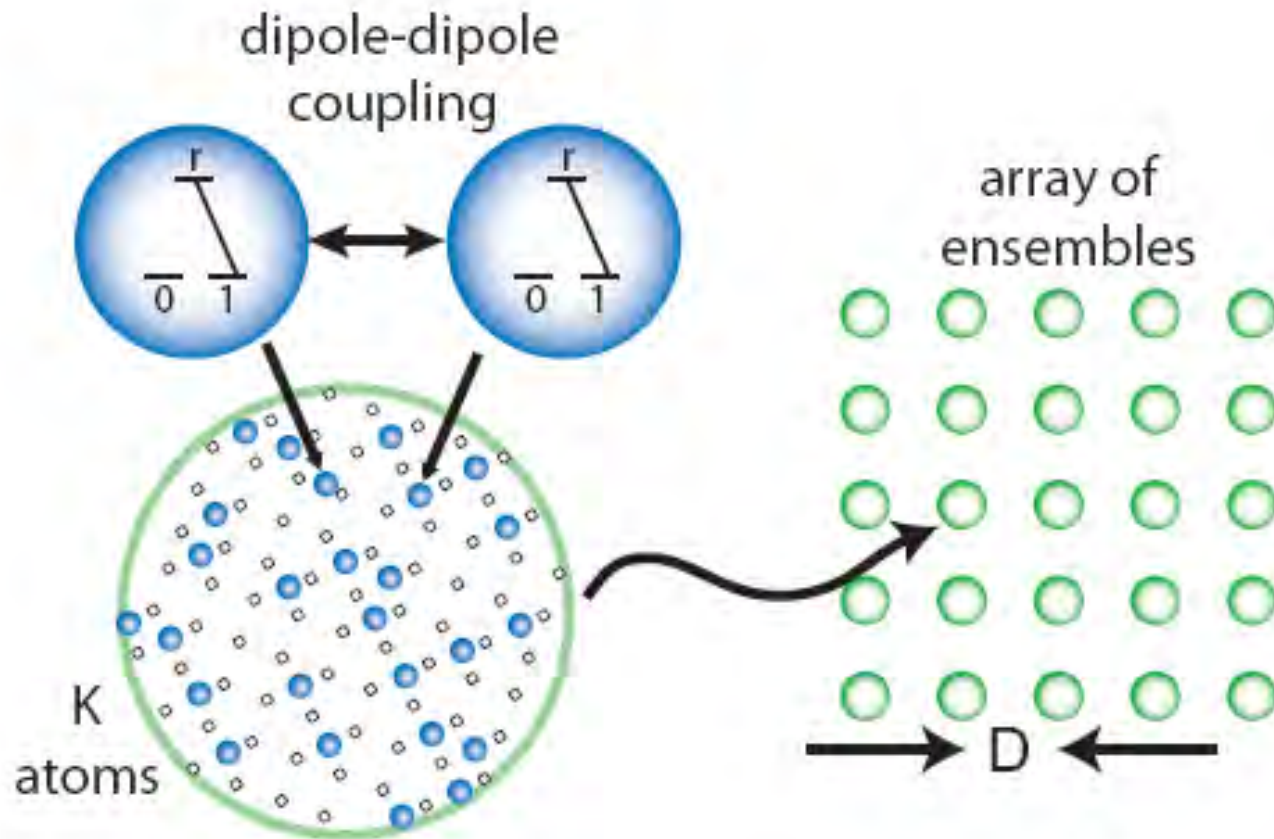


Isenhower, et al. Opt. Lett. 2009

1000 qubit scale processor

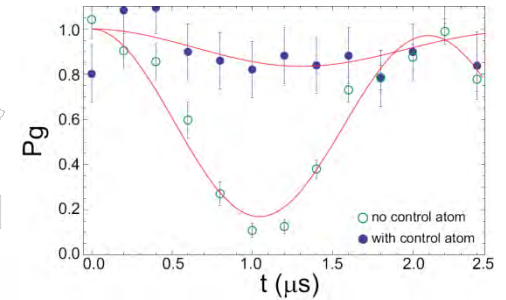
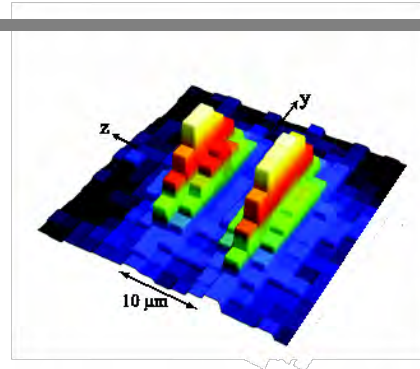
1000 qubit register based on:

18 sites, each with a 60 qubit ensemble with $K \sim 100$ atoms/ensemble



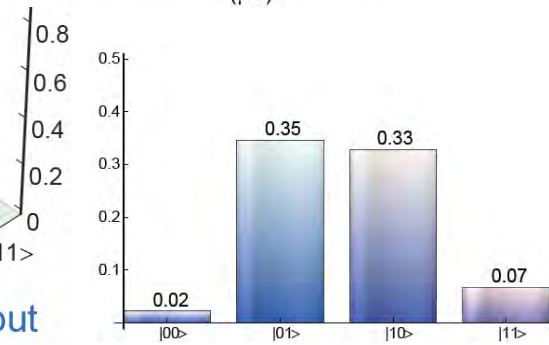
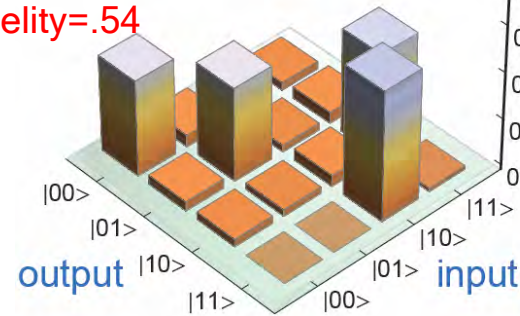
Summary

Rydberg blockade at $>10\ \mu\text{m}$

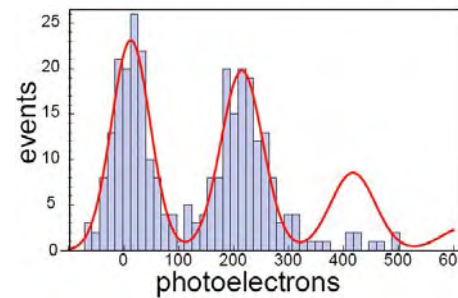


CNOT gate
entanglement

fidelity=.54



Single atom loading



collective encoding

